

SOME STUDIES on CRYSTAL and MOLECULAR STRUCTURE.

A thesis submitted for the degree of

Doctor of Philosophy

at the University of Glasgow

by

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### Summary.

The work which is described in this thesis is principally concerned with X-ray crystallography. Some other work, which involved molecular orbital calculations, was also carried out and is described in Appendix I.

Three compounds in all were studied by X-ray diffraction techniques, and the results of these investigations are summarised below.

(1) Dimethylcortierocin: This compound, which has the formula  $\text{CH}_3\text{O.CO.}(\text{CH}=\text{CH})_6\text{.COOCH}_3$ , and which is the dimethyl ester of an acid isolated from a fungus, crystallises as yellow needles from chloroform. The crystals are triclinic, space group  $C_1^1 - P 1$  or  $C_1^1 - P \bar{1}^*$ , and the unit cell dimensions are  $a = 9.06$ ,  $b = 7.52$ ,  $c = 6.07 \text{ \AA}$ ;  $\alpha = 104.8^\circ$ ,  $\beta = 104.0^\circ$ ,  $\gamma = 103.7^\circ$ . Unfortunately due to all the crystals which could be obtained being twinned, only the projection on (001) was studied. However this has been sufficient to show that the compound adopts the trans configuration throughout the molecule.

(2) n - Dodecylammonium Chloride: The samples of this compound were prepared, and all the photographs taken, in Sweden. Preliminary work was carried out on it in this department by C.J. McCron. The crystals are monoclinic, space group  $C_2^2 - P2_1$ , and the unit cell dimensions are  $a = 5.68$ ,  $b = 7.16$ ,  $c = 17.86 \text{ \AA}$ ;  $\beta = 92^\circ 16'$ . The projections on (010) and (100) have been studied and sub-cell theory has been applied to the calculation of the structure factors. The figures of merit are 15.8% and 21.2% respectively, the higher value being due to difficulties which were encountered in that projection.

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- |  |  |
|--|--|
| <p>(1) <u>Dimethylcortierocin</u>: This compound</p> | <p>Formula</p> <p><math>\alpha</math> of an acid</p> <p>on chloroform.</p> <p><math>\bar{1}</math>*, and the</p> <p><math>\alpha = 104.8^\circ</math>,</p> |
|--|--|

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space group  $C_1 - P1$  or  $C_2 - P1$ ,

ADD

though the results of the analysis definitely indicate the latter space group,

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The bond lengths etc. have been calculated and no abnormalities discovered. The standard deviations of these bond lengths are also given.

(3) L - Serine Phosphoric acid: This compound is an important amino acid which possesses an especially interesting feature in the phosphate group. The crystals are orthorhombic, space group  $D_2^4 - P2_1^2 2_1$ , and the unit cell dimensions are  $a = 7.79$ ,  $b = 10.24$ ,  $c = 9.09$  Å. The projections on (010) and (100) have been studied and a variety of methods was employed in an attempt to affect an analysis of the structure. Unfortunately none of these was successful. Consequently complete three-dimensional data have been collected and a three-dimensional Patterson synthesis has been computed on X-RAC. It is hoped that this will lead to a complete analysis of the structure being possible.

In Appendix I is an account of some molecular orbital calculations which were performed on dimethylcorticoic acid. These show that the molecule has a large amount of resonance energy and that this leads to the formal double bonds being considerable longer, and the formal single ones being considerably shorter, than "pure" double or single bonds respectively. The bond orders are compared with those in the corresponding conjugated hydrocarbon and a discussion of probably reactivity is also given.

Appendix II contains a short note on the re-determination of the film absorption factor for Ilford Industrial G film with  $Cu K_{\alpha}$  radiation.

Acknowledgements.

I should like to thank my two supervisors, Professor J.M. Robertson and Dr. V. Vand,\* for their generous help and cooperation during the progress of this research, and to thank Dr. T.H. Goodwin, who gave me a great deal of useful advice on how to write up the molecular orbital calculations.

My gratitude is due also to Professor Ray Pepinsky and his staff at the Pennsylvania State University who evaluated the sections of a three-dimensional Patterson synthesis on X-RAC, and to the Department of Scientific and Industrial Research for their maintenance grant, during the tenure of which this work was carried out.

\* Now at the Pennsylvania State University, State College, Pa., U.S.A.

## Introduction.

### 1. General:

It is less than fifty years since X-rays were first employed to study matter in the solid state, and during this period our knowledge of the intimate structure of solids has increased enormously. Before the advent of X-ray crystal analysis and other allied techniques it was not possible to make any direct study of the details of molecular structure as it exists in the solid state. Chemical studies could, in general, only be made if the solids were first vapourised, melted, or dissolved in a suitable solvent, and there could be no guarantee that this had not brought about some fundamental change. This position has now been completely changed due to these newer methods and, in certain instances, our ideas concerning the solid state revolutionised. Also certain theories, such as that of hydrogen bonding, have been confirmed in a way which would not have been possible by any other known method.

The study of organic crystals by these methods has enabled us to confirm very nearly all the structural deductions of the organic chemist and to differentiate between structures of equal chemical plausibility.<sup>1, 2.</sup> Further than this, accurate measurements have been made of a whole range of organic molecules,<sup>3, 4.</sup> and attempts are now being made to elucidate the structures of molecules so complex that they have not yet yielded to the ordinary methods of chemistry.<sup>1, 5, 6, 7.</sup>

It is also a consequence of these studies that nowadays a great deal of attention is being paid to the more theoretical

aspects of chemistry, since the theoreticians have now got a sound basis on which to work.<sup>8</sup>

## 2. Methods of Structure Solution:

When we record the diffraction pattern of a crystal the knowledge of the relative phases of the reflections is lost. Consequently the process of structure solution cannot in general be direct. Various methods have been devised to assist in overcoming this difficulty, but none of these is infallible. An account of some which found application in the present work is given below.

### 2 (a). Trial and Error:

This is the simplest and oldest method. It consists of postulating a set of atomic positions which conform with space group requirements and which are chemically plausible. A general survey of the intensities is usually of great assistance in this.<sup>11</sup> The structure factors which are calculated on the basis of these positions ( $F_c$ ) are then compared with the observed ones ( $F_o$ ), which are obtained from the measured intensities. The overall agreement is tested by means of the discrepancy or figure of merit (R) which is given by,

$$R = \frac{\sum |F_o - F_c|}{\sum |F_c|} \times 100\%$$

It is not possible to lay down cast-iron rules about the maximum acceptable value of R. The only reliable guide is individual experience.

We can only use this method when dealing with fairly simple or regular structures. If it were applied as outlined above

to complex ones it would involve a vast amount of computational work. The method has, however, been used with great success in many instances.<sup>3, 9.</sup>

## 2 (b) Utilisation of the Patterson Synthesis:

In this treatment we use only the observed data in order to evaluate a function  $P(UVW)$  which gives information about interatomic vectors. It is normally called the Patterson<sup>12</sup> synthesis and can be represented by the Fourier series,<sup>13</sup>

$$P(UVW) = \frac{1}{V_c} \sum_{-\infty}^{\infty} h \sum_{-\infty}^{\infty} k \sum_{-\infty}^{\infty} l |F(hkl)|^2 \cos 2\pi(hU + kV + lW)$$

Evaluation of this function for a unit cell of  $N$  atoms leads to  $N^2$  peaks, which represent weighted interatomic vectors and of which  $N$  are at the origin, so that there will normally be a great deal of overlapping due to the limited resolving power of the wavelengths employed. Consequently various modifying functions<sup>9</sup> have been proposed in attempts to improve the resolution, which will obviously be much poorer if the function is evaluated in only one or two dimensions instead of three.

It is possible in the course of structure analysis to employ the vector map which this method gives, because, as Wrinch first showed, we can always recover the fundamental set of points from its vector set.<sup>9</sup> In theory, therefore, it should always be possible to interpret a Patterson synthesis. Unfortunately, however, due to the complication of non-resolution of peaks, which has already been discussed, a Patterson synthesis does not yield a set of points, and a further difficulty arises from the fact that two or more fundamental sets of points, which Patterson has called

homometric, can have the same weighted vector set. This means that we cannot always obtain the vector set of points from the map and, when we can, we cannot distinguish between homometric pairs from the map alone.

When this method is employed in practice our task is somewhat simplified because we can always make use of additional pieces of information, such as bond lengths and known configurations. Nevertheless even with this help success cannot be guaranteed, although it seems probable that trial and error methods based on inspection of the Patterson synthesis are the most powerful available at present. The usual application of the method are: <sup>9</sup>

(1) to locate a heavy atom, whence the analysis may proceed by Fourier methods if we use the phase constants calculated for this atom as a first approximation. <sup>14, 15</sup> (see Section 3 (a))

(2) to determine the fundamental set of points from superposition methods. <sup>16, 17.</sup> This process may be assisted by the presence of a heavy atom. <sup>5</sup>

(3) to determine interatomic vectors from known interatomic distances. <sup>18.</sup>

## 2 (c) Relationships Between Structure Factors.

In a crystal the electron density can never be negative and this alone is sufficient to impose limitations on the magnitudes and phases of the structure factors. <sup>9, 19, 20.</sup> Hence it is possible to derive relationships between structure factors which are known as Harker-Kasper inequalities. When only a centre of symmetry is present the most powerful of these is, <sup>23</sup>

$$\left| U_H \pm U_K \right| \leq \left( 1 \pm U_{H+K} \right) \left( 1 \pm U_{H-K} \right)$$

Here H and K refer to Miller indices and the U's to the sharpened structure factors on a unitary scale (i.e.  $F(000) = 1$ ). When further elements of symmetry are present more inequalities can be deduced.<sup>9, 13.</sup>

Statistical relationships between the signs (S) of the structure factors in centrosymmetrical crystals have also been discussed.<sup>21, 22, 23.</sup> If the unitary structure factors involved are about 1,5 times the root mean square value then the relationship,

$$S_H = S(S_{K_1} \cdot S_{H+K_1})$$

is probably, but not necessarily, true.

Unfortunately the practical utility of these relations is limited because of their insistence on strong unitary structure factors. In complicated structures, where the method would be of greatest use, the average value of the structure factors is small and there are few much greater than average. Nevertheless successful use has been made of the method.<sup>23, 24.</sup>

### 3. Methods of Refining Atomic Parameters:

In this section we shall assume that approximate atomic coordinates have been arrived at by some means or another and we shall consider how these can be improved.

#### 3. (a) The $F_o$ Synthesis:<sup>9</sup>

The electron density ( $\rho$ ) in a crystal is periodic and can be represented by the Fourier series,

$$\rho(xyz) = \frac{1}{V_c} \sum_{-\infty}^{\infty} h \sum_{-\infty}^{\infty} k \sum_{-\infty}^{\infty} l \left| F_o(hkl) \right| \cos. 2\pi \left\{ \frac{hx}{a} + \frac{ky}{b} + \frac{lz}{c} - \alpha(hkl) \right\}$$

which can be evaluated in one, two, or three dimensions.

In practice one evaluates the series using observed

coefficients  $\{ F_o(hkl) \}$ , and phase angles  $\{ \alpha(hkl) \}$  calculated for the postulated structure. If this has not been too wide of the mark a density distribution will be obtained which should enable us to choose improved coordinates. The process is now one of successive approximation to the true structure.

If there is a centre of symmetry present the complexity of the problem is greatly reduced since, if the centre of symmetry is chosen as origin, the phase angles can only have the values 0 or  $\pi$ . In these cases the principal thing is the determination of the correct signs and the refinement can proceed much more rapidly.

Even if all the phase angles included in the synthesis are correct, the coordinates normally obtained from an  $F_o$  synthesis are not the true ones. This is because the  $F_o$  values are subject to random errors of measurement as well as to certain systematic ones, and because the series may not contain all terms of appreciable magnitude due to the limitations of the wavelengths employed. It is possible to correct for errors caused by series termination and absorption, but not for extinction and random errors, although their effects may be minimised.

### 3. (b) The $(F_o - F_c)$ Synthesis: 9,25,26

We compute this synthesis from a Fourier series similar to that given in section 3. (a), but with  $(F_o - F_c)$  replacing  $F_o$  as its coefficients and normally only in two dimensions. It has properties which make it particularly useful in the final stages of crystal structure analysis as a complete picture of the causes of discrepancies between  $F_o$  and  $F_c$  is obtained.

From the resulting difference map we are able to decide on the magnitude and direction of the required atomic shifts, which are in the direction of the steepest ascents and which eliminate errors due to termination of series. Errors due to wrongly assumed atomic scattering factors and anisotropic thermal motion are also detectable as well as peaks due to hydrogen atoms. In the recent very accurate study of salicylic acid<sup>27</sup> the effects of bonding were even allowed for.

It should be emphasised that several ( $F_o - F_c$ ) syntheses are normally required before all these corrections can be properly allowed for, and that fine details will only show up if the initial intensities have been measured with sufficient accuracy. In the most accurate two dimensional work Geiger-counters are used for this purpose. 26, 27, 28.

#### 4. Sub-Cell Theory:<sup>29, 30.</sup>

Regions of structural periodicity are often to be found within a molecule, e. g. hydrocarbon chains and proteins. Thus, in the unit cell (axes  $a_k$ ) of their crystals, there may be one or more regions where the electron density is periodic and can be described by a sub-cell (axes  $A_j$ ). The suffixes mean that the symbols refer to all three axes,  $k$  being used for the main-cell and  $j$  for the sub-cell.

The equivalent positions in a sub-cell region are given by,

$$x_k = u_k + t_j A_j + X_j$$

where  $x$  and  $X$  refer to main and sub-cell parameters respectively,

$t_j$  enumerates the required sub-cell, and  $u_k$  gives the coordinates of the nearest corner of the block-shaped sub-cell region referred to the main-cell origin. The centre of gravity ( $x_m$ ) referred to the main cell origin is given by the similar expression,

$$x_m = u_k + \frac{1}{N_j} \sum (N_j - 1) A_j,$$

where  $N_j A_j$  gives the edges of the sub-cell region.

The sub-cell region has finite volume and therefore each of its reciprocal lattice points will be surrounded by an interference function,

$$I = \prod_j \frac{\sin \pi N_j H_j}{\sin \pi H_j},$$

where  $H_j$  refers to the sub-cell Miller indices. To refer the sub-cell diffraction amplitude to the origin of the main-cell, I must be multiplied by a modulating factor  $\exp. (2\pi i x_m h_k)$ ,

where  $h_k$  refers to the main-cell Miller indices. Then, since the contribution of the different parts of the main-cell to the structure factors are additive, we can write for  $F_h$ , the main cell structure factor, that  $F_h = F_h^o + \prod_j \frac{\sin \pi N_j H_j}{\sin \pi H_j} F_H \cdot \exp(2\pi i x_m h_k)$ , where  $F_h^o$  is the contribution of atoms not included in the sub-cell region and  $F_H$  is the structure factor of the sub-cell.

The main and sub-cell translations, diffraction orders, etc. can be interrelated by the same transformation matrix,<sup>13</sup> e.g.

$$(H_j) = (s_j^k)(h_k),$$

where  $s_j^k$  are the transformation coefficients.

When the number of non-sublattice atoms is small it is sometimes possible to use this theory in structure solution,<sup>30, 31</sup>

and a suggestion has been made that it might be possible to apply it to some very large molecules.<sup>3</sup> In our case, however, it was employed only to simplify the structure factor calculations.

### 5. Accuracy of Crystal Structure Determination:<sup>4</sup>

In order to discuss quantitatively the results of our investigations we must have some means of assessing the experimental errors, and proper significance tests to apply before drawing conclusions from comparisons. These errors result chiefly from errors in the  $F_o$  values and imperfections in the model used to calculate the  $F_o$ 's. Errors in cell dimensions and approximation errors in computation are normally of much less consequence.

Cruickshank<sup>32</sup> has studied these errors and has arrived at the following formulae for the standard deviation of electron density  $\sigma(\rho_o)$  and of the x-coordinate of the n-atom  $\sigma(x_n)$ ,

$$\sigma(\rho_o) = \frac{1}{V_c} \left\{ \sum_3 (F_o - F_c)^2 \right\}^{\frac{1}{2}}, \text{ and } \sigma(x_n) = \frac{2\pi \left\{ \sum_3 h^2 (F_o - F_c)^2 \right\}^{\frac{1}{2}}}{a V_c C_n}$$

Here  $C_n$  represents the central curvature,  $\frac{\partial^2 \rho}{\partial x^2}$ , at the centre of the n-th atom,  $\sum_3$  means summation over all the  $F_o$ 's including  $F(hkl)$  and  $F(\overline{hkl})$ , and the other symbols  $a$ ,  $V_c$ , and  $h$ , have their usual significance. With regard to reflections which are accidentally absent, Cruickshank<sup>32</sup> suggests leaving them out if series termination has been corrected for by an  $F_o$  { or the equivalent  $(F_o - F_c)$  } synthesis.

These standard deviations are liable to be an underestimate if many of the phases are badly wrong. They also require modification in the following circumstances:

(1) The structure is non-centrosymmetrical. Each must then be

multiplied by two if the structure does not have a centrosymmetrical projection<sup>33</sup> on a plane or line and by a value between one and two if it does.<sup>18</sup>

(2) The number of unknown parameters is not much less than the number of independent  $F_o$  values.<sup>34</sup>

The standard deviation of a bond  $\sigma(d_{12})$ , provided that

$$\sigma(x_n) \approx \sigma(y_n) \approx \sigma(z_n),$$

is given by

$$\sigma^2(d_{12}) = \sigma^2(x_1) + \sigma^2(x_2)$$

since errors at right angles to the bond do not affect its length.

The value of  $\sigma(d_{12})$  will be greater, however, should the atoms concerned be related by symmetry.

It is often desirable to be able to compare bond lengths, either experimental with theoretical ones or two sets of experimentally derived ones. In the latter case let the determinations be  $l_1$  and  $l_2$  with estimated standard deviations  $\sigma(l_1)$  and  $\sigma(l_2)$  respectively. Then, following the treatment of Cruickshank and Robertson,<sup>8</sup> we can write that,

$$t_o = \frac{(l_1 - l_2)}{\{\sigma^2(l_1) + \sigma^2(l_2)\}^{\frac{1}{2}}}$$

The following arbitrary conventions can now be applied to decide on the significance of any difference,

$t_o \leq 1,960$	not significant,
$1.960 < t_o \leq 2.576$	possibly significant,
$2,576 < t_o \leq 3,291$	significant,
$3.291 < t_o$	highly significant.

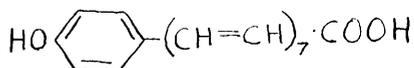
When a theoretical bond length is used for comparison the expression

is modified by considering that it has zero standard deviation.

The Crystal and Molecular Structure of Dimethylcorticrocin.

Introduction.

Corticrocin, a new pigment, was first isolated by Erdtman<sup>35</sup> from the mycelium of a mycorrhiza fungus. He reported it to be a polyene dicarboxylic acid, of formula  $\text{HO.CO.}(\text{CH}=\text{CH})_6\text{.COOH}$ , whose role in fungal metabolism is not yet known. Since then, cortisalin, another compound which possesses a polyene chain without isoprenoid methyl substitution, has been isolated from a fungus by Grippenberg.<sup>36</sup> It has the formula:



and it therefore seems possible that such compounds are of general importance among the fungi.

The acids of general structure  $\text{HO.CO.}(\text{CH}=\text{CH})_n\text{.COOH}$ , where  $n = 3, 4, 5,$  and  $7$ , were prepared by Kuhn and his collaborators.<sup>37</sup> Corticrocin itself, which is the missing link in this series, has also been synthesised recently.<sup>38</sup> None of these workers, however, were able to establish the geometrical configuration of the compounds, although it appears very likely from the study of its ultra-violet absorption spectra that dimethylcorticrocin is trans throughout.<sup>35</sup> The present work shows that this is indeed the case.

The bond lengths etc. in corticrocin are also of great interest because of the conjugated chain. Unfortunately, due to the adverse crystal habit, it was only possible to obtain one two dimensional projection of the structure so that no measurements of the molecule could be made. Theoretical bond lengths have however been calculated, and are to be found in the appendix.

Experimental.

A supply of corticocin and of its dimethyl ester was obtained from Erdtman.<sup>35</sup> It was decided to work on the dimethyl ester rather than the free acid since the latter was either insoluble or only sparingly soluble in most solvents.<sup>35</sup> The ester was recrystallised from chloroform in the form of yellow needles by very slow evaporation. When examined under the polarising microscope these showed two extinction positions, with the exception of one part which remained dark all the time, ~~also marked dichreism.~~<sup>\*</sup> ~~This proved to be due to twinning.~~<sup>DELETE</sup>

A previous investigator (Dr. G. Todd) obtained c-axis rotation and oscillation photographs, as well as the corresponding zero-layer moving film, without interference from the twin. Unfortunately this has not been possible with the other axes. Accordingly attempts were made to cut the crystals but the fragments obtained still showed evidence of twinning. Different solvents and other corticocin derivatives were also investigated but the crystals obtained still <sup>\*</sup> showed ~~pronounced dichreism.~~<sup>DELETE</sup> This meant that only the projection of the structure on (001) could be examined because of the interference of the twin with the photographs about the other axes, and with the first-layer Weissenberg photographs about the c-axis.

Unit cell dimensions were determined purely from photographs about the c-axis which was itself measured from a rotation photograph. The remaining unit cell parameters were determined from the superposition of a first-layer precession photograph<sup>39</sup> on a zero-layer one. The precession photographs were taken using  $M. K_{\alpha}$  radiation ( $\lambda = 0.711 \text{ \AA}$ )

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Twinning was later found to be the cause of this difference.

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\* see opposite leaf

and the remainder with Cu K $\alpha$  ( $\lambda = 1.542 \text{ \AA}$ )

Crystal Data.

Dimethylcorticoic acid,  $\text{CH}_3\text{O.CO.}(\text{CH} = \text{CH})\text{.COOCH}_3$ , molecular weight 274.2, melting point  $225^\circ \text{C}$ , density calculated  $1.225 \text{ g.cm}^{-3}$ , found  $1.237 \text{ g.cm}^{-3}$ .  
Triclinic, elongated in the direction of the c-axis.

$$a = 9.06, b = 7.52, c = 6.07 \text{ \AA};$$

$$\alpha = 104.8^\circ, \beta = 104.0^\circ, \gamma = 103.7^\circ$$

No absent spectra, space group  $\text{C}_2^1 - P 1$  or  $\text{C}_2^1 - P \bar{1}$ .\* One molecule per unit cell.

Intensity data for the (h k 0) reflections were recorded on a moving film photograph using the Robertson multiple film technique.<sup>40</sup> Measurement was by eye estimation and Lorentz and polarisation corrections were applied. This gave a set of structure factors on a relative scale. The film used was Ilford Industrial G which has an absorption factor of 3.22 (see Appendix II)

Analysis of the c-axis Projection.

The true unit cell has a projection on (001) which is not particularly convenient so that one was chosen which has the dimensions,

$$a' = 16.76, b' = 7.28 \text{ \AA}, \beta' = 95.7^\circ$$

The advantage of this choice is that the molecule now lies along the centre of the projection unit, and that the asymmetric unit is one quarter of the projected area which is face-centred.

An approximate structure was obtained by trial and error. This was refined using first  $F_o$ , and then  $(F_o - F_c)$  syntheses.

\* see opposite leaf

and the remainder with Cu  $K_{\alpha}$  ( $\lambda = 1.542 \text{ \AA}$ )

Crystal Data.

Dimethylcortierocoin,  $\text{CH}_3\text{O.CO.}(\text{CH} = \text{CH})\text{.COOCH}_3$ , molecular weight 274.2, melting point  $225^{\circ} \text{ C}$ , density calculated  $1.225 \text{ g.cm}^{-3}$ , found  $1.237 \text{ g.cm}^{-3}$ .  
Triclinic, elongated in the direction of the c-axis.

IN LINE 8 AFTER

space group  $C_1^1 - P_1$  or  $C_1^1 - P\bar{1}$ ,

ADD

Though the latter space group is definitely indicated by the results of the analysis.

molecule

recorded

on a moving film photograph using the Robertson multiple film technique.<sup>40</sup> Measurement was by eye estimation and Lorentz and polarisation corrections were applied. This gave a set of structure factors on a relative scale. The film used was Ilford Industrial G which has an absorption factor of 1.22 (see Appendix II)

Analysis of the c-axis Projection.

The true unit cell has a projection on (001) which is not particularly convenient so that one was chosen which has the dimensions,

$$a' = 16.76, b' = 7.28 \text{ \AA}, \beta' = 95.7^{\circ}$$

The advantage of this choice is that the molecule now lies along the centre of the projection unit, and that the asymmetric unit is one quarter of the projected area which is face-centred.

An approximate structure was obtained by trial and error. This was refined using first  $F_o$ , and then  $(F_o - F_c)$  syntheses.

\* see opposite leaf

Hydrogen atoms, which were placed using a model combined with inspection of a difference map, were included in the later stages of the work.

The chief difficulty associated with this projection was due to the lack of resolution of individual atoms. This necessitated the use of  $(F_o - F_c)$  syntheses from an early stage and probably added to the amount of work necessary for refinement. The discrepancy R stands at 21.1% when all the observed reflections are taken into account and an  $F_o$  synthesis using signs calculated for the final atomic coordinates is given in Figure I. The positions of the carbon and oxygen atoms are indicated on this, while the positions chosen for hydrogen atoms are shown by small crosses. The atomic coordinates are given in Table I.

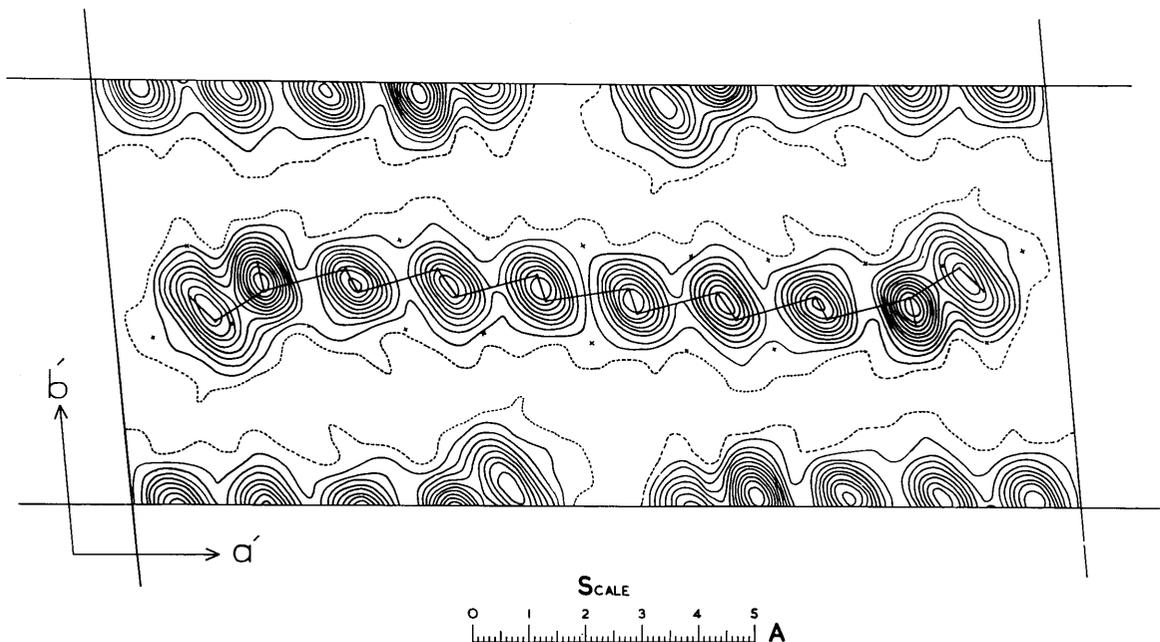


Figure I.

Electron density projected on (001). The contour line interval is one electron  $\text{\AA}^{-2}$ , the one electron line being broken.

Table I. Atomic Coordinates.

These are referred to the projection unit origin at centre of symmetry. Equivalent positions:

$$(x, y); (\bar{x}, \bar{y}); \left(\frac{1}{2} + x, \frac{1}{2} + y\right); \left(\frac{1}{2} - x, \frac{1}{2} - y\right)$$

Atom	$\frac{x}{a}$	$\frac{y}{b}$
$C_1$	0.045	0.016
$C_2$	0.051	0.955
$C_3$	0.138	0.008
$C_4$	0.155	0.942

Atom	$\frac{x}{a}$	$\frac{y}{b}$
C <sub>5</sub>	0.241	0.998
C <sub>6</sub>	0.252	0.943
C <sub>7</sub>	0.340	0.995
C <sub>8</sub>	0.419	0.012
O <sub>1</sub>	0.395	0.067
O <sub>2</sub>	0.344	0.927
H <sub>1</sub>	0.113	0.093
H <sub>2</sub>	0.000	0.885
H <sub>3</sub>	0.195	0.085
H <sub>4</sub>	0.099	0.868
H <sub>5</sub>	0.296	0.077
H <sub>6</sub>	0.192	0.874
H <sub>7</sub>	0.377	0.070
H <sub>8</sub>	0.462	0.107
H <sub>9</sub>	0.416	0.890

The atomic scattering curve used for both carbon and oxygen atom was a McWeeney<sup>41</sup> curve for an isotropic carbon atom with a temperature factor of  $\exp. \left[ -2.5 \left( \frac{\sin \theta}{\lambda} \right)^2 \right]$ . No allowances were made for the different scattering powers of oxygen and carbon etc. The curve for hydrogen was also an approximation, being a carbon curve similar to the above with a factor of  $\exp. \left[ -4 \left( \frac{\sin \theta}{\lambda} \right)^2 \right]$  and weighted for hydrogen. All these approximations probably contributed to raise the discrepancy, but would be unlikely to change any signs or affect the atomic positions by more than a very small amount. Since no bond lengths or angles have been measured this will not affect our conclusions.

### Discussion.

The electron density map obtained, and the reasonable agreement between observed and calculated structure factors, show that Erdtman's structural deductions<sup>35</sup> were correct and that the dimethyl ester does indeed have the trans configuration. It is disappointing, however, that no measurements of the molecule could be made due to all the crystals which could be obtained being twinned.

The atomic positions and the electron density map seem to indicate that there is a lack of regularity along the chain. No definite pronouncements on this are possible though, unless single crystals of this, or a related compound, can be prepared. An accurate study of the molecular configuration would then be possible.

Table II.

Observed and calculated structure Factors; values of  $F(h\ k\ 0)$ .Indexing for projection unit,  $F(h\ k\ 0)$  absent when  $h + k \neq 2n$ .

$h\ k$	$\sin\theta$	$F_o$	$F_c$	$h\ k$	$\sin\theta$	$F_o$	$F_c$
00	0	-	292.0	$\overline{12}, 2$	0.574	6.1	6.7
20	0.092	24.0	-18.3	$\overline{10}, 2$	0.489	42.3	-38.0
40	0.185	47.0	-33.8	$\overline{8}, 2$	0.408	13.5	12.4
60	0.277	10.5	-10.5	$\overline{6}, 2$	0.332	7.3	-11.0
80	0.370	10.9	9.9	$\overline{4}, 2$	0.268	33.6	-31.5
10,0	0.462	40.5	-41.4	$\overline{2}, 2$	0.223	31.9	-28.3
12,0	0.554	14.7	14.5	0 2	0.213	152.1	159.9
14,0	0.647	3.1	5.3	2 2	0.240	4.5	4.0
16,0	0.739	6.7	-6.0	4 2	0.296	9.4	-5.5
$\overline{15}, 1$	0.689	5.9	4.8	6 2	0.367	10.7	-8.8
$\overline{13}, 1$	0.599	6.3	-11.7	8 2*	0.447	6.9	1.2
$\overline{11}, 1$	0.509	13.8	-18.3	10,2	0.528	25.8	-31.6
$\overline{9}, 1$	0.419	6.9	-13.7	12,2	0.614	15.5	15.4
$\overline{7}, 1$	0.330	14.9	-4.9	14,2	0.701	5.2	-4.8
$\overline{5}, 1$	0.243	27.4	26.6	16,2	0.789	8.1	-6.5
$\overline{3}, 1$	0.166	46.5	36.5	18,2	0.878	< 2.5	-3.0
$\overline{1}, 1$	0.112	6.5	-4.7	20,2	0.967	7.4	9.5
1 1	0.120	10.6	26.3	$\overline{13}, 3$	0.690	6.6	-7.2
3 1	0.184	22.5	19.7	$\overline{11}, 3$	0.574	< 2.8	-1.5
5,1	0.266	23.8	20.9	$\overline{9}, 3$	0.489	< 2.5	19.2
7 1	0.351	5.9	6.0	$\overline{7}, 3$	0.431	13.1	-9.0
9 1*	0.440	3.9	-0.5	$\overline{5}, 3$	0.373	3.5	10.8
11,1	0.531	41.0	-38.7	$\overline{3}, 3$	0.334	23.7	25.6
13,1	0.621	7.1	-8.9	$\overline{1}, 3$	0.317	8.1	-11.1
15,1	0.712	8.9	6.5	13	0.327	20.2	28.3
17,1	0.802	< 2.8	3.0	3 3	0.360	< 2.0	-3.4
19,1	0.895	4.0	-3.4	5 3	0.413	< 2.2	3.7
				7 3	0.477	9.3	12.7

h k	sin $\theta$	F <sub>c</sub>	F <sub>o</sub>	h k	sin $\theta$	F <sub>c</sub>	F <sub>o</sub>
9 3	0.550	4.6	10.8	13,5	0.841	2.6	-4.4
11,3	0.628	38.7	-39.3				
13,3	0.708	9.0	-6.3	0 6	0.638	10.1	8.4
15,3	0.790	7.0	3.6	2 6	0.653	< 2.9	-7.6
17,3	0.875	< 2.5	2.7	4 6	0.681	< 3.0	7.1
19,3	0.963	2.9	-3.1	8 6	0.721	< 3.0	-2.9
				10,6	0.769	2.9	-3.7
$\overline{10},4$	0.595	13.2	-20.3				
$\overline{8} 4$	0.535	3.7	6.1	$\overline{9} 7$	0.813	6.2	3.4
$\overline{6} 4$	0.483	7.3	-7.7	$\overline{7} 7$	0.782	< 2.9	-2.8
$\overline{4} 4$	0.447	7.9	-12.4	$\overline{5} 7$	0.756	< 2.9	3.5
$\overline{2} 4$	0.426	9.8	-12.0	$\overline{3} 7$	0.743	< 3.0	4.6
0 4	0.425	39.8	38.0	$\overline{1} 7$	0.739	4.7	-0.7
2 4	0.444	< 2.3	7.4	1 7	0.748	5.1	4.9
4 4	0.481	6.4	9.0	3 7	0.770	< 2.9	1.6
6 4	0.531	< 2.6	-3.3	5 7	0.801	< 2.8	0
8 4	0.591	8.8	-3.3	7 7	0.841	< 2.6	2.9
10,4	0.659	6.6	-16.7	9 7	0.887	7.6	8.8
12,4	0.732	8.4	5.9				
14,4	0.809	2.7	-1.2	080*	0.850	2.3	0
16,4	0.888	2.6	-2.8				
$\overline{9} 5$	0.640	5.1	-7.8				
$\overline{7} 5$	0.594	< 2.8	-3.6				
$\overline{5} 5$	0.558	< 2.7	-0.4				
$\overline{3} 5$	0.536	< 2.6	11.0				
1 5	0.528	6.3	-6.0				
1 5	0.538	15.5	15.3				
3 5	0.564	< 2.7	-9.7				
5 5*	0.601	5.1	0				
7 5	0.650	< 2.9	0				
9 5	0.708	13.3	12.9				
11,5	0.772	22.6	-24.0				

\* These planes were left out of the Fourier shown in Figure I on account of uncertain sign. All non-observed planes were also left out.

The Crystal and Molecular Structure of n-Dodecylammonium Chloride.

Introduction.

Some preliminary details of this work have already been published by Gordon, Stenhagen, and Vand,<sup>42</sup> who detail the transitions undergone by the chloride on heating it to 75°C. They also give unit cell data and the results of a preliminary study of the projection on (010).

The N-mono-n-alkyl substituted ammonium halides are of interest because they exhibit surface activity and polymorphism, the latter property suggesting that the hydrocarbon chains can pack in different ways. It is also thought that the lower members exhibit some form of chain rotation since their crystals are tetragonal<sup>43,44</sup> at normal temperatures. Other points of interest concern the dimensions of the hydrocarbon chain and the packing of the chloride ions with the terminal groups, ionic ammonium and non-ionic methyl.

Crystal Data.

The samples were prepared and all photographs taken in Sweden using  $\text{CuK}\alpha$  radiation ( $\lambda = 1.542 \text{ \AA}$ ). The unit cell dimensions were obtained from rotation, oscillation, and Weissenberg photographs.

n-Dodecylammonium chloride, molecular formula  $\text{C}_{12}\text{H}_{25}\text{NH}_3\text{Cl}$   
 molecular weight 221.5, crystallizes as thin yellowish needles which decompose at about 100°C.<sup>45</sup> The crystals are monoclinic with,

$$a = 5.68, b = 7.16, c = 17.86 \text{ \AA}, \beta = 92^\circ 16'$$

Two molecules per unit cell. Calculated density 1.015 gm.cm.<sup>-3</sup>

Absent reflections;  $(0k0)$  when  $k$  is odd. Space group  $C_2^2 - P2_1$ .  
 The space group  $C_{2h}^2 - P2_1/m$  is also possible, but requires the atoms to lie in special positions, and was for this reason rejected. The above data are in good agreement with those of Clark and Hudgens.<sup>46</sup>

The intensities of the  $(h0l)$  and  $(0kl)$  reflections were estimated using the Robertson multiple film technique.<sup>40</sup> The absorption factor of the Gevaert film used is 1.93.<sup>47</sup> Lorentz and polarisation factors were applied to the observed intensities giving a set of observed structure factors which were later scaled by comparison with the calculated values.

#### Analysis of the Structure.

$(h0l)$  zone. Approximate  $x$ -coordinates were obtained by McCron<sup>48</sup> who used one-dimensional Fourier syntheses on the  $(00l)$  planes. He also obtained an indication of the  $x$ -coordinates. This structure was modified in the light of the bromide structure,<sup>42</sup> and on re-estimating intensities (McCron had used a wrong film factor), a discrepancy of 32% was obtained. An  $F_0$  synthesis of the electron density projected on  $(010)$  was then carried out from which it was seen that the hydrocarbon chain was arranged in a regular fashion. Consequently it was decided to use sub-cell theory in the calculation of the structure factors for these atoms. The sub-cell dimensions, projected on  $(010)$ , were directly determined from this electron density map, and the sub-cell was chosen so that it contained two  $CH_2$  groups related by a centre of symmetry. There were two such regions, one associated with each chain in the unit

cell, and with their centroids related by normal space group symmetry. Further details of the sub-cell are given in another section.

An ( $F_0 - F_c$ ) synthesis was then carried out, over the whole projections, using sub-cell theory to calculate only the carbon atom contributions to the structure factors. The best positions of the sub-cell atoms, including hydrogen, were then found using trial and error systematically. The best position of the  $Cl^-$  ion was found directly from the difference map. Uncertainty exists, however, about the exact position of the nitrogen atom due to the x- and z- coordinates of it and the  $Cl^-$  ion being almost the same. The nitrogen coordinates were chosen therefore from consideration of the  $F_0$  synthesis and the hydrocarbon chain. The hydrogen atoms of the  $NH_3$  group were positioned from a model and the final figure of merit over the observed reflections was 15.8%. The final x- and z- coordinates are given in Table III and the projection of the structure on (010) is shown in Figures II and III. The calculated and observed structure factors are in Table IV.

The atomic scattering curves employed were:

- (a) carbon and nitrogen (as  $N^+$ ). A McWeeney<sup>41</sup> curve for an isotropic carbon atom with a temperature factor of  $\exp. \left[ -3.9 \left( \frac{\sin \theta}{\lambda} \right)^2 \right]$
- (b) chlorine (as  $Cl^-$ ). A James and Brindley curve<sup>49</sup> with a factor of  $\exp. \left[ -3.9 \left( \frac{\sin \theta}{\lambda} \right)^2 \right]$
- (c) hydrogen. A McWeeney<sup>41</sup> hydrogen curve with a factor of  $\exp. \left[ -4.2 \left( \frac{\sin \theta}{\lambda} \right)^2 \right]$

(0 k l) zone. It was found that all the even and odd C - C bonds as projected on (010) were approximately equal. Since the chain can be assumed regular and planar, and the plane of the chain is inclined to the (010) plane, it follows that the long axis of the chain must lie approximately parallel to (010). The angle between (010) and the plane containing the chain may then be calculated from the expected and the projected bond lengths. This method is very similar to that employed in finding the orientation of many polynuclear hydrocarbons using data derived from a single projection. 50

The projection on (100) has no centre of symmetry and the origin can lie anywhere on the b-axis,<sup>13</sup> which means that the y-co-ordinate of any one atom can be arbitrarily chosen, and that the positions of the others are fixed relative to it. As the exact y-separation of the even and odd carbon atoms was not known, even though the chain had been assumed planar, it was decided to fix the chloride ion at y=0. The problem is now reduced to one of finding the best y-co-ordinates for the other atoms.

The dimensions of the sub-cell for this projection were found from the dimensions projected on (010) and space group considerations. There are again two related sub-cell regions, and each sub-cell has a centre of symmetry. The use of this theory meant that it was much easier to compute structure factors for trial structures involving systematic movement of

the whole chain. The z-coordinates used were the same as those determined from the projection on (010) and the best y-coordinates which could be found by making use of the above facts gave an agreement of 21.9% between observed and calculated structure factors.

An  $(F_o - F_c)$  synthesis then showed that all the even numbered carbon atoms required to be slightly moved in the same direction. This was done by trial and error and was found to have very little effect on the agreement between the calculated and observed structure factors. It was found that shifting all these even numbered atoms by  $0.06\text{\AA}$  gave the best agreement, reducing the figure of merit from 21.9% to 21.2%, but that a further shift of  $0.02\text{\AA}$  only increased it to 21.8%. Since this work indicated that the agreement between the calculated and observed values for the structure factors was not particularly sensitive to changes in atomic coordinates, it was decided not to proceed any further with the refinement.

The reasons for the above facts are uncertain, especially as the  $(F_o - F_c)$  synthesis did not yield any large peaks or troughs, however they may have their origin in:

- (a) the number of reflections available not being so great as in the other zone.
- (b) the fact that the projection does not possess a centre of symmetry.
- (c) a possible unfavourable crystal form.

The scattering curves used were the same as in the other zone.

The projected structure is shown schematically in Figure IV and the y-coordinates are in Table III. The calculated and observed structure factors are given in Table IV.

$(F_o - F_e)$  Synthesis on the Projection on (100)<sup>25</sup>

The normal Fourier coefficients used in our computations are:

$$\left. \begin{array}{l} \text{A terms} - F_o \cos \alpha \\ \text{B terms} - F_o \sin \alpha \end{array} \right\} \text{ - where } \alpha \text{ is a calculated value for the phase constant.}$$

In the computation of the difference map,  $F_o$  was replaced by  $(F_o - F_e)$  with appropriate sign, after which normal procedures were used. The value of  $\alpha$  was left unchanged.

Cochrane<sup>25</sup> gives a formula for the shift required by any atom which lies on a slope in the  $(F_o - F_e)$  synthesis. In the projection on (100) the magnitude of the required shift was found, by trial and error, to be 1.5 times greater than the shift which would have been required had the origin been at a centre of symmetry. This is due to phase-angle lag and the value of 1.5 was taken to be the constant required in the expression for standard deviation of coordinates in the  $(Ok\ell)$  zone.<sup>13</sup>

Application of Sub-Cell Theory.

There are two hydrocarbon chains in each unit cell, and the sub-cell was chosen so that it included only one of these chains, which means that there are two sub-cell regions related by the normal space-group symmetry for  $C_2^2 - P2_1$ . The only symmetry element present in the sub-cell is a centre of symmetry so that the sub-cell space group is  $C_1^1 - P\bar{1}$ , although it contains two right

angles. The sub-cell dimensions, which were found as previously discussed, are:

$$A = 5.68, \quad B = 7.16, \quad C = 2.526\text{\AA}, \quad \beta_s = 114^\circ 9.$$

(capitals have been used to differentiate these translations from the main cell ones, similarly  $\beta$  has a subscript s). Two  $\text{CH}_2$  groups per sub-cell. Each sub-cell region contains six sub-cells, the centroid of one region being at,

$$\frac{x}{a} = 0.296, \quad \frac{y}{b} = 0.452, \quad \frac{z}{c} = 0.463,$$

while the centroid of the other region is at the equivalent position for  $C_2^2 - P2_1$ .

The transformation matrix,<sup>13</sup> where  $s_j^k$  are the transformation coefficients, was determined by measuring the sub-cell edges, in fractional coordinates, as vectors relative to the edges of the main cell. This employs the following relationship

where  $x_k$  are the main-cell vectors and  $X_j$  are the sub-cell ones.

$$(x_k) = (s_j^k) (X_j)$$

The numerical values of the transformation matrix are,

$$\begin{vmatrix} s_1^1 & s_2^1 & s_3^1 \\ s_1^2 & s_2^2 & s_3^2 \\ s_1^3 & s_2^3 & s_3^3 \end{vmatrix} = \begin{vmatrix} 1 & 0 & -0.167 \\ 0 & 1 & 0 \\ 0 & 0 & 0.129 \end{vmatrix}$$

The main-cell structure factors were calculated, as set out in the introduction, from the relationship,

$$F_h = F_h^o + \prod_j \frac{\sin \pi N_j H_j}{\sin \pi H_j} F_H \cdot \exp(2\pi i \sum_m h_k)$$

This was used in the two projections as follows:

(1)  $(h0\ell)$  zone. The plane group involved is  $p2$ .  $F_h^0$  was calculated normally.  $F_H$  values were calculated normally also after the sub-cell indices had been transformed into equivalent main-cell ones by making use of the relationship,

$$(H_j) = (s_j^k) (h_k),$$

which is discussed in the introduction. The term  $\exp(2\pi i \sum_m h_k)$  was simplified to  $2\cos 2\pi \sum_m h_k$  since the centroid at  $\bar{x}_m$  is related by a centre of symmetry to one at  $\bar{x}_m$ . The value of the interference function was found graphically by plotting it against  $H_3$  values (which corresponds to the normal Miller index  $\ell$ ) since this is the only index which varies. The  $H_3$  value for each plane may be easily found.

(2)  $(0k\ell)$  zone. The plane group here is  $pg$  which has no centre of symmetry.  $F_H$  can be found as before, however, since the sub-cell has a centre of symmetry.  $F_h^0$  and  $\exp(2\pi i \sum_m h_k)$  must be evaluated using the structure factor equations for  $pg$ ,<sup>13</sup> while the evaluation of the interference function is entirely analogous with the method used in the  $(h0\ell)$  zone.

The use of this method substantially shortens the structure factor calculations as compared with the orthodox methods, the only disadvantage being that the chain must be assumed strictly periodic.

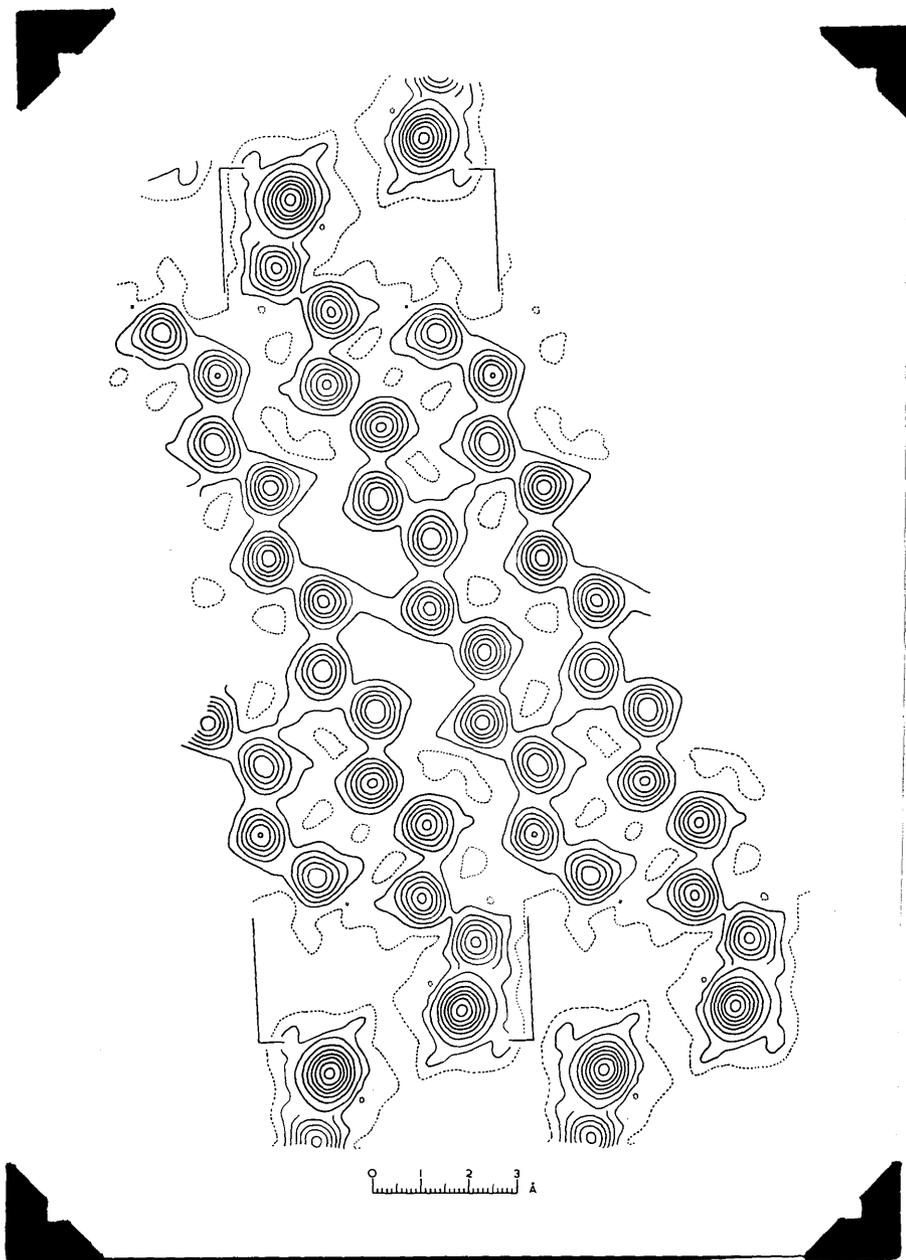


Figure II.

Electron density projected on (010). Interval of contours, one electron per  $\text{\AA}^2$ , except round the chloride ions where only every fourth is plotted. The one electron line is dotted.

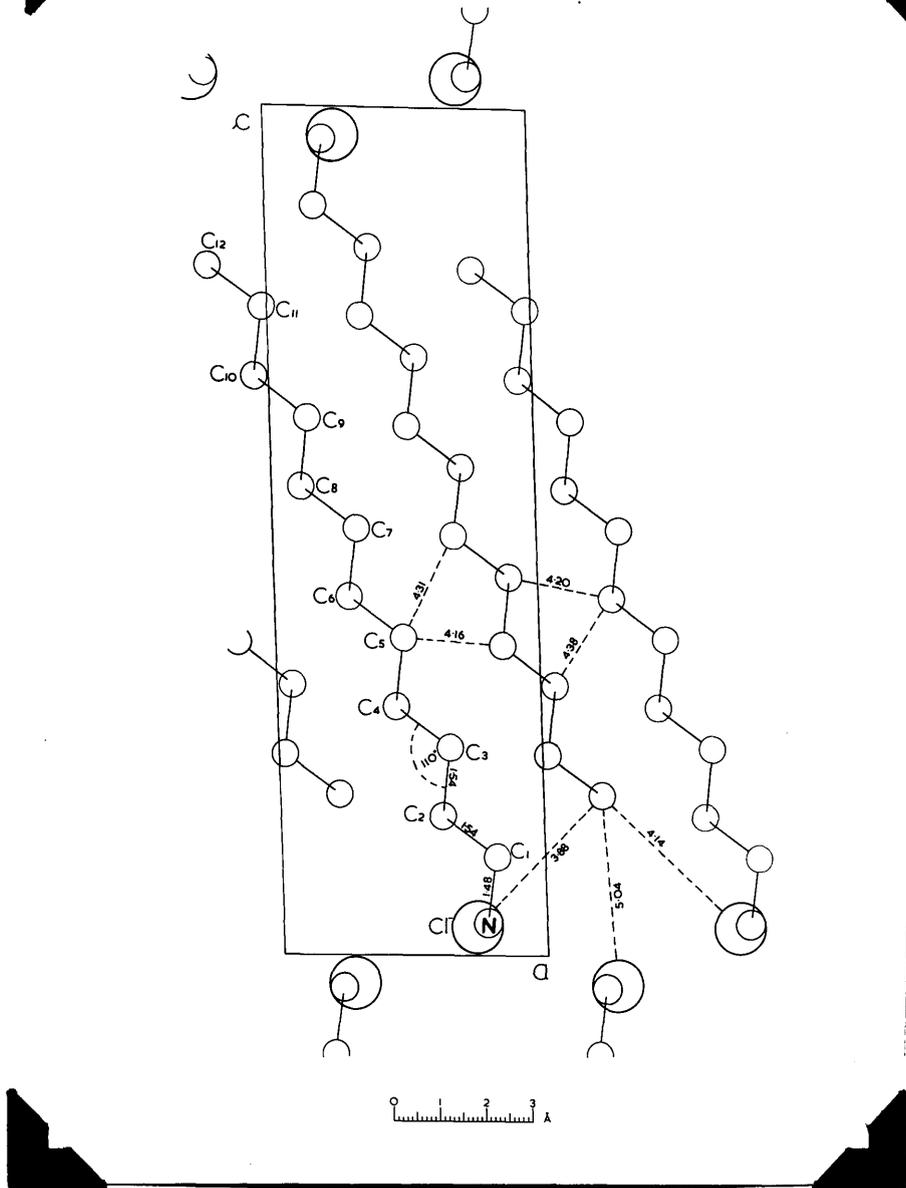


Figure III.

Diagram of structure projected on (010) showing certain interatomic distances. The spatial arrangement is indicated.



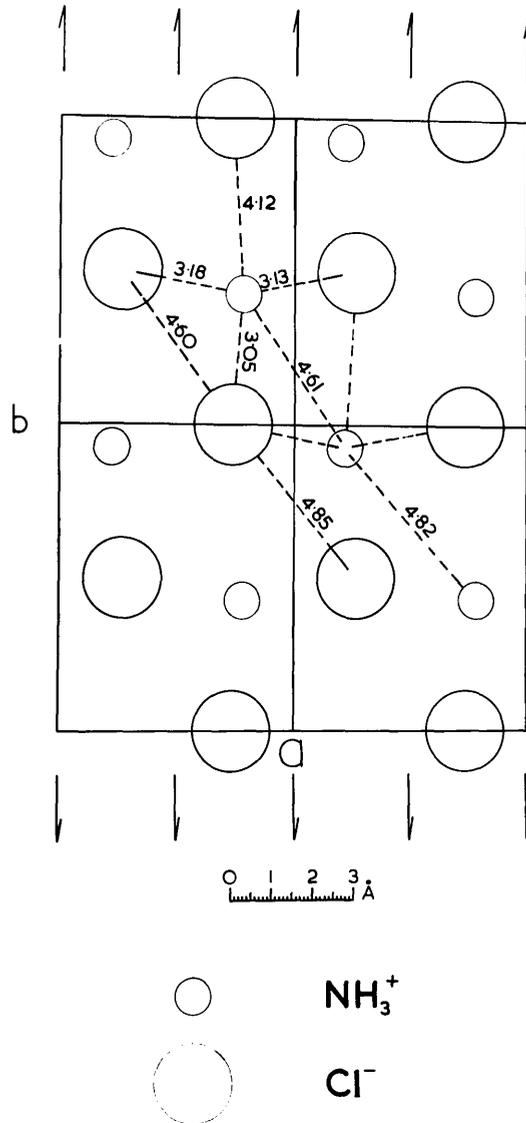


Figure V.

Diagram showing the arrangement of the ionic groups, viewed perpendicular to the  $a$   $b$  crystal face.

Table III. Atomic Coordinates.(1) Sub-Cell Coordinates:

	$\frac{X}{A}$	$\frac{Y}{B}$	$\frac{Z}{C}$
C	0.930	0.965	0.188
H <sub>1</sub>	0.917	0.817	0.172
H <sub>2</sub>	0.735	0.025	0.015

(2) Main-Cell Coordinates: The sub-cell coordinates were converted into main-cell ones, as set out in the introduction, using the relationship,

$$x_k = u_k + t_j A_j + X_j.$$

The coordinates of  $u_k$  are  $\frac{x}{a} = 0.714$ ,  $\frac{y}{b} = 0.452$ ,  $\frac{z}{c} = 0.141$ ;

and the sub-cell coordinates can be converted to equivalent main-cell ones by the use of the transformation matrix. The coordinates of all the atoms, excepting sub-cell hydrogen are then as follows:

	$\frac{x}{a}$	$\frac{y}{b}$	$\frac{z}{c}$
C	0.736	0.000	0.034
H	0.778	0.425	0.038
C <sub>1</sub>	0.815	0.487	0.117
C <sub>2</sub>	0.513	0.417	0.165
C <sub>3</sub>	0.648	0.487	0.246
C <sub>4</sub>	0.446	0.417	0.294
C <sub>5</sub>	0.481	0.487	0.375
C <sub>6</sub>	0.279	0.417	0.423
C <sub>7</sub>	0.314	0.487	0.504
C <sub>8</sub>	0.112	0.417	0.552

	$\frac{x}{a}$	$\frac{y}{b}$	$\frac{z}{c}$
$C_9$	0.147	0.487	0.633
$C_{10}$	-0.055	0.417	0.681
$C_{11}$	-0.020	0.487	0.762
$C_{12}$	-0.222	0.417	0.810
$N$ - $(H_1$	0.761	0.273	0.036
$(H_2$	0.612	0.481	0.015
$(H_3$	0.112	0.520	0.007
$H$	0.193	0.515	0.130

( at end of chain)

These coordinates were referred to orthogonal coordinates (in  $\hat{A}$ ) in order to calculate interatomic distances. These are referred to axes  $a$ ,  $b$ , and  $c$  where  $b$  <sup>DELETE</sup> is chosen normal to  $a$  in the  $a-b$  plane and  $c$  is mutually perpendicular to  $a$  and  $b$ . The relationships are,

$$x' = x + z \cos \beta$$

$$y' = y$$

$$z' = z \sin \beta$$

where  $x'$ ,  $y'$ , and  $z'$  are the orthogonal coordinates in  $\hat{A}$  and  $x$ ,  $y$ , and  $z$  are the unit cell coordinates also in  $\hat{A}$ .

#### Errors in Coordinates.

The observed electron density ( $\rho_0$ ) near the centre of an atom can be closely approximated to by the relation<sup>9</sup>,

$$\rho_0 = Z \left( \frac{p}{\pi} \right)^{\frac{3}{2}} \exp. \left[ -pr^2 \right]$$

where  $Z$  is the atomic number of the atom in question,  $r$  is the distance from its centre in  $\hat{A}$ , and  $p$  is a constant for that

particular atom. We can thus calculate a value for  $p$  from measurement on an electron density map and the value so obtained for carbon from the (010) projection is 3.2

It is also possible to obtain a value for the curvature at the centre of the  $n$ th atom  $\left\{ \left( \frac{\partial^2 e_0}{\partial x^2} \right)_n \right\}$  by the use of this relationship, i.e.

$$\left( \frac{\partial^2 e_0}{\partial x^2} \right)_n = 2 p Z \left( \frac{p}{\pi} \right)^{\frac{3}{2}}$$

We can now utilise the formula for standard deviation of a coordinate which is discussed in the introduction. This gives us for carbon,

$$\sigma(x_n) = \sigma(z_n) = \pm 0.045 \text{ \AA}$$

$$\sigma(y_n) = \pm 0.049 \text{ \AA}$$

all the carbon atoms having the same standard deviation. The position of the nitrogen atom is much less certain due to overlap, however it will not have a smaller standard deviation than the carbon values. For chlorine the values are much smaller,

$$\sigma(x_{cl^-}) = \sigma(y_{cl^-}) = \sigma(z_{cl^-}) = \pm 0.003 \text{ \AA}$$

The standard deviations of bond lengths and angles can now be calculated.<sup>4</sup> We use the root mean square value for the standard deviation of carbon coordinates for this purpose,<sup>51</sup> since the value for the  $y$  coordinate differs from those for the  $x$  - and  $z$  - coordinates.

atoms.	standard deviation of bond length.	angle.
C-C and C-N	$\pm 0.065 \text{ \AA}$	$\pm 4^\circ$
$cl^-$ -C and N- $cl^-$	$\pm 0.049 \text{ \AA}$	$\pm 1^\circ$
$cl^-$ - $cl^-$	$\pm 0.004 \text{ \AA}$	—

This method of calculation neglects the fact that the carbon positions were obtained from a sub-cell, i.e. treated collectively, and it therefore refers to individual distances. The averaged distances, as such, are probably more accurate. Indeed it is possible that the standard deviation for the averaged bond length is less, by a factor of the square root of the number of bonds over which averaging has taken place, than the standard deviation of any one bond. This means that, if we average all six carbon-carbon bonds, the standard deviation of the average value is  $\pm 0.027 \text{ \AA}$ , with  $\pm 0.03 \text{ \AA}$  being acceptable as a maximum value for the average standard deviation.

#### Interatomic Distances and Angles.

(1) Hydrocarbon chain. This must be very regular as shown by a low figure of merit being obtained upon application of sub-cell theory which gives only any variation in alternate bonds and all the angles equal, i.e. no fine details (if any) along the chain are shown up. The actual distances and angles are:

$C_1 - N$	$1.48 \pm 0.065 \text{ \AA}$
$C_n - C_{n+1}$	$1.54 \pm 0.03 \text{ \AA}$
all angles	$110^\circ \pm 2.5^\circ$
$C_n - C_{n+2}$	$2.526 \pm 0.030 \text{ \AA}$

These bond lengths and angles were obtained by the use of sub-cell theory and are averaged over the whole chain. They agree well with the normally accepted values.<sup>31,57.</sup>

The interchain distances of closest approach have also been calculated and are shown in Figure III. The chains are about

4.2Å apart, which separation is in good agreement with the values found in n-hexatriacontane<sup>31</sup> and polythene.<sup>52</sup> The packing is in the manner most economical as regards space, the chains being interlocked, as can be seen upon inspection of Figures III and IV. This type of arrangement is typical of hydrocarbon chains in general and accounts for the fact that there is a certain constancy about their angles of tilt to (001). For our compound the tilt is 66° (deduced from the sub-cell angle  $\beta_2$ ). This value is exactly the same as that for the hydrobromide,<sup>42</sup> though in the opposite direction, and is closely similar to the 62½° tilt in n-hexatriacontane.<sup>31</sup> Values about 70° have also been found for a series of silver salts of fatty acids,<sup>58</sup> although other angles of tilt are possible if the chains slip over one another so that different zig-zags are interlocked. Such angles have been measured.<sup>57, 59, 60.</sup>

The distances between the terminal methyl groups of the hydrocarbon chain and the chloride ions are also shown in Figure III. The closest approach is 3.88± 0.05Å which is in good agreement with the 3.81Å, the sum of their van der Waals radii.<sup>53\*</sup>

(2) Ionic groups. The Cl<sup>-</sup> and NH<sub>3</sub><sup>+</sup> ions form a layer between the rows of hydrocarbon chains. The arrangement can best be seen viewed perpendicular to the ab crystal face as in Figure V. In each vertical row the Cl<sup>-</sup> and NH<sub>3</sub><sup>+</sup> ions have approximately the same x-coordinates and the rows are related by screw axes, so that they are alternately above and below the ab face.

Each nitrogen atom is surrounded by three chloride ions at 3.05, 3.13, and 3.18Å (average 3.12Å) and by one at 4.12Å which

\* See leaf opposite

4.2Å apart, which separation is in good agreement with the values found in n-hexatriacontane<sup>31</sup> and polythene.<sup>52</sup> The packing is in the manner most economical as regards space, the chains being interlocked, as can be seen upon inspection of Figures III and IV. This type of arrangement is typical of hydrocarbon chains in general and accounts for the fact that there is a certain constancy about their angles of tilt to (001). For our compound the tilt is 66° (deduced from the sub-cell angle  $\beta_2$ ). This value is exactly the same as that for the hydrobromide,<sup>42</sup> though in the opposite direction, and is closely similar to the 62½° tilt in n-hexatriacontane.<sup>31</sup> Values about 70° have also been found for a series of silver salts of fatty acids,<sup>58</sup> although other angles of tilt are possible if the chains slip over one another so that different zig-zags are interlocked. Such angles have been measured.<sup>57, 59, 60.</sup>

LINE 18 AFTER

van der Waal radii.

ADD

(This term is used since the contact is between an ionic and a non-ionic grouping)

terminal methyl groups of the are also shown in Figure 5Å which is in good or van der Waals radii.<sup>53\*</sup> form a layer between the

~~the arrangement can best be seen~~  
viewed perpendicular to the ab crystal face as in Figure V. In each vertical row the Cl<sup>-</sup> and NH<sub>3</sub><sup>+</sup> ions have approximately the same z-coordinates and the rows are related by screw axes, so that they are alternately above and below the ab face.

Each nitrogen atom is surrounded by three chloride ions at 3.05, 3.13, and 3.18Å (average 3.12Å) and by one at 4.12Å which

\* See leaf opposite

is a highly significant difference (see introduction section 5). The fourth chloride ion can therefore exert only a very weak attraction on the substituted ammonium ion. However the distance of 3.05 Å is significantly shorter than 3.21 Å, the sum of the ~~van der Waals~~ <sup>ionic</sup> radii,<sup>53</sup> but the other two distances do not differ significantly. This seems to indicate that at least one hydrogen bond is formed between the ammonium group and the chloride ion. Binnie and Robertson found evidence of similar behaviour in hexamethylene diamine dihydrochloride<sup>54</sup> where the interionic distances <sup>are</sup> of 3.01, 3.07, 3.15, and 3.24 Å.<sup>55\*</sup>

~~correspond very well to those found in our compound.~~

#### Discussion.

The fact that we have been able to apply sub-cell theory shows that the hydrocarbon chains, although occupying general positions in the space group, are built up of a series of regular repetitions. This has been called hypersymmetry,<sup>56</sup> and it has been shown that where this is present the figure of merit for a trial structure is likely to be higher and probably will not drop quite so low on refinement. It is possible, therefore, that some of the present disagreement between calculated and observed structure factors is due to this.

The interatomic distances and angles have already been discussed. These are shown not to differ appreciably from the expected values. As regards the general packing of the molecules two arrangements appear possible, the one actually adopted and one with two chains between successive ionic layers, as for example

\* see leaf opposite

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LINE 10 AFTER

3.24 Å

ADD

It will be noticed that the three shorter of these distances correspond very well with the three short interionic contacts found in our compound.

) ammonium group

found evidence of

hydrochloride<sup>54</sup>

3.15, and 3.24 Å.<sup>55\*</sup>

compound.

The fact that we have been able to apply sub-cell theory shows that the hydrocarbon chains, although occupying general positions in the space group, are built up of a series of regular repetitions. This has been called hypersymmetry,<sup>56</sup> and it has been shown that where this is present the figure of merit for a trial structure is likely to be higher and probably will not drop quite so low on refinement. It is possible, therefore, that some of the present disagreement between calculated and observed structure factors is due to this.

The interatomic distances and angles have already been discussed. These are shown not to differ appreciably from the expected values. As regards the general packing of the molecules two arrangements appear possible, the one actually adopted and one with two chains between successive ionic layers, as for example

\* see leaf opposite

in the fatty acids. 57, 59, 60. The favoured arrangement, however, allows more room for the chloride ion and this is probably the reason for its adoption.

TABLE IV.

Observed and calculated Structure factors.

(1) Values of  $F(h0l)$ .

$h\ell$	$2\sin\theta$	$F_o$	$F_c$	$h\ell$	$2\sin\theta$	$F_o$	$F_c$
00	0	-	248.0	1, $\overline{15}$	1.313	5.7	7.8
01	0.086	*	2.8	1, $\overline{14}$	1.230	4.5	-8.7
02	0.172	11.8	11.0	1, $\overline{13}$	1.146	3.4	-3.1
03	0.258	19.3	18.2	1, $\overline{12}$	1.061	7.0	-7.0
04	0.344	22.0	22.0	1, $\overline{11}$	0.980	3.2	-11.7
05	0.430	19.7	20.0	1, $\overline{10}$	0.898	13.3	-15.4
06	0.516	13.4	12.6	1 $\overline{9}$	0.817	16.3	-17.9
07	0.602	2.7	1.8	1 $\overline{8}$	0.738	18.4	-19.1
08	0.688	2.7	6.7	1 $\overline{7}$	0.658	19.3	-17.5
09	0.774	9.6	-7.6	1 $\overline{6}$	0.580	26.8	-27.4
0,10	0.860	10.4	-11.2	1 $\overline{5}$	0.507	29.8	-27.4
0,11	0.946	11.1	-11.8	1 $\overline{4}$	0.437	31.5	-25.6
0,12	1.032	13.0	-11.7	1 $\overline{3}$	0.373	29.5	-30.1
0,13	1.118	10.2	-11.6	1 $\overline{2}$	0.321	27.1	-24.8
0,14	1.204	9.9	-12.6	1 $\overline{1}$	0.284	13.3	-14.8
0,15	1.290	18.6	-18.2	1 0	0.271	< 1.1	1.3
0,16	1.376	2.7	1.3	1 1	0.288	13.5	15.5
0,17	1.462	3.4	-4.9	1 2	0.326	28.2	27.2
0,18	1.548	3.6	-5.0	1 3	0.379	37.9	31.6
0,19	1.634	3.4	-3.6	1 4	0.443	30.7	29.0
				1 5	0.513	24.9	22.7
1, $\overline{22}$	1.907	4.2	7.8	1 6	0.588	15.3	17.8
1, $\overline{21}$	1.823	< 2.0	1.7	1 7	0.667	10.9	10.5
1, $\overline{20}$	1.738	< 2.3	2.2	1 8	0.743	8.4	9.1
1, $\overline{19}$	1.653	3.0	3.1	19	0.828	59.9	59.2
1, $\overline{18}$	1.570	3.7	4.1	1,10	0.908	3.6	4.8
1, $\overline{17}$	1.486	5.7	5.3	1,11	0.990	5.5	6.9
1, $\overline{16}$	1.400	6.5	6.5	1,12	1.073	5.6	6.0

\* not observed due to the construction of the camera.

$h\ell$	$2\sin\theta$	$F_o$	$F_e$	$h\ell$	$2\sin\theta$	$F_o$	$F_e$
1,13	1.153	3.1	4.2	2 6	0.753	8.5	-11.9
1,14	1.240	< 2.6	1.3	2 7	0.818	2.7	- 6.5
1,15	1.322	< 2.7	1.4	2 8	0.883	< 2.0	- 0.2
1,16	1.408	3.1	-2.0	2 9	0.953	3.7	4.4
1,17	1.491	4.6	-6.2	2,10	1.024	3.1	1.6
1,18	1.576	3.6	-3.7	2,11	1.099	12.3	15.5
1,19	1.661	3.0	-3.2	2,12	1.173	11.5	11.9
1,20	1.747	2.5	-2.9	2,13	1.249	8.6	9.6
				2,14	1.331	6.2	7.4
2, <u>18</u>	1.632	2.9	3.6	2,15	1.408	4.7	6.0
2, <u>17</u>	1.552	3.9	5.1	2,16	1.488	3.6	5.2
2, <u>16</u>	1.470	3.7	6.1	2,17	1.568	3.8	4.5
2, <u>15</u>	1.390	4.6	5.9	2,18	1.648	3.0	5.6
2, <u>14</u>	1.311	7.0	6.3				
2, <u>13</u>	1.233	32.1	34.5	3, <u>19</u>	1.815	3.4	-3.6
2, <u>12</u> <sup>s</sup>	1.158	4.0	-0.3	3, <u>18</u>	1.740	2.8	-3.6
2, <u>11</u>	1.082	2.9	2.6	3, <u>17</u>	1.662	3.4	-3.7
2, <u>10</u>	1.009	6.1	4.9	3, <u>16</u>	1.587	3.6	-3.6
2 <u>9</u>	0.939	7.1	6.9	3, <u>15</u>	1.513	< 2.7	-2.8
2 <u>8</u>	0.870	6.7	8.2	3, <u>14</u>	1.442	< 2.8	-1.0
2 <u>7</u>	0.803	4.5	6.8	3, <u>13</u>	1.371	< 2.8	2.3
2 <u>6</u>	0.774	< 1.8	3.5	3, <u>12</u>	1.303	12.3	11.9
2 <u>5</u>	0.688	15.4	-15.5	3, <u>11</u>	1.239	6.2	1.9
2 <u>4</u>	0.639	16.7	-14.6	3, <u>10</u>	1.174	11.6	8.1
2 <u>3</u>	0.598	27.3	-20.3	3 <u>9</u>	1.115	11.6	10.7
2 <u>2</u>	0.568	29.3	-24.8	3 <u>8</u>	1.059	11.6	11.6
2 <u>1</u>	0.549	31.2	-27.6	3 <u>7</u>	1.005	11.5	10.7
2 0	0.542	33.4	-29.4	3 <u>6</u>	0.958	13.0	9.5
2 1	0.551	36.0	-34.0	3 <u>5</u>	0.915	18.3	10.5
2 2	0.572	51.3	-53.7	3 <u>4</u>	0.879	54.1	51.5
2 3	0.607	19.8	19.0	3 <u>3</u>	0.850	6.6	-2.3
2 4	0.648	14.9	-11.7	3 <u>2</u>	0.830	3.4	4.4
2 5	0.698	15.0	-14.6	3 <u>1</u>	0.818	3.4	2.5

$h\ell$	$2\sin\theta$	$F_0$	$F_0$	$h\ell$	$2\sin\theta$	$F_0$	$F_0$
3 0	0.813	<2.0	1.4	4 1	1.089	10.9	10.3
3 1	0.820	3.3	-1.9	4 2	1.100	8.3	8.6
3 2	0.833	8.5	-5.9	4 3	1.119	8.4	8.4
3 3	0.858	11.1	-9.5	4.4	1.143	8.2	7.0
3 4	0.890	12.5	-12.1	4 5	1.173	4.3	10.9
3 5	0.928	13.6	-12.2	4 6	1.209	3.6	3.4
3 6	0.970	11.7	-9.9	4 7	1.250	<2.6	2.5
3 7	1.020	10.3	-10.8	4 8	1.295	<2.7	0.3
3 8	1.073	10.6	-10.6	4 9	1.340	3.0	-2.0
3 9	0.132	11.0	-11.8	4,10	1.397	4.0	-3.5
3,10	1.194	11.7	-14.2	4,11	1.451	4.6	-4.9
3,11	1.259	16.8	-21.4	4,12	1.512	5.0	-4.0
3,12	1.327	12.4	12.5	4,13	1.570	6.2	-8.4
				4,14	1.635	3.7	-3.4
4,18	1.873	3.0	1.5				
4,17	1.804	2.7	-2.8	5,16	1.912	2.6	-4.8
4,16	1.735	3.6	-4.1	5,15	1.852	<1.8	-0.3
4,15	1.670	5.8	-5.1	5,14	1.796	<2.1	-0.1
4,14	1.606	6.0	-5.7	5,13	1.740	<2.3	-0.3
4,13	1.541	5.9	-6.3	5,12	1.689	<2.4	-1.2
4,12	1.482	7.0	-7.4	5,11	1.639	2.6	-2.5
4,11	1.427	12.5	-12.4	5,10	1.591	4.4	-3.9
4,10	1.370	3.8	5.2	5 9	1.548	6.3	-3.9
4 9	1.320	6.6	-4.2	5 8	1.509	7.6	-6.1
4 8	1.272	6.3	-5.4	5 7	1.472	7.1	-6.5
4 7	1.230	3.0	-4.4	5 6	1.441	7.8	-6.7
4 6	1.191	<2.6	-1.0	5 5	1.414	9.0	-6.6
4 5	1.159	4.1	2.1	5 4	1.392	9.5	-5.7
4 4	1.130	6.4	7.2	5 3	1.375	11.1	-9.3
4 3	1.110	12.3	14.7	5 2	1.361	12.6	-15.0
4 2	1.095	11.8	8.1	5 1	1.357	16.4	15.1
4 1	1.085	13.0	11.3	5 0	1.355	3.8	-1.9
4 0	1.084	11.8	11.8	5 1	1.360	<2.8	-2.7

$h\ell$	$2\sin\theta$	$F_o$	$F_o$	$h\ell$	$2\sin\theta$	$F_o$	$F_o$
5 2	1.369	<2.8	-1.2	6 $\bar{4}$	1.657	3.0	-0.8
5 3	1.382	3.4	1.3	6 $\bar{3}$	1.641	4.2	-1.5
5 4	1.402	4.6	3.9	6 $\bar{2}$	1.631	5.7	-4.0
5 5	1.430	4.7	5.6	6 $\bar{1}$	1.628	6.5	-4.8
5 6	1.458	6.4	6.2	6 0	1.626	5.1	-3.1
5 7	1.492	7.0	6.9	6 1	1.630	5.9	-5.1
5 8	1.530	6.1	6.1	6 2	1.639	3.1	-4.7
5 9	1.572	3.4	5.2	6 3	1.650	3.6	-4.6
5,10	1.618	3.1	4.4	6 4	1.669	3.5	-4.8
5,11	1.666	3.0	3.9	6 5	1.690	3.4	-5.4
5,12	1.719	2.3	3.2	6 6	1.712	3.4	-6.4
5,13	1.770	<2.2	2.2	6 7	1.742	3.2	-8.8
5,14	1.829	3.6	-3.1	6 8	1.778	6.1	6.4
5,15	1.886	2.4	3.6				
				7 $\bar{5}$	1.933	2.4	2.6
6, $\bar{12}$	1.910	2.6	3.2	7 $\bar{4}$	1.920	2.5	2.5
6, $\bar{11}$	1.866	2.2	3.3	7 $\bar{3}$	1.908	2.4	2.2
6, $\bar{10}$	1.825	2.3	3.2	7 $\bar{2}$	1.898	2.0	1.9
6 $\bar{9}$	1.789	2.2	3.4	7 $\bar{1}$	1.893	1.8	1.5
6 $\bar{8}$	1.753	2.3	5.7	7 0	1.897	<1.6	-0.9
6 $\bar{7}$	1.722	2.3	1.3	7 1	1.899	3.2	-2.8
6 $\bar{6}$	1.697	<2.4	0	7 2	1.904	2.7	-2.5
6 $\bar{5}$	1.674	<2.5	-0.9				

<sup>s</sup> This plane was left out of the final  $F_o$  synthesis on account of uncertain sign.

(2) Values of  $F$  ( $Ok\ell$ ).

$k\ell$	$2\sin\theta$	$A'$	$B'$	$ F_o $	$ F_o $
11	0.233	-4.5	-6.8	5.5	8.2
1 2	0.280	-7.2	-8.8	6.3	11.4
1 3	0.338	-7.8	-3.8	5.1	8.7
1 4	0.409	-5.8	7.1	9.0	9.2

$k \ell$	$2 \sin \theta$	$A'$	$B'$	$ F_o $	$ F_e $
1 5	0.483	-2.4	20.8	16.0	20.9
1 6	0.561	1.3	32.8	24.3	32.8
1 7	0.642	6.6	47.8	41.8	48.2
1 8	0.727	-6.5	-27.0	20.6	27.8
1 9	0.808	-1.8	11.7	11.5	11.8
1,10	0.891	-1.0	11.2	9.8	11.2
2 0	0.430	111.7	-62.1	108.1	127.7
2 1	0.440	8.2	8.7	10.9	11.9
2 2	0.467	9.9	5.7	13.8	11.4
2 3	0.503	13.1	1.3	16.4	13.1
2 4	0.552	14.8	-2.1	17.2	14.9
2 5	0.611	13.9	-3.3	13.4	14.3
2 6	0.676	9.9	-2.2	12.2	10.9
2 7	0.743	2.1	1.5	18.1	2.6
2 8	0.817	8.9	-5.3	22.2	10.3
2 9	0.889	-4.6	2.3	< 9.3	5.1
2,10	0.969	-8.6	3.2	10.5	9.4
2,11	1.042	-10.3	6.1	12.3	12.0
2,12	1.121	-10.2	5.5	12.9	11.6
2,13	1.204	- 9.3	1.8	< 11.6	9.5
2,14	1.284	- 7.0	1.7	< 11.9	7.2
2,15	1.368	-10.5	2.7	16.0	10.8
3 1	0.656	- 4.8	0.3	< 7.6	4.8
3 2	0.672	- 8.3	1.6	< 7.8	8.5
3 3	0.700	- 8.8	4.9	< 7.9	10.1
3 4	0.734	- 5.7	10.4	10.8	11.8
3 5	0.780	- 1.2	15.7	13.4	15.7
3 6	0.830	4.7	20.9	16.1	21.4
3 7	0.880	14.3	27.2	23.5	30.7

$k l$	$2 \sin \theta$	$A'$	$B'$	$ F_o $	$ F_c $
4 0	0.860	26.9	-29.9	44.5	40.1
4 1	0.868	8.9	7.4	< 9.0	11.6
4 2	0.880	9.7	4.7	< 9.2	10.8
4.3	0.900	8.8	2.1	< 9.4	9.0
4 4	0.929	7.9	-0.6	< 9.6	7.9
4 5	0.965	6.4	-2.0	< 9.8	6.7
4 6	1.008	4.8	-1.5	< 10.3	5.0
4 7	1.054	1.4	1.2	11.2	1.8
4 8	1.106	5.4	-6.0	17.2	8.1
5 1	1.080	-1.9	3.6	< 10.8	4.1
5 2	1.090	-3.4	4.5	< 9.6	5.6
5 3	1.107	-3.9	7.7	< 10.9	8.6
5 4	1.130	-2.7	8.8	< 11.1	9.2
5 5	1.160	0.1	9.4	< 11.2	9.4
5 6	1.193	4.2	10.6	12.1	11.4
5 7	1.233	11.6	10.7	12.4	15.8

Studies on the Crystal and Molecular Structure ofL-Serine Phosphoric Acid.Introduction.

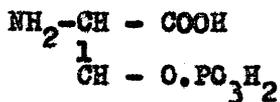
Of recent years a great deal of work has been carried out on the structural analysis of the individual amino acids which go to make up the polypeptide chains of the proteins. The information thus gained has enabled model structures of the proteins themselves to be proposed.<sup>61</sup> However there still remains much work to be done in order to determine fully the nature of all the individual amino acids which are associated with the polypeptide chains. L-Serine phosphoric acid, which occurs naturally in silk fibroin, vitellin (from egg-yolk), and casein, is an important compound to study by X-ray diffraction methods not only for this reason, but also because it contains an orthophosphate group directly attached to an organic residue. So far no X-ray work on this type of linkage has been reported.

Part I. Two-Dimensional Work.Experimental.

Crystals of the compound were obtained from Ågren, de Verdier, and Glomset,<sup>62</sup> and take the form of flattened needles showing extinctions under the polarising microscope in directions parallel to, and perpendicular to, the needle axis. The crystal data were obtained from rotation, oscillation, and Weissenberg photographs taken about all three axes using  $\text{Cu K}\alpha$  radiation ( $\lambda = 1.542 \text{ \AA}$ )

Crystal Data.

L-serine phosphoric acid,  $C_3H_7O_6NP$ , structural formula,



Observed density  $1.707 \text{ g.cm}^{-3}$ , calculated  $1.712 \text{ g.cm}^{-3}$ , m. pt.  $167^\circ \text{ C}$  (with decomposition). The crystals are orthorhombic with

$$a = 7.79 \pm 0.01, b = 10.24 \pm 0.04, c = 9.09 \pm 0.02 \text{ \AA}$$

Absent spectra, (h00) when h is odd, (0k0) when k is odd, (00l) when l is odd. Space group  $D_2^4 - P 2_12_12_1$ . Four molecules per unit cell.

Intensity data for the (hk0), (h0l), and (0kl) reflections were recorded using the Robertson multiple film technique<sup>40</sup>, and were estimated by eye. The normal Lorentz and polarisation factors were applied and structure factors on a relative scale obtained. The film used was Ilford Industrial G which has an absorption factor of 3.22 (see Appendix II).

Attempts at Structure Determination.

The projection on (010). A sharpened Patterson synthesis with the origin removed<sup>12</sup> was carried out and the highest peaks assumed to correspond to phosphorus-phosphorus vectors. After the position of the phosphorus had been determined in this way the origin was shifted from the one given in International Tables,<sup>13</sup> which is not at a centre of symmetry, to one at  $\frac{1}{4}$  of the way along one axis which is at a centre of symmetry in projection. This is because although the three screw axes are not concurrent, in the three main projections any two will appear concurrent.

An  $F_0$  synthesis was then carried out using only planes where the phosphorus contribution was more than one quarter of the possible value. It proved to be impossible to fit a model onto the peaks which appeared, although the phosphorus atom was clearly defined, and another  $F_0$  synthesis, which used signs derived from the peaks which appeared, gave an electron density map worse than the first one.

A new line of attack using Harker-Kasper inequalities and statistical relationships between signs was then tried. The phosphorus position was assumed to be correct and signs were allotted to the four planes to which it made the greatest contribution. Nineteen further signs of strong reflections were then determined and all twenty-three were used in another  $F_0$  synthesis.

Positions for all the atoms were then chosen in regions which appeared positive in both sets of electron density maps, although only the phosphorus atom was positively identified. Structure factors calculated from these gave a discrepancy of 43% and another  $F_0$  synthesis using the new signs gave well defined peaks except in one instance. The map is shown in Figure VI. Unfortunately it was again impossible to fit a model onto the peaks, even the position of the oxygen atoms in the  $PO_4$  group being uncertain. The coordinates of the peaks which were obtained are given in Table V.

At this point it was decided to investigate the a-axis projection in order to see if the two could be correlated. The projection on (100). Here again the origin was shifted in order that it might be at a centre of symmetry.

The first method of analysis tried involved the use of

inequality and statistical relationships as before, two of the  $(00\ell)$  planes being sufficiently strong to be included with known sign (allowing for the change in origin). Only the signs of reflections with  $(k + \ell)$  even could be determined in this way, although the odd ones could be interrelated. An  $F_0$  synthesis was then carried out using only reflections with  $(k + \ell)$  even and thus introducing false symmetry. A large peak appeared for phosphorus at  $y = \frac{1}{2}$ , and an attempt was made, using the map, to determine some of the signs of the odd planes with the aid of structure factor graphs. Another  $F_0$  synthesis was then done which gave reasonably resolved peaks.

A model was then used in an attempt to interrelate the two projections. Only one plausible structure could be found by this method and  $(0k\ell)$  structure factors calculated for it bore no resemblance to the  $F_0$  values so that this line of attack was abandoned.

A sharpened Patterson synthesis with the origin removed<sup>12</sup> was then computed and interpreted by superposition methods.<sup>9</sup> This gave seven atoms only, one of which appeared to be phosphorus, and an  $F_0$  synthesis using 40 out of 119 possible terms was then carried out. This gave a large peak, at the position chosen for the phosphorus atom surrounded by four smaller ones which were taken to be the oxygens of the  $PO_4$  group.

Using signs determined by these five atoms another  $F_0$  synthesis gave peaks onto which it proved possible to fit a molecular model. On calculating structure factors a discrepancy of 40% was obtained but it did not prove possible to improve the agreement

using either  $F_0$  synthesis or trial and error methods.

The best Fourier synthesis obtained is shown in Figure VII and the last set of atomic coordinates obtained are in table V.

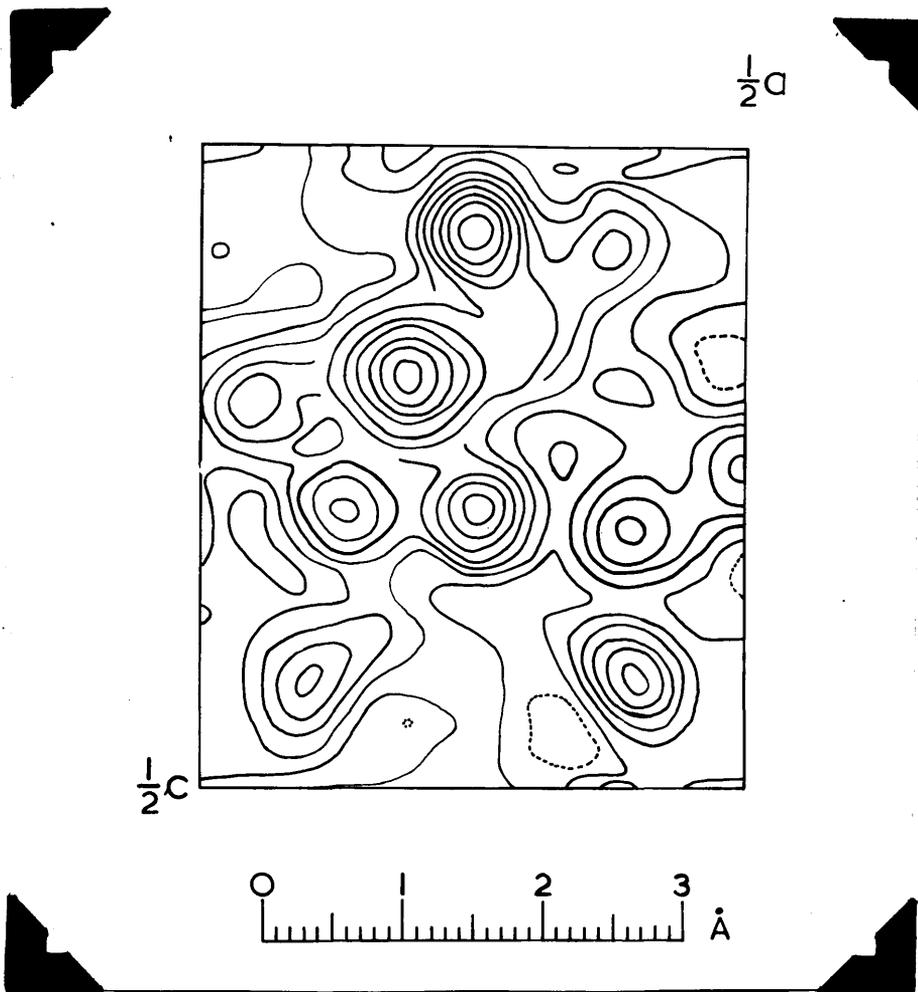
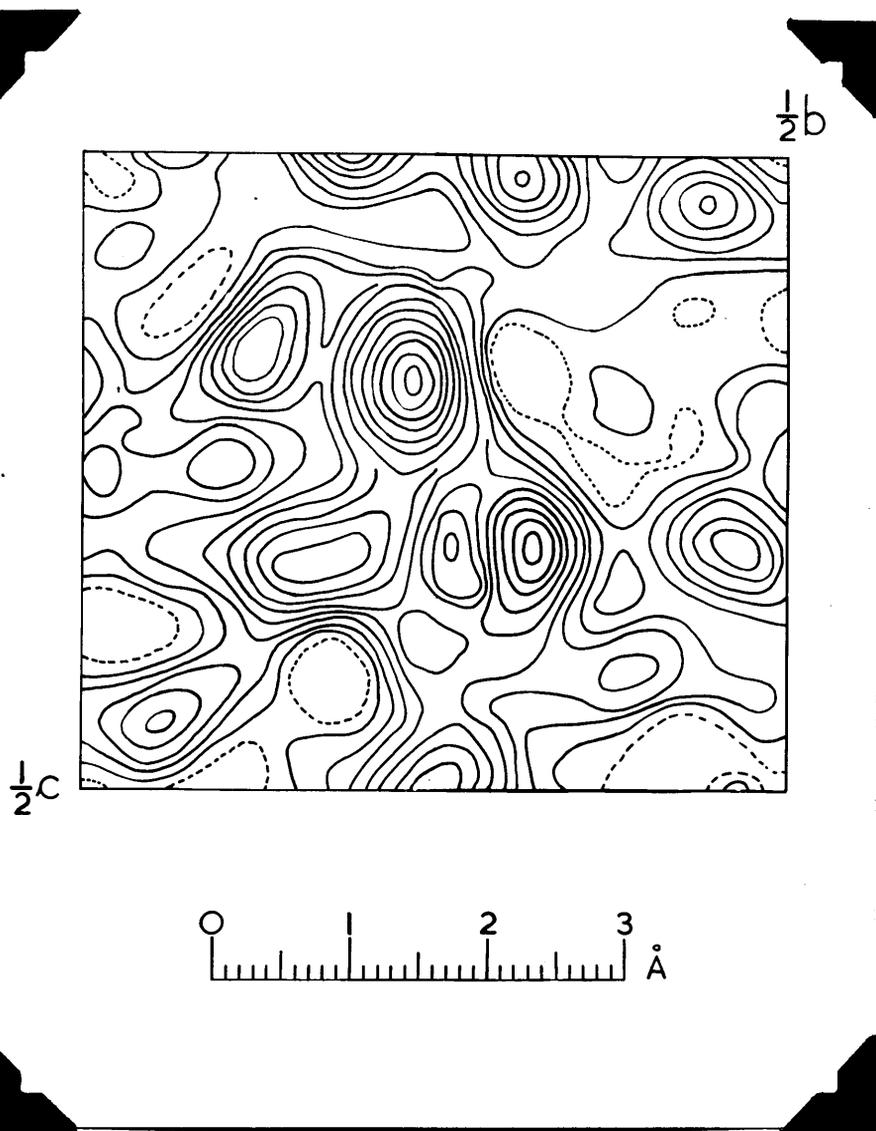


Figure VI.

Fourier map projected on (010). Contour scale arbitrary, the lowest line being dotted. In the region of the phosphorus peak only every second contour is plotted.



**Figure VII.**

Fourier map on (100). Arbitrary contour scale, lowest line dotted, and with only every second line plotted in neighbourhood of the phosphorus peak.

TABLE V.

Atomic Coordinates. Since the structure could not be refined these will, at best, be very approximate if not completely wrong.

The origin is that given in International Tables.<sup>13</sup>

(1) Projection on (010).

Peak	$\frac{x}{a}$	$\frac{z}{c}$
P	0.191	0.179
1	0.105	-0.202
2	0.245	-0.217
3	0.370	-0.215
4	0.102	-0.086
5	0.403	-0.083
6	0.250	0.066
7	0.378	0.079
8	0.049	0.199
9	0.388	0.188

(2) Projection on (100). The atoms are in postulated positions which correspond to peaks on the Fourier map.

	Atom	$\frac{x}{b}$	$\frac{z}{c}$
PO <sub>4</sub> group	P	0.233	0.178
	O <sub>1</sub>	0.191	0.312
	O <sub>2</sub>	0.123	0.155
	O <sub>3</sub>	-0.189	0.483
	O <sub>4</sub>	-0.181	0.191
	O <sub>5</sub>	-0.035	0.191
	O <sub>6</sub>	0.191	0.387
	C <sub>1</sub>	0.150	0.326
	C <sub>2</sub>	0.057	0.447
	C <sub>3</sub>	-0.112	0.094
	H	-0.057	0.463

## Part II. Three-Dimensional Work.

### Introduction.

The results which are given in Part I led us to the belief that it would be exceedingly difficult and laborious to solve the structure of serine phosphoric acid by two-dimensional methods. Consequently it was decided to collect three-dimensional data with a view to getting a three-dimensional Patterson synthesis done, as it appeared to be the most promising method of approach.<sup>17,18.</sup>

### Experimental.

The unit cell data etc. is given in Part I. Three-dimensional intensity data were obtained from equi-inclination Weissenberg photographs taken with unfiltered  $\text{CuK}_\alpha$  radiation about each of the three principal axes for all equi-inclination angles less than  $30^\circ$ . Only three reflections, the (676), (686) and (677), all of which are very high order, were outwith the range of this method. Altogether nineteen different photographs were taken and the intensities were estimated by eye using the multiple film technique,<sup>40</sup> after which Tunell,<sup>63</sup> Lorentz, and polarisation factors were applied in the usual manner. The structure factors derived from the different photographs were all put on the same scale by utilising the fact that separate photographs of layers of the reciprocal lattice at right angles to one another have certain spots in common which enables an average scaling factor for each photograph to be found. The scaling was made approximately absolute by reference to a set of  $F(h0l)$  values calculated during the two dimensional work and where more than one observed value was available an average was taken after

the scaling had been completed.

The film used was Ilford Industrial G which has an absorption factor of 3.22 (see Appendix II) when the X-ray beam is normal to the film. \* ~~However for angles of incidence other than~~ <sup>DELETE</sup> ~~90° this factor G is increased. 64~~ If  $G_0$  is the factor for 90° incidence and  $G'$  that for incidence at an angle  $\omega$  then,

$$\log G' = \sec \omega \log G_0.$$

It can also be shown that in equi-inclination work  $\omega = \mu$  where  $\mu$  is the equi-inclination angle.

The sizes of the crystals used in this work were:

(1) for photographs taken about the a-axis

(0.24 x 0.27) x 0.55 m.m.

(2) about the b-axis

(0.38 x 0.43) x 1.24 m.m.

(3) about the c-axis

(0.64 x 0.58) x 0.30 m.m.

In each case the axis of rotation was the one not included in the brackets.

The values of  $F(hk\ell)$  which were obtained are given in Table VI.

#### The Three-Dimensional Patterson Synthesis.

It was intended that this should be computed on X-RAC, consequently a scheme had to be worked out for presenting the data in a suitable form. This was done as follows:

\* see opposite leaf

the scaling had been completed.

The film used was Ilford Industrial G which has an

factor of 1.00 (see Appendix II) when the X-ray beam is

IN LINE 4 REPLACE THE DELETED PORTION BY,

However when the X-ray beam is not normal to the film this factor  $G$  is increased.<sup>64</sup> If  $G_0$  is the factor when the beam is normal to the film and  $G'$  that when the angle between the normal and the beam is  $\omega$  then,

~~is other than~~  
~~factor for 90°~~  
~~ray~~

$$\omega = \mu \quad \text{where}$$

$\mu$  is the equi-inclination angle.

The sizes of the crystals used in this work were:

(1) for photographs taken about the a-axis

$$(0.24 \times 0.27) \times 0.55 \text{ m.m.}$$

(2) about the b-axis

$$(0.38 \times 0.43) \times 1.24 \text{ m.m.}$$

(3) about the c-axis

$$(0.64 \times 0.58) \times 0.30 \text{ m.m.}$$

In each case the axis of rotation was the one not included in the brackets.

The values of  $F(hk\ell)$  which were obtained are given in Table VI.

#### The Three-Dimensional Patterson Synthesis.

It was intended that this should be computed on X-RAC, consequently a scheme had to be worked out for presenting the data in a suitable form. This was done as follows:

\* see opposite leaf

X-RAC computes two-dimensional functions directly from the series<sup>10</sup>

$$\sum_1^{20} H_1 \sum_{-20}^{20} H_2 A_{H_1 H_2} \cos 2\pi (H_1 x_1 + H_2 x_2) + \sum_1^{20} H_1 \sum_{-20}^{20} H_2 B_{H_1 H_2} \sin 2\pi (H_1 x_1 + H_2 x_2)$$

Three-dimensional functions are computed from the same series in the form of sections at discrete levels of the third coordinate so that the three-dimensional Patterson function must be put in this form.

The Patterson function for orthorhombic crystals is,<sup>13</sup>

$$P(UVW) = \frac{8}{V_c} \sum_0^{\infty} h \sum_0^{\infty} k \sum_0^{\infty} l |F(hkl)|^2 \cos 2\pi h U \cos 2\pi k V \cos 2\pi l W$$

$$= 2 \sum_0^{\infty} h \sum_0^{\infty} l \left\{ \frac{4}{V_c} \sum_0^{\infty} k |F(hkl)|^2 \cos 2\pi k V \right\} \cos 2\pi h U \cos 2\pi l W$$

The part within the curly brackets may be written as  $C_{hyl}$  in shorthand and can be evaluated by summing one-dimensional Fourier series at sufficiently small intervals of  $y$ . In the present instance intervals of  $\frac{y}{60}$  were used and  $C_{hyl}$  was summed from  $\frac{0}{60}y$  to  $\frac{30}{60}y$ . The three dimensional function can now be evaluated in the form of sections at intervals of  $\frac{y}{60}$  since the Patterson function can be written as,

$$P(UVW) = \sum_0^{\infty} h \sum_0^{\infty} l C_{hyl} \cos 2\pi (hU + lW) + \sum_0^{\infty} h \sum_0^{\infty} l C_{hyl} \cos 2\pi (\bar{h}U + lW)$$

which is a form suitable for X-RAC (since the crystal is orthorhombic  $C_{hyl} = C_{\bar{h}yl}$  and only one eighth of the total cell volume need be covered by the summation).

In the actual Patterson synthesis the coefficients employed were the  $F(hkl)^2$  values without modification since it was felt that the presence of the phosphorus atom would lead to undesirable diffraction effects were any of the usual sharpening functions employed.<sup>9</sup> The origin was not removed since this is

unnecessary in three-dimensional work and the terms were all multiplied by a factor of  $\frac{14}{V_c}$ , where  $V_c$  is the volume of the unit cell in  $\text{\AA}^3$ , so that the terms for summation by X-RAC would not exceed  $100^{10}$ .

### Interpretation of the Patterson Synthesis.<sup>9</sup>

The equivalent positions in the space group  $D_2^4 - P2_12_12_1$  are  $(x, y, z)$ ,  $(\frac{1}{2} - x, \bar{y}, \frac{1}{2} + z)$ ,  $(\frac{1}{2} + x, \frac{1}{2} - y, \bar{z})$ ,  $(\bar{x}, \frac{1}{2} + y, \frac{1}{2} - z)$ . There are only three essentially distinct vectors between these positions and they give rise to peaks in the Patterson synthesis at  $(\frac{1}{2} + 2x, 2y, \frac{1}{2})$ ,  $(\frac{1}{2}, \frac{1}{2} + 2y, 2z)$ , and  $(2x, \frac{1}{2}, \frac{1}{2} + 2z)$ , all of which are in the Harker sections. The phosphorus-phosphorus vectors ought to stand out clearly against the vectors between the lighter atoms, consequently an investigation of the Harker sections was made. When the largest peaks were chosen in these sections the following coordinates were obtained for the phosphorus atom;

$$\frac{x}{a} = 0.185, \frac{y}{b} = 0.025, \frac{z}{c} = -0.172.$$

These coordinates differ somewhat from the ones previously given and are likely to be nearer the truth.

At the moment these atomic coordinates for the phosphorus atom are being employed in an attempt to solve the structure by the method of Beavers and Robertson.<sup>65</sup> This method depends on the presence of a heavy atom and they employed it successfully in the case of strychnine hydrobromide.

As an approximation the Patterson synthesis, when a heavy atom (W) is present, may be considered as consisting solely of vectors

of the W - W and W - V type, where  $V_1, V_2, \dots, V_n$  are lighter atoms. For any one atom  $V_R$  in  $P2_12_12_1$  there occur four W -  $V_R$  peaks whose coordinates may be readily calculated from the equivalent positions, the first of them being  $(x_W - x_R, y_W - y_R, z_W - z_R)$ . The position of W (in this case phosphorus) has been found, hence  $V_R$  can be located directly from the Patterson synthesis peaks by four transcriptions of the synthesis, with its origin transferred to the four equivalent positions of the heavy atom. This will give rise to  $16n$  possible atomic positions, of which  $12n$  will be randomly spaced, and the remaining  $4n$  will show fourfold coincidences on  $n$  points in space. (Here  $n$  is the number of light atoms). These fourfold coincidences should allow the positions of the lighter atoms to be recognised. This method is very similar to that employed in the structural analysis of DL-serine.<sup>17</sup>

In Figure VIII the X-RAC photograph showing the section at  $y = \frac{1}{2}$  for the Patterson synthesis is given. The large peak, which was taken to represent a phosphorus-phosphorus vector, can be clearly seen.

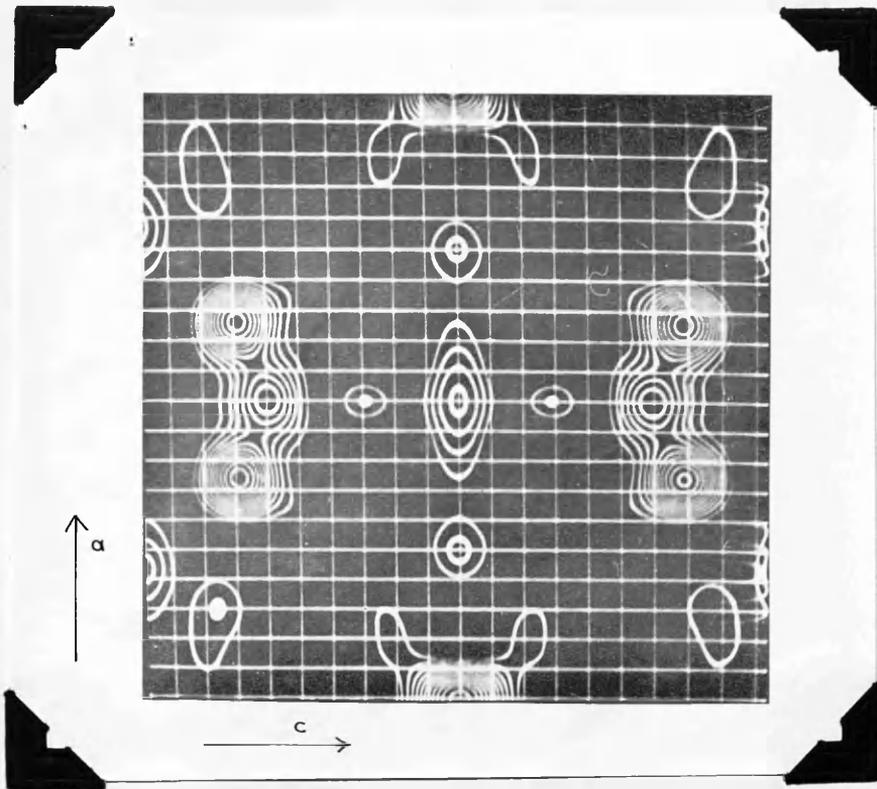


Figure VIII.

Section of the three-dimensional Patterson synthesis at  $y = \frac{1}{2}$ . Whole cell shown. The contours are at arbitrary intervals and go from 0 - 10. The mesh is at intervals of  $\frac{1}{20}$  of the unit cell edge.

Table VI.Observed structure factors, values of  $F(hkl)$ 

$hkl$	$2\sin\theta$	$ F_o $	$hkl$	$2\sin\theta$	$F_o$	$hkl$	$2\sin\theta$	$ F_o $
200	0.398	63.3	012	0.373	15.6	064	1.129	5.0
400	0.796	15.9	022	0.455	24.6	074	1.252	16.3
600	1.194	20.0	032	0.565	17.9	084	1.380	16.4
800	1.592	10.8	042	0.691	9.6	094	1.514	3.2
			052	0.826	23.4	0,10,4	1.649	2.7
002	0.340	67.0	062	0.964	4.3	0,11,4	1,787	2.6
004	0.680	6.8	072	1.105	5.6	0,12,4	1,927	4.4
006	1.020	38.8	082	1.249	7.5			
008	1.360	29.8	092	1.395	8.6	015	0.864	6.8
00,10	1.700	9.5	0,10,2	1.541	1.9	025	0.902	25.1
			0,11,2	1,687	< 1.7	035	0.962	6.3
020	0.301	59.4	0,12,2	1.836	4.0	045	1.041	17.4
040	0.601	34.4				055	1.135	17.9
060	0.902	11.1	013	0.532	25.5	065	1.239	5.2
080	1.202	5.6	023	0.592	14.6	075	1.352	24.1
0,10,0	1.503	9.3	033	0.680	30.7	085	1.462	9.5
0,12,0	1.804	5.9	043	0.789	11.3	095	1.598	5.3
			053	0.909	26.2	0,10,5	1.726	11.2
011	0.228	14.5	063	1.036	4.0			
021	0.346	31.1	073	1.169	39.9	016	1.031	18.8
031	0.482	1.4	083	1.305	6.6	026	1.064	28.4
041	0.625	20.8	093	1.446	22.7	036	1.115	< 2.6
051	0.771	13.0	0,10,3	1.587	6.0	046	1.184	5.2
061	0.918	2.1	0,11,3	1.730	2.6	056	1.267	< 1.5
071	1.066	6.7	0,12,3	1.874	3.7	066	1.361	1.4
081	1.214	20.8				076	1.470	< 3.0
091	1.364	13.4	014	0.696	29.1	086	1.586	< 2.8
0,10,1	1.513	16.7	024	0.744	4.7	096	1.702	< 2.6
0,11,1	1.662	4.9	034	0.815	5.1	0,10,6	1.824	< 2.1
0,12,1	1.812	16.4	044	0.907	14.1	0,11,6	1.856	3.1
0,13,1	1.961	4.2	054	1.014	10.1			

017	1.200	9.1	0,1,11	1.876	1.9	102	0.394	45.1
027	1.227	12.3	0,2,11	1.894	7.9	112	0.422	6.8
037	1.272	9.6	0,3,11	1.924	5.3	122	0.491	27.3
047	1.333	<2.8	0,4,11	1.964	6.3	132	0.598	13.3
057	1.409	4.8				142	0.718	36.1
067	1.493	4.2	110	0.250	16.7	152	0.849	19.0
077	1.594	<2.7	120	0.360	21.3	162	0.984	14.3
087	1.702	8.4	130	0.492	47.9	172	1.123	12.6
097	1.810	2.8	140	0.632	33.4	182	1.266	18.1
0,10,7	1.926	14.1	150	0.780	26.6	192	1.411	5.9
			160	0.927	5.1	1,10,2	1.556	5.0
018	1.369	13.8	170	1.071	5.2	1,11,2	1.702	4.2
028	1.393	18.6	180	1.220	5.6	1,12,2	1.849	11.3
038	1.433	10.2	190	1.370	9.2	1,13,2	1.995	6.3
048	1.487	3.6	1,10,0	1.519	4.9			
058	1.554	3.5	1,11,0	1.668	11.6	103	0.548	8.3
068	1.632	2.0	1,12,0	1.818	4.4	113	0.568	15.7
078	1.724	<2.5				123	0.625	38.9
088	1.824	3.2	101	0.263	12.5	133	0.709	21.1
098	1.926	3.6	111	0.303	9.8	143	0.813	37.9
			121	0.400	23.3	153	0.931	18.4
019	1.538	4.1	131	0.522	23.1	163	1.055	12.9
029	1.559	<2.2	141	0.656	7.6	173	1.186	5.7
039	1.595	<2.7	151	0.797	7.3	183	1.322	10.7
049	1.644	2.5	161	0.939	4.5	193	1.462	1.4
059	1.705	6.2	171	1.085	12.3	1,10,3	1.602	11.7
069	1.776	6.3	181	1.231	4.4	1,11,3	1.744	<1.2
079	1.864	10.7	191	1.381	27.1	1,12,3	1.888	4.3
089	1.956	4.5	1,10,1	1.528	3.1			
			1,11,1	1.677	12.6	104	0.710	4.2
0,1,10	1.707	5.1	1,12,1	1.826	6.3	114	0.726	10.8
0,2,10	1.727	4.7	0,13,1	1.973	4.7	124	0.771	13.8
0,3,10	1.759	8.8				134	0.841	28.7
0,4,10	1.803	1.8				144	0.897	14.1
0,5,10	1.859	5.7				154	1.034	27.3
0,6,10	1.924	1.9				164	1.148	18.8

174	1.267	6.0	107	1.208	8.2	1,0,10	1.711	12.0
184	1.396	9.6	117	1.217	2.5	1,1,10	1.718	2.0
194	1.532	3.7	127	1.245	10.2	1,2,10	1.738	6.8
1,10,4	1.664	1.6	137	1.289	2.9	1,3,10	1.770	3.7
1,11,4	1.801	< 1.4	147	1.349	6.2	1,4,10	1.814	4.8
1,12,4	1.941	5.8	157	1.423	8.1	1,5,10	1.869	4.1
105	0.873	6.8	167	1.507	2.7	1,6,10	1.934	4.6
115	0.886	13.7	177	1.607	11.8			
125	0.924	23.5	187	1.714	3.6	1,0,11	1.883	1.8
135	0.982	16.5	197	1.821	8.0	1,1,11	1.888	2.8
145	1.059	16.8	1,10,7	1.936	2.8	1,2,11	1.907	3.1
155	1.152	7.7	108	1.375	19.2	1,3,11	1.936	2.4
165	1.255	4.6	118	1.383	9.2	1,4,11	1.977	4.1
175	1.366	16.1	128	1.408	11.3	210	0.426	11.8
185	1.486	2.9	138	1.447	9.0	220	0.500	28.9
195	1.612	11.0	148	1.501	9.0	230	0.602	33.5
1,10,5	1.740	8.6	158	1.567	5.8	240	0.721	4.5
1,11,5	1.872	4.0	168	1.644	3.8	250	0.852	9.0
106	1.040	5.0	178	1.735	5.4	260	0.988	10.8
116	1.051	21.7	188	1.835	3.5	270	1.127	17.9
126	1.083	13.5	109	1.542	3.4	280	1.268	5.0
136	1.133	12.6	119	1.550	3.1	290	1.411	10.9
146	1.201	16.4	129	1.571	13.5	2,10,0	1.558	9.3
156	1.284	12.9	139	1.607	7.3	2,11,0	1.701	13.7
166	1.377	4.8	149	1.655	3.3	2,12,0	1.850	6.4
176	1.484	11.0	159	1.716	6.7	201	0.437	27.6
186	1.599	2.1	169	1.786	< 1.2	211	0.463	41.3
196	1.714	2.6	179	1.875	2.8	221	0.531	26.9
1,10,6	1.835	3.9	189	1.966	5.2	231	0.628	23.9
1,11,6	1.867	8.6				241	0.743	22.8

251	0.870	18.2	204	0.790	12.1	296	1.748	5.5
261	1.002	1.5	214	0.804	8.0	2,10,6	1.867	2.2
271	1.140	9.6	224	0.846	9.7			
281	1.279	14.3	234	0.909	13.8	207	1.257	< 1 .7
291	1.421	16.5	244	0.992	12.8	217	1.266	29.4
2,10,1	1.567	10.0	254	1.091	23.4	227	1.293	11.1
2,11,1	1.709	7.5	264	1.199	20.3	237	1.335	17.8
2,12,1	1.882	10.0	274	1.316	28.8	247	1.393	4.2
			284	1.438	21.0	257	1.465	7.4
202	0.526	12.6	294	1.566	6.9	267	1.547	3.6
212	0.548	8 .9	2,10,4	1.700	10.0	277	1.643	3.0
222	0.607	9.2	2,11,4	1.831	4.3	287	1.748	7.1
232	0.693	10.6	2,12,4	1.971	4.6	297	1.853	3.8
242	0.799	13.9				2,10,7	1.967	7 .5
252	0.918	18.8	205	0.940	20.0			
262	1.044	17.5	215	0.952	1.4	208	1.419	9.9
272	1.177	17.3	225	0.987	14.9	218	1.427	17.2
282	1.313	11.7	235	1.043	14.4	228	1.451	12.9
292	1.451	9.2	245	1.116	10.1	238	1.489	3.0
2,10,2	1.594	9.1	255	1.204	13.1	248	1.541	6.1
2,11,2	1.734	8.5	265	1.303	6.1	258	1.606	3.7
2,12,2	1.881	8.0	275	1.411	16.8	268	1.681	3.3
			285	1.526	7.4	278	1.769	8.2
203	0.649	17.0	295	1.647	6.8	288	1.867	< 1.2
213	0.666	24.6	2,10,5	1.774	5.6	298	1.967	5.5
223	0.716	8.0	2,11,5	1.901	4.1			
233	0.790	16.2				209	1.582	6.4
243	0.884	11.2	206	1.098	14.2	219	1.589	5.6
253	0.993	7.9	216	1.109	11.7	229	1.611	10.7
263	1.111	11.9	226	1.139	14.3	239	1.645	1.6
273	1.237	20.6	236	1.187	4.5	249	1.692	6.6
283	1.366	7.9	246	1.252	12.2	259	1.752	1.5
293	1.466	21.1	256	1.331	14.3	269	1.821	1.5
2,10,3	1.669	10.5	266	1.420	5.6	279	1.906	5.9
2,11,3	1.775	5.2	276	1.523	14.3			
2,12,3	1.919	2.5	286	1.635	7.3			

2,0,10	1.748	2.0	391	1.491	16.7	304	0.908	9.7
2,1,10	1.755	4.2	3,10,1	1.629	2.6	314	0.920	7.6
2,2,10	1.772	2.4	3,11,1	1.768	8.6	324	0.957	7.9
2,3,10	1.805	4.6	3,12,1	1.912	5.5	334	1.013	10.4
2,4,10	1.851	5.9				344	1.089	9.4
2,5,10	1.914	5.2	302	0.690	25.5	354	1.179	21.9
2,6,10	1.967	4.7	312	0.706	26.2	364	1.219	12.7
			322	0.753	8.0	374	1.388	8.7
2,0,11	1.913	4.3	332	0.824	6.9	384	1.507	5.2
2,1,11	1.919	2.4	342	0.915	20.9	394	1.629	9.8
2,2,11	1.937	8.0	352	1.021	15.5	3,10,4	1.756	1.5
2,3,11	1.966	3.3	362	1.135	16.7	3,11,4	1.887	9.1
			372	1.258	5.4			
310	0.618	24.2	382	1.387	13.0	305	1.040	< 1.3
320	0.670	32.1	392	1.519	6.9	315	1.051	7.4
330	0.750	3.0	3,10,2	1.655	12.4	325	1.082	14.2
340	0.850	10.6	3,11,2	1.792	10.2	335	1.134	20.6
350	0.961	2.6	3,12,2	1.934	4.1	345	1.201	14.1
360	1.084	18.5				355	1.284	14.0
370	1.211	5.7	303	0.789	16.5	365	1.376	8.4
380	1.345	12.3	313	0.804	12.1	375	1.479	8.0
390	1.481	7.2	323	0.845	6.5	385	1.591	3.3
3,10,0	1.620	17.1	333	0.909	31.5	395	1.707	4.5
3,11,0	1.760	< 1.5	343	0.974	5.9	3,10,5	1.829	3.2
3,12,0	1.904	9.9	353	1.085	22.5	3,11,5	1.954	6.4
			363	1.198	7.5			
301	0.625	23.7	373	1.314	14.6	306	1.182	1.7
311	0.643	18.4	383	1.438	7.9	316	1.192	13.4
321	0.694	20.4	393	1.566	8.5	326	1.220	5.3
331	0.771	24.6	3,10,3	1.698	2.4	336	1.265	4.5
341	0.867	30.6	3,11,3	1.832	2.4	346	1.326	7.0
351	0.978	20.3				356	1.401	1.9
361	1.096	12.7				366	1.487	13.0
371	1.223	6.4				376	1.587	7.5
381	1.356	8.7				386	1.694	9.5

396	1.804	3.4	3,0,10	1.802	4.1	402	0.869	6.4
3,10,6	1.945	5.9	3,1,10	1.808	2.4	412	0.882	18.7
			3,2,10	1.827	3.5	422	0.920	13.6
307	1.332	9.3	3,3,10	1.857	5.7	432	0.979	6.1
317	1.341	7.7	3,4,10	1.857	1.8	442	1.057	15.6
327	1.366	14.5	3,5,10	1.953	5.4	452	1.149	9.9
337	1.406	3.2				462	1.252	27.0
347	1.461	17.3	3,0,11	1.964	1.5	472	1.364	7.9
357	1.531	< 1.5	3,1,11	1.970	3.1	482	1.482	11.6
367	1.608	8.5				492	1.609	5.3
377	1.702	8.1	410	0.810	5.1	4,10,2	1.736	4.7
387	1.804	< 1.4	420	0.852	5.7	4,11,2	1.871	7.4
397	1.906	6.4	430	0.918	8.3			
			440	1.000	14.9	403	0.950	8.0
308	1.488	< 1.8	450	1.098	2.9	413	0.962	11.3
318	1.496	8.5	460	1.205	11.8	423	0.997	11.4
328	1.518	3.2	470	1.321	5.5	433	1.051	13.2
338	1.555	< 1.8	480	1.443	4.0	443	1.124	24.9
348	1.605	6.4	490	1.573	8.4	453	1.212	4.2
358	1.667	2.9	4,10,0	1.702	< 2.3	463	1.310	15.4
368	1.740	8.4	4,11,0	1.840	14.5	473	1.416	5.8
378	1.824	5.6				483	1.536	5.3
388	1.919	6.9	401	0.819	7.6	493	1.653	3.8
			411	0.833	39.6	4,10,3	1.777	5.7
309	1.642	3.1	421	0.873	14.1	4,11,3	1.909	4.6
319	1.649	9.8	431	0.935	27.5			
329	1.669	2.1	441	1.016	19.2	404	1.050	10.5
339	1.703	6.9	451	1.117	8.9	414	1.061	4.8
349	1.748	4.0	461	1.218	5.7	424	1.093	8.3
359	1.806	5.6	471	1.332	4.5	434	1.143	9.2
369	1.873	5.0	481	1.453	10.6	444	1.209	25.3
379	1.957	1.5	491	1.582	3.0	454	1.292	9.2
			4,10,1	1.711	3.9	464	1.384	30.6
			4,11,1	1.848	< 1.5	474	1.485	8.2

484	1.595	9.1	447	1.553	6.7	530	1.095	9.0
494	1.713	4.1	457	1.618	7.0	540	1.167	10.4
4,10,4	1.833	2.8	467	1.692	4.7	550	1.250	3.5
4,11,4	1.951	4.1	477	1.782	3.8	560	1.347	23.2
			487	1.852	4.2	570	1.450	7.0
405	1.168	5.3	497	1.977	1.7	580	1.562	10.8
415	1.178	9.1				590	1.682	9.4
425	1.206	8.0	408	1.579	3.1	5,10,0	1.806	9.4
435	1.252	16.3	418	1.536	< 1.7	5,11,0	1.932	6.5
445	1.313	14.5	428	1.596	8.4			
455	1.389	13.3	438	1.642	5.8	501	1.012	16.5
465	1.476	10.4	448	1.639	3.3	511	1.023	13.2
475	1.570	2.2	458	1.752	4.8	521	1.056	21.3
485	1.674	3.5	468	1.818	4.6	531	1.108	21.5
495	1.787	7.0	478	1.899	4.6	541	1.177	13.5
4,10,5	1.902	2.8	488	1.990	2.2	551	1.261	16.3
						561	1.355	5.3
406	1.295	2.9	409	1.728	5.1	571	1.460	3.6
416	1.300	5.2	419	1.735	6.0	581	1.571	5.6
426	1.330	8.7	429	1.754	10.5	591	1.691	2.0
436	1.407	5.6	439	1.736	5.7	5,10,1	1.814	5.2
446	1.427	4.5	449	1.830	9.0	5,11,1	1.940	4.5
456	1.498	10.1	459	1.885	3.1			
466	1.578	6.3				502	1.055	4.2
476	1.672	6.8	4,0,10	1.820	3.8	512	1.065	4.6
486	1.775	3.9	4,1,10	1.826	< 1.1	522	1.097	3.8
496	1.879	8.4	4,2,10	1.904	1.8	532	1.147	9.8
			4,3,10	1.933	< 1.4	542	1.214	19.6
407	1.432	< 2.0	4,4,10	1.974	5.0	552	1.338	10.4
417	1.440	23.9				562	1.388	15.2
427	1.463	11.5	510	1.010	19.7	572	1.489	5.5
437	1.501	17.0	520	1.041	14.4	582	1.598	11.2

592	1.716	8.2	555	1.512	5.1	509	1.829	6.9
5,10,2	1.838	7.8	565	1.591	5.8	519	1.835	8.5
5,11,2	1.962	9.2	575	1.681	3.0	529	1.854	5.2
503	1.121	7.3	585	1.778	1.5	539	1.884	7.7
513	1.131	23.2	595	1.884	2.2	549	1.925	< 1.2
523	1.161	5.8	506	1.430	9.2	559	1.978	4.5
533	1.208	22.8	516	1.438	9.5	610	1.206	6.2
543	1.272	11.4	526	1.462	11.1	620	1.232	18.9
553	1.350	6.1	536	1.499	6.7	630	1.280	2.8
563	1.439	12.0	546	1.554	7.5	640	1.340	18.9
573	1.537	3.5	556	1.616	2.9	650	1.414	5.4
583	1.643	8.6	566	1.691	13.6	660	1.500	12.4
593	1.757	4.9	576	1.775	< 2.6	670	1.595	< 2.8
5,10,3	1.877	5.1	586	1.872	6.9	680	1.698	6.9
504	1.210	3.4	507	1.555	11.0	690	1.810	2.6
514	1.220	7.9	517	1.562	7.2	6,10,0	1.922	7.1
524	1.247	11.5	527	1.584	13.7	601	1.210	4.1
534	1.291	14.7	537	1.619	10.0	611	1.220	5.9
544	1.351	10.5	547	1.667	11.6	621	1.247	13.9
554	1.425	15.5	557	1.728	2.5	631	1.291	18.3
564	1.509	9.3	567	1.798	2.8	641	1.351	9.5
574	1.601	6.6	577	1.879	1.9	651	1.425	6.0
584	1.703	2.2	587	1.972	4.3	661	1.509	12.1
594	1.814	3.1	508	1.689	4.4	671	1.604	3.8
5,10,4	1.930	3.0	518	1.696	2.9	681	1.707	3.6
505	1.312	< 1.7	528	1.716	2.1	691	1.818	2.4
515	1.321	23.9	538	1.748	3.7	6,10,1	1.930	3.5
525	1.346	11.4	548	1.800	6.4	602	1.246	9.3
535	1.387	16.2	558	1.847	3.9	612	1.255	23.7
545	1.443	15.4	568	1.915	6.1	622	1.282	3.9

632	1.325	2.8	625	1.500	2.1	609	1.943	2.4
642	1.383	12.1	635	1.537	11.4	619	1.949	2.8
652	1.459	3.5	645	1.587	9.6	629	1.966	5.4
662	1.538	14.3	655	1.650	5.5			
672	1.631	4.9	665	1.724	11.3	710	1.402	22.0
682	1.732	3.1	675	1.807	2.6	720	1.430	6.8
692	1.842	3.0	685	1.898	8.1	730	1.468	9.2
6,10,2	1.952	3.0				740	1.520	6.9
			606	1.572	11.2	750	1.588	3.1
693	1.300	11.9	616	1.579	6.9	760	1.662	5.8
613	1.309	13.4	626	1.600	12.5	770	1.750	9.6
623	1.334	4.1	636	1.635	3.1	780	1.842	5.8
633	1.376	20.1	646	1.685	9.2	790	1.948	6.3
643	1.432	13.6	656	1.741	5.5			
653	1.502	10.2	666	1.812	4.9	701	1.408	< 2.0
663	1.582	11.3	676	1.894	*	711	1.416	3.7
673	1.674	10.1	686	1.985	*	721	1.439	3.1
683	1.773	8.2				731	1.478	7.4
693	1.880	7.3	607	1.690	6.2	741	1.531	2.3
			617	1.700	6.2	751	1.608	8.6
604	1.379	6.7	627	1.717	3.4	761	1.672	4.0
614	1.387	19.3	637	1.749	6.6	771	1.758	10.8
624	1.412	3.5	647	1.794	3.6	781	1.850	5.8
634	1.451	5.2	657	1.850	3.9	791	1.956	11.8
644	1.504	15.1	667	1.915	4.7			
654	1.571	4.6	677	1.992	*	702	1.439	11.7
664	1.648	10.6				712	1.447	5.3
674	1.734	3.3	608	1.812	6.5	722	1.470	14.4
684	1.829	3.5	618	1.818	2.4	732	1.499	6.2
694	1.933	1.6	628	1.837	6.3	742	1.559	12.4
			638	1.867	6.1	752	1.624	4.5
605	1.469	10.8	648	1.909	5.0	762	1.698	6.5
615	1.477	9.6				772	1.783	6.2
						782	1.873	3.4
						792	1.977	5.2

703	1.488	2.0	707	1.836	2.5	803	1.676	6.8
713	1.496	1.6	717	1.842	2.3	813	1.683	4.2
723	1.518	<1.4	727	1.860	2.7	823	1.702	2.3
733	1.555	6.7	737	1.891	5.8	833	1.735	7.2
743	1.605	1.8	747	1.932	2.5	843	1.780	1.5
753	1.667	6.1	708	1.950	4.5	853	1.837	9.4
763	1.740	8.4	718	1.956	5.1	863	1.903	3.1
773	1.823	2.6	810	1.600	4.5	873	1.977	9.1
783	1.911	13.6	820	1.620	11.6	804	1.736	3.4
704	1.555	14.3	830	1.655	< 2.6	814	1.743	4.2
714	1.562	7.0	840	1.702	10.2	824	1.762	2.1
724	1.584	14.5	850	1.761	<1.3	834	1.794	7.7
734	1.619	3.4	860	1.831	11.1	844	1.837	2.8
744	1.667	8.6	870	1.910	8.0	854	1.892	8.6
754	1.728	1.5	801	1.603	6.3	864	1.956	2.5
764	1.798	6.0	811	1.610	<1.6	805	1.809	8.2
774	1.877	4.4	821	1.631	6.0	815	1.815	2.1
784	1.963	3.6	831	1.665	4.6	825	1.834	5.1
705	1.637	<1.7	841	1.712	2.5	835	1.864	3.5
715	1.644	6.2	851	1.771	7.4	845	1.934	2.1
725	1.665	7.1	861	1.839	<1.1	855	1.959	3.3
735	1.698	6.7	871	1.918	10.1	806	1.897	9.5
745	1.744	2.5	802	1.630	6.2	816	1.903	2.0
755	1.802	5.4	812	1.637	8.4	826	1.921	6.8
765	1.869	3.0	822	1.657	5.9	836	1.950	1.6
775	1.945	5.2	832	1.691	5.7	910	1.798	11.4
706	1.730	6.2	842	1.737	5.3	920	1.816	3.1
716	1.737	10.1	852	1.795	5.7	930	1.847	<2.0
726	1.756	7.9	862	1.863	6.9	940	1.889	3.8
736	1.786	6.4	872	1.940	8.9	950	1.943	2.9
746	1.831	3.1						
756	1.887	4.6						
766	1.951	0.7						

901	1.800	5.1
911	1.806	6.4
921	1.825	7.4
931	1.855	8.1
941	1.898	3.0
951	1.951	6.4
902	1.827	14.7
912	1.833	5.8
922	1.852	6.7
932	1.882	3.7
942	1.923	4.7
952	1.976	3.6
903	1.865	1.7
913	1.871	4.5
923	1.889	6.4
933	1.919	6.8
904	1.920	14.6
914	1.926	4.5
924	1.943	8.3
905	1.986	3.7
915	1.992	1.3

\* Reflection not recorded.

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Appendix I.Molecular Orbital Description of Dimethylcortierocin.Introduction.

It is now well known that the intimate description of a chemical bond is essentially electronic, but Heisenberg<sup>1</sup> has shown that our knowledge of so minute a particle can never be as precise as a classical treatment requires. Wave mechanics, which is based on Schrödinger's<sup>2</sup> wave equation, provides a solution to this problem. The equation itself can only be completely solved for some rather special cases, so that when dealing with molecules we must use approximations. Of these the two chief ones are the valence bond and molecular orbital (usually abbreviated to m.o.) methods. We shall deal with the latter exclusively.

In the m. o. treatment we consider that each electron moves in an orbital which is essentially polycentric and which is governed by a set of quantum numbers. Molecules can thus be built up by feeding electrons into the allowed orbitals. Most of these electrons can be localised but some, which are called mobile or  $\pi$ , and which occur in conjugated and aromatic systems, have orbitals extending over several atoms.

When we deal with conjugated compounds it is usual to consider only the  $\pi$  electrons and the simplest way to obtain their m. o.'s is by taking linear combinations of the contributing atomic orbitals.<sup>3, 4.</sup> Hence the complete  $\pi$  electron wave-function for a chain containing  $r$  atoms, each contributing one  $\pi$  electron from an atomic orbital  $\psi_r$ , is,

$$\psi = \sum_r c_r \psi_r,$$

where the coefficients  $c_r$  are chosen so as to minimise the energy.

We now introduce an approximate Hamiltonian  $H$  for each  $\pi$  electron which neglects all the other  $\pi$  electrons. The wave equation for one  $\pi$  electron then becomes,

$$H \psi = E \psi$$

The Variation Principle<sup>5</sup> can now be applied to give us  $r$  secular equations, each of which has the form,

$$\sum_s (H_{rs} - S_{rs} E) c_s = 0,$$

where  $s$  also enumerates the atoms in the chain, and where

$$H_{rs} = \int \psi_r H \psi_s d\tau \quad \text{and} \quad S_{rs} = \int \psi_r \psi_s d\tau$$

We can simplify these equations by assuming that;

- (1)  $H_{rs} = 0$  when  $r \neq s$  and  $r$  and  $s$  are not immediate neighbours.
- (2) for the same kinds of atoms all the  $H_{rr}$  and all the  $H_{rs}$  are equal.
- (3)  $S_{rs} = 0$  when  $r \neq s$  and that all the atomic orbitals are separately normalised, i.e.  $S_{rr} = 1$ .

When the secular equations are solved we obtain  $r$  values for the energy  $E$  which represent the  $r$  possible energy levels of a  $\pi$  electron in the molecule. Similarly there are  $r$  sets of values of  $C_r$  which correspond to these energy levels and each set must be normalised so that,

$$\sum_r c_r^2 = 1.$$

The coefficients also have the property of orthogonality, i.e.

$$\sum_r c_{rj} c_{rk} = 0,$$

where  $j$  and  $k$  are different energy levels.

Using these  $C_r$  values we can obtain important information about the molecule and for this we consider all the occupied orbitals. The familiar interpretation of  $\psi^2$  is electron density, so that we

define the  $\pi$  electron charge on the  $r$ th atom by,

$$q_r = \sum_{i=1}^j m_i (c_r^i)^2$$

where  $m_i$  = number of electrons in orbital  $i$  and there are  $j$  energy levels. Similar reasoning leads to the mobile bond order being defined as,

$$p_{rs} = \sum_{i=1}^j m_i c_r^i c_s^i$$

The total bond order is  $1 + p_{rs}$  since there is always a  $\sigma$  bond which is supposed to be of order 1.

Bond lengths can be implied from bond orders if we use an empirical curve for correlation with certain standard compounds. The curve used in this work for carbon-carbon bonds is due to Coulson<sup>6</sup> and is given in Figure AI. The accuracy is about  $\pm 0.03\text{\AA}$ .<sup>7</sup> No equally satisfactory curve exists for oxygen-carbon bonds, but one can be constructed<sup>8</sup> and is also given in Figure AI. These bond orders can also be used to deduce the free valence ( $N_r$ ) of the  $r$ th atom from,

$$N_r = N_{\max} - \sum_s p_{rs}$$

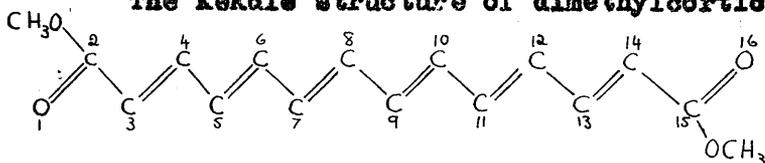
For trigonally hybridised carbon  $N_{\max}$  is given by the square root of the number of similarly hybridised carbons to which it is attached.

Lastly it can be said that from calculations made along these general lines the beginnings of a theory of chemical reactivity are appearing.<sup>10,11,12,13.</sup>

Calculations.

This work is based on the theories set out by Coulson in his book "Valence",<sup>14</sup> and by B. and A. Pullman<sup>15</sup> in "Les Théories Electronique de la Chimie Organique". The mathematical treatment of the determinants was developed by Goodwin and Vand.<sup>16</sup>

The Kekulé structure of dimethylcorticoecin is:



In the calculations the lone pair of electrons on the methoxy oxygen has been neglected since, from chemical properties, they are less mobile than the other  $\pi$  electrons.

It is clear that we may now deal with the molecule according to the methods set out in the introduction. For carbon we put all the  $H_{rr} = E_0$  and all the  $H_{rs} = \beta$ , but since oxygen is more electronegative it will claim a greater share of the  $\pi$  electrons, so that we put  $H_{11} = H_{16,16} = E_0 + \beta$  and  $H_{12} = H_{15,16} = 2\beta$ .

We now have for the sixteen secular equations,

$$\begin{array}{rcl}
 c_1(\overline{E_0 + \beta} - E) + 2c_2\beta & & = 0 \\
 2c_1\beta + c_2(E_0 - E) + c_3\beta & & = 0 \\
 c_2\beta + c_3(\overline{E_0 - E}) + c_4\beta & & = 0 \\
 & \dots & \\
 2c_{15}\beta + c_{16}(\overline{E_0 + \beta} - E) & & = 0
 \end{array}
 \left. \vphantom{\begin{array}{rcl} c_1(\overline{E_0 + \beta} - E) + 2c_2\beta \\ 2c_1\beta + c_2(E_0 - E) + c_3\beta \\ c_2\beta + c_3(\overline{E_0 - E}) + c_4\beta \\ & \dots & \\ 2c_{15}\beta + c_{16}(\overline{E_0 + \beta} - E) \end{array}} \right\} \text{--- I}$$

If the  $c$ 's are to be non-trivial the determinant of their coefficients must be zero. This gives us,





unusual for the roots to exceed  $\pm 3$  when this type of compound is being dealt with. All the roots were refined using Taylor's theorem. Their values are given in Table AI.

Table AI.

Number of Root.	Value (= x).	Factor used.	Number of Root.	Value (= x).	Factor used.
16	2.417	A	8	-0.293	A
15	2.060	S	7	-0.727	S
14	1.884	A	6	-1.129	A
13	1.662	S	5	-1.480	S
12	1.366	A	4	-1.755	A
11	1.000	S	3	-1.935	S
10	0.593	A	2	-2.737	A
9	0.158	S	1	-2.737	S

These values, when we apply the normalisation and orthogonality conditions, lead in turn to the values of  $C_x$  which are in Table A II. The values for only one half the molecule are given, the others being related as stated previously. Only the coefficients for the eight lowest levels have been evaluated since, in the ground state, the rest will be unoccupied.

Table A II.

Root	Fac- tor.	Atoms.							
		1	2	3	4	5	6	7	8
1	S	0.509	0.442	0.192	0.083	0.037	0.015	0.006	0
2	A	0.509	0.442	0.192	0.083	0.037	0.015	0.006	0
3	S	0.065	0.030	-0.070	-0.167	-0.253	-0.322	-0.371	-0.396
4	A	0.112	0.042	-0.150	-0.305	-0.386	-0.372	-0.267	-0.097
5	S	0.142	0.034	-0.234	-0.380	-0.328	-0.106	0.171	0.352
6	A	0.161	0.010	-0.310	-0.361	-0.097	0.251	0.381	0.179
7	S	0.172	-0.024	-0.361	-0.239	0.187	0.375	0.085	-0.313
8	A	0.180	-0.064	-0.379	-0.047	0.365	0.154	-0.320	-0.247

We may also use the roots to calculate the total energy of the  $\pi$  electrons which is  $16E_0 + 25.586\beta$ . The difference between this energy and the energy of the Kekulé structure is  $5.586\beta$ , which represents the resonance energy of the molecule and is comparable with the values for phenanthrene or anthracene.<sup>10</sup>

The bond orders etc. which were calculated using the coefficients of Table A II are given in Table A III. In Figure A II the same values are given in the form of a conventional molecular diagram.

Table A III.

(1) charge distribution and free valence.

	Atom.							
	1	2	3	4	5	6	7	8
charges	1.286	0.798	1.054	0.936	1.000	0.962	0.984	0.972
resultant charges	-0.286	+0.202	-0.054	+0.064	0	+0.038	+0.016	+0.028
free valence	-	-	0.256	0.096	0.158	0.128	0.148	0.142

Table A III.

(2) bond orders and lengths.

	Bond.							
	1-2	2-3	3-4	4-5	5-6	6-7	7-8	8-9
bond order	1.896	1.386	1.790	1.528	1.728	1.558	1.708	1.564
bond length	1.22	1.45	1.37	1.42	1.38	1.41	1.39	1.41

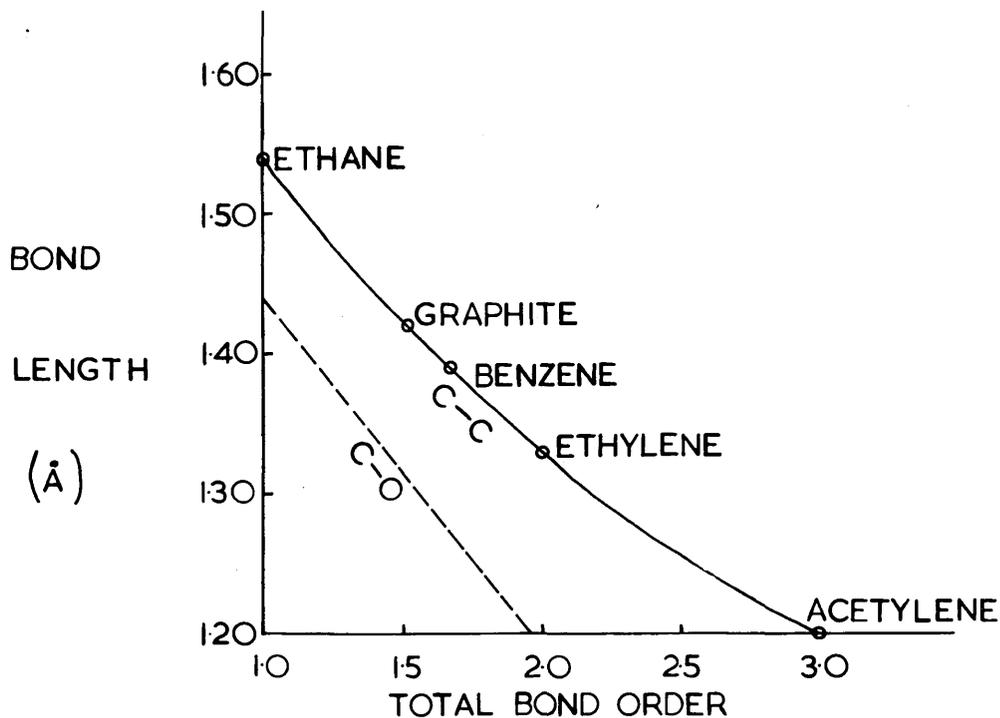
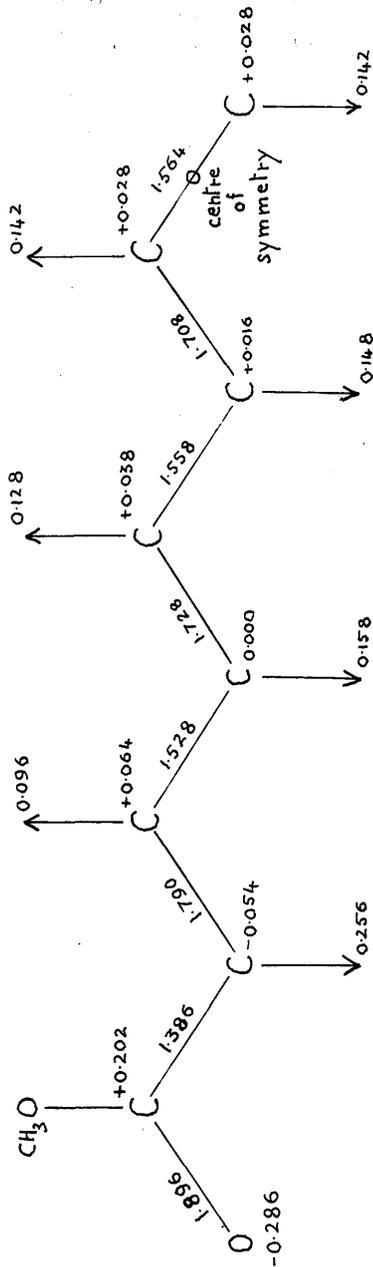
**Figure A I.****Order-length curves used for carbon-carbon and carbon-oxygen bonds.**

Figure A II.  
Molecular Diagram for Dimethylcorticoin.



Discussion.

It is unfortunate that no experimental bond lengths are available for comparison with our calculations. However the bond orders can be compared with those of the corresponding hydrocarbon as Coulson<sup>6</sup> has obtained a general formula for these compounds.\* Our compound shows a greater degree of alternation between the "single" and "double" bonds but the differences<sup>is</sup> are only large for the 2-3 and 8-9 <sup>bond</sup> <sup>it</sup> bonds where they amounts to about  $\frac{1}{3}$  and  $\frac{1}{8}$  of the mobile order. ~~respectively.~~

Predictions of reactivity can also be made which show the carbon (3) to be the most reactive except towards anionic radicals which will favour the central atoms. No reactions have yet been reported for dimethylcortierocin, but this is in line with general chemical considerations.

\* see leaf opposite

IN LINE 5 AFTER

for these compounds.

ADD

Coulson's formula is,

$$p_r = \frac{1}{2n+1} \left\{ \operatorname{cosec} \frac{\pi}{4n+2} + (-1)^{r-1} \operatorname{cosec} \left( 2r+i \right) \frac{\pi}{4n+2} \right\}$$

where  $p_r$  is the mobile bond order of the  $r$ th bond and  $n$  is the number of carbon atoms.

The total bond orders which are calculated by this method are:-

1.851, 1.476, 1.730, 1.564, 1.697, 1.585, 1.683, and 1.595

where the sequence is the same as that given in the molecular diagram.

which will favour the central atoms. No reactions have yet been reported for dimethylcortierocin, but this is in line with general chemical considerations.

\* see leaf opposite

lengths are

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References for Appendix I.

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Appendix II.Re-Determination of the Film Absorption Factor.<sup>1</sup>

Previous workers in this department have found that this factor does not change appreciably for different batches of any one make of film. However it was decided to check up on this by redetermining the factor for  $\text{CuK}_\alpha$  radiation (the factor varying with the wavelength).

Monochromatic radiation was obtained by reflecting  $\text{Ni}$  filtered  $\text{CuK}_\alpha$  from a large calcite crystal, which gives a very intense (200) reflection, and a short exposure oscillation photograph was taken with a pack of four films. A series of wedges of known intensity gradation were also exposed on a separate film and all the films were developed together. The relative intensities of the spots on the different films were then found using a photometer calibrated by means of the wedges.

The film employed was Ilford Industrial G, each sheet of which was kept wrapped in its black paper covering, and which was found to have an absorption factor of 3.22. This value agrees, within the limits of the probably experimental error, with the value of 3.3 which had been employed previously.

This work was carried out in collaboration with G.A. Sim and we are indebted to Dr. S.T.R.S. Mitchell for the use of his photometer.

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