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# A Thosis

### Entitled

The Structure of Yohimbine and other Organic Molecules by X-Ray Crystal Analysis

Submitted for the Degree of

Doctor of Philosophy

in the

University of Glasgov

by

Kathleon Ann Kerr

May 1966.

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#### STRACEY

This thosis describes work carried out ander the supervision of Professor J.M. Tobertson since Scieber 1962. It is divided into two sections.

Soction I describes two structures solved by the heavy-atom method. In both cases, it is the actual structure of the melecule in the chemical sense and the stereochemistry that is of interest, and for this reason refinement was terminated at a fairly early stage.

Part I of this section nevicws certain aspects of heavy-atom and Pattorson methods.

Part II is concerned with the structure of the methiodide of the 7-acctory-7-N-yohimbine. The configuration of the acctory group is the point of major intersect hows. The results of this analysis have been used by N. Finch and W.I. Taylor of the C.I.B.A. Pharmacoutical Company to establish the storeochemistries of a number of rolated compands. This work is also an independent preef of the structure of yohimbine.

Partill is an investigation of the stereochemistration of the monomimum of benail through the coystal structure analysis of the part-branchemeters of the S-curles. Show is some cridence in the literature that crussed the Meisenheimer's enigtual assignment should be covered. but the present work confirme bis sectored.

Bootion II departies on correcte determination of bond longthe in 1.1° binephily? and in perylone. Noth moloculos can be described as compased of two naphthelesic residues joined by long bond(s). The perglane molecule is planar but the binaphthyl molecule is not. Valonce bond theory predicts that the areastic character of the porylone moleculo is localised in the asphthelouic residues and that the Linking bonds are pure single bonds with If this is sto the dimensions of no R-bond charactor. the molecule should be very similar to those of 1,1 binaphthyl whore the dikedral angle between the pleases of the nambthalonic residues is 68°- large enough to proclude any inforaction between the two aromatic systems

The relinements of the two structures are described, and the dimensions of the two solocules are compared with other solecules of similar bond type.

Appendix I is a short note concerning the choice of unit cell in a triclinic lattice.

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# SECTION I

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Two Heavy-Atom Structures

#### INTROLEUCETON

This section of the thesis describes the structural analysis of 7-acetoxy-72-yohimbino and the <u>para</u>bromobenzoate of bensil meconime. It is divided into three parts.

Part one discusses some aspects of the theory of heavyatom and Patterson methods in the context of the problem to which they were applied. The structures in point illustrate different aspects of the cooperative use of the two methods. This is not intended to be a general discussion of the theory or application of either method.

Part two is concerned with the structure of the methiodide of the acetoxy-indolenine of ychimbine. This is an independent proof of the structure and stereochemistry of yohimbine and is in agreement with the previous assignment by chemical means. The feature of interest in this analysis is the stereochemistry at C(7), the junction between the indolenine meiety and the aliphatic ring system. This feature is discussed in terms of the work that has been done on exidative transformations of indole alkaloids at CIBA in Summit, New Jersey.

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Cortain phonome suggest that the accopted configuration of the  $\alpha$  and  $\beta$  isomors should be reversed. The structuro was solved by the use of a "minimum function" derived from a sharpened Patterson synthesis. The configuration of the molecule agrees with that currently accepted in the literature. Since the molecule crystallizes in the triclinic system, the choice of unit cell is not dictated by symmetry considerations. The cell chosen is the Bravais reduced cell. A short discussion of the problem of choosing the unique cell in the triclinio system is included in Appendix 1.

PART I

#### 1.1 The Phase Problem

It is well known that the ideal diffraction pattern can be calculated for a structure containing N atoms. The structure factor expression gives the amplitude and phase of the scattered wave and has the form

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(1)  $\underline{F}(\underline{H}) = \sum_{j=1}^{N} f_j(\underline{H}) \exp (2\pi i \underline{H} \cdot \underline{r}_j) = |F(\underline{H})| \exp (i \operatorname{se}(\underline{H}))$ 

where  $F(\underline{H})$  is the structure factor associated with the plane of Miller index  $\underline{H} = (h, k, \ell); \underline{r}_{i} = (\underline{x}_{i}, \underline{y}_{i}, \underline{z}_{i})$ , the fractional coordinates of the j<sup>on</sup> atom;  $\varphi(\underline{H})$  is the phase angle associated with  $F(\underline{H})$  and  $f_{j}(\underline{H})$  is the scattering factor of the j<sup>th</sup> atom. The intensition of the diffraction maxima in the ideal pattern are obtained from  $F(\underline{H})^{\ell}$ , the square of the scattering amplitude. The positions of the maxima can be calculated from a knowledge of the size and geometry of the unit cell.

The problem facing the crystallographer is the reverse of the one stated above; given the diffraction pattern, calculate the distribution of the atoms in the unit cell. If we regard the electron density function in the unit cell as the sum of N atomic densities, then

(2) 
$$p(\underline{x}) = \sum_{j=1}^{N} p(\underline{x}-\underline{x}_{j}) = \sum_{j=1}^{N} p_{j}(\underline{x}_{j})$$

whore  $\rho(\mathbf{g})$  is the density function;  $\mathbf{g}$  is the vector  $(\mathbf{x},\mathbf{y},\mathbf{z})$ ; and  $\mathbf{g}_{j}^{+} = (\mathbf{g} \cdot \mathbf{r}_{j})$ . The function  $\rho_{j}(\mathbf{z}_{j}^{+})$  is the atomic density function of the j<sup>th</sup> atom defined with the atomic centre  $\mathbf{g}(\mathbf{j})$  as origin. The scattering factor,  $f_{j}(\underline{\mathbf{H}})_{\rho}$  is the Fourier transferred to for the atomic density function of the j<sup>th</sup> atom. Let  $(\mathbf{g} \cdot \mathbf{g}_{j}) = \mathbf{z}_{j}^{+}$ . Then, by definition,  $(3) \quad f_{j}(\underline{\mathbf{H}}) = \int_{-\infty}^{\infty} \rho_{j}(\mathbf{g}_{j}^{+}) \exp((2\mathbf{r} \cdot \mathbf{H} \cdot \mathbf{g}_{j}^{+})) d\mathbf{V} \mathbf{g}_{j}^{+}$ . Again,  $\rho_{j}(\mathbf{g}_{j}^{+})$  is the atomic density function referred to  $\mathbf{g}_{j}$  as origin. By substituting this expression into the generation (1), it can be shown that the structure factor expression is the Fourier transform of the unit cell contents.

(4) 
$$F(\underline{H}) = \int \rho(\underline{r}) \exp(2\pi i \underline{H} \cdot \underline{r}) dV \underline{r}$$
.

By the properties of Fourier series,  $F(\underline{H})$  is exactly equal to the value of the appropriate Fourier coefficients of  $\rho(\underline{x})$ , so that

(5) 
$$\rho(\mathbf{r}) = \begin{pmatrix} \mathbf{l} & \boldsymbol{\Sigma} & \boldsymbol{F}(\underline{H})_{\rho} \exp(-2\pi i \underline{H} \cdot \underline{r}) \\ \mathbf{V} & \mathbf{H} \\ \mathbf{V} & \mathbf{H} \end{pmatrix}$$

The maxima in this function represent the positions of atoms in the unit cell. If the Fourier series could be calculated directly from the experimental data, solution crystal structures would be trivial. However, the observed intensities yield only the structure amplitudes. [F(H)]. The phases are not observed. This is the "phase problem" of exystallography.

-See

Referring back to equation (1), the scattering factors and structure amplitudes can be regarded as known since the former are available in tables and the latter can be obtained from the experiment. This means that the structure factor expression can be formulated in terms of two types of unknowns, the atomic positions or the phases. If either is known, the other can be calculabed. This suggests two possible approaches to the phase "Diroct" methods as developed by Sayre, problem. Zachariasen, Karle and Hauptman and others, attempt to find the phases directly by finding relationships among the structure amplitudes. "Indirect" methods seek to solve the phase problem through the atomic positions.

## 1.2. METHODS INVOLVING HEAVY ATOMS

The heavy-atom method and isomorphous replacement are distinct techniques, but both methods have developed from Robertson's work on the phthalocyanines.

"Free" phthalocyanine and its nickel derivative are isomorphous; they have the same unit coll dimensions and the same distribution of light atoms in the exystal. In the metal derivative, the metal atom is situated at the contre of the molecule. Since the phthalocyamines crystallize in space group P2,/a with two molecules in the unit cell, the molecular centres must lie on special positions. Thus the metal atom lies on a centre of symmetry and makes a positive contribution to all reflections. By comparing the structure amplitudes of "free" phthalocyamine and its mickel isomorph, Robertson (1935,1936) was able to determine the phases of mearly all of the observed reflections from phthalocyamine. This is the carliest example of the solution of a crystal structure without recourse to chemical information.

The first direct application of the heavy-atom method was the solution of the structure of platinum phthalocyanine by Robertson and Voodvard (1940). The platinum compound is not isomorphous with the two phthalocyanines described above, although it has the same space group. A Fourier synthesis using the observed structure amplitudes and the phase of the heavy atom alone revealed the positions of all of the light atoms other than hydrogen.

If the scattering power of one atom is very much larger than that of the other atoms in the structure and if the position of the heavy atom is known, the structure factor expression can be factored into light atom and heavy atom components;

(6) 
$$P(\underline{H}) = \mathcal{I}_{h}(\underline{H}) \exp \left(2\pi i \underline{H}_{\cdot,\underline{J}_{h}}\right) + \sum_{j=1}^{N-1} \mathcal{I}_{j}(\underline{H}) \exp \left(2\pi i \underline{H}_{\cdot,\underline{J}_{j}}\right).$$

, é

With a suitable choice of heavy atom, it will so dominate the scattering that for most planes the phase angle calculated from the first term in the series is a suitable approximation to the true phase. If a Fourier synthesis is calculated with the observed structure amplitudes as coefficients and with phases calculated from the heavy atom position, it will reveal more atomic positions. These can be used in subsequent phasing calculations to give a better approximation to the true phases. The process is continued until the positions of all atoms are known.

Lipson and Cochran (1957) suggest that the method is most effective when the scattering power of the heavy atom is about equal to that of the light atoms If a very heavy atom is used, the in the structure. contribution of the light atoms to the total scattering is so small that experimental errors and diffraction Under these conditions it effects become important. is difficult to obtain accurate parameters for the light atoms. If the heavy atom is too light, the phase of the heavy atom is not a good approximation to the true phase and the Fourier map may reveal no new atomic positions. Sim (1957,1959) has calculated

the proportion of structure factors correctly phased by a beavy atom or a group of heavy atoms in space group PN, and has developed formulae for the distribution of phase angles around the phase of the heavy atom in space group PL. His tables provide a good guide in the selection of a derivative suitable for X-Ray analysis.

#### 1.3. THE PATTERSON FUNCTION

In 1935 Patterson defined the function (7)  $P(\underline{r}') = V_0 \int \rho(\underline{r}) \rho(\underline{r} + \underline{r}') dV_{\underline{r}}$ . Substituting equation 5 in equation 7 we get (8)  $P(\underline{r}') = V_0 \int \rho(\underline{r}) \rho(\underline{r} + \underline{r}') dV_{\underline{r}}$ . (8)  $P(\underline{r}') = V_0 \int \rho(\underline{r}) \rho(\underline{r} + \underline{r}') dV_{\underline{r}}$ .  $X \exp(-2\pi \underline{H}' \cdot \underline{r}') dV_{\underline{r}}$ .

This function is zero except when  $\underline{H} = -\underline{H}^{0}$ ; when this condition applies

(9) 
$$P(\underline{r}^{\circ}) = \bigvee_{\underline{H}}^{\underline{l} \leftrightarrow \infty} F(\underline{H}) F(-\underline{H}) \exp(-2\pi i \underline{H}^{\circ} \underline{r}^{\circ})$$
  
=  $\frac{1}{\sqrt{2}} \sum_{\underline{H}}^{\underline{l} \leftrightarrow \infty} |F(\underline{H})|^2 \exp(2\pi i \underline{H}_{\circ} \underline{r}^{\circ})^{\circ}$ 

Thus a Patterson function is a Fourier series with the square of the structure amplitudes as coefficients. As defined in equation 7, it is the self-convolution of the electron density. The maxima of this function represent interatomic vectors. If there are N atoms in the unit coll, there will be N<sup>2</sup> peaks in the Patterson

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mup with N of these superimposed at the origin and one half of the remaining N(N-1) related to the other half by a centre of symmetry at the origin. The weight of the j<sup>th</sup> atom is its atomic number  $\mathbf{Z}_{i}$  , and the weight of the ij<sup>th</sup> Patterson peak is Z<sub>i</sub>Z<sub>i</sub>. Botk the atomic functions  $\rho(r_{ij})$  and the interatomic function P(r'ij), are Gaussian, but the half-width of a Patterson peak is the root mean square of the half-widths of the atomic functions of the corresponding atoms. Tais leads to chance overlap of vector peaks and tends te complicate the interpretation of the Patterson function. Mowever, inversion centers and translations in the crystal give rise to anti-parallel and parallel vectors respectively. This results in exact superposition of certain maxima, thus increasing their weights relative to background, and is often useful in interpreting a vector map. If there are a small number of heavy atoms in the structure, the outstanding vectors in the Patterson map will be the heavy atom heavy atom vectors and the heavy atom-light atom vectors. Peaks representing interactions between light atoms Norma a fluctuating background of vector density for a structure of moderate complexity. If an atom is heavy enough to be distinguished in the Patterson map, its position in the unit cell can be determined from a knowledge of

 $= \sqrt[6]{a}$ 

# <u> Table 1</u>

Patterson Vectors for the Space Group  $P2_12_12_1$ 

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<u>General Positions</u>			• • • • • • • • • • • • • • • • • • •
1) x, y, z			•
2) $\frac{1}{2}$ -x, -y, $\frac{1}{2}$ +z	`		
3) ż.z., ż.y,z			
4) ~1, żty, ż~2			
Vector Positions			
2-1) \$-2x, -2y, \$		(1-2)	<b>}+2x</b> , 2y, }
3-1) ½, ½-2y, -2z		(1-3)	10 1+2y 28
4-1) -2x, ½, ½-23		(1-4)	2x, 1, 1+22
4-2) ½, ½+2y, -2s	·	(2-3)	-2x, <b>},</b> }+2z
4-3) ½-2x, 2y, ½		(2-4)	1, 1-2y, 2s
$3-2) 2x, \frac{1}{2}, \frac{1}{2}-2z$		(3-4)	<b>}+2x</b> , −2y, ≜
Projection on (010)			
position	W	eight	
$\frac{1}{2}$ - 2x, $\frac{1}{2}$		2	
÷20 ~22		2	
-2x, <del>]</del> -2z	• •	3	
2x, 1-2z	· .	1	

1+2z, 1

23.5 3428

-22. h+23

ł, 2z

the equivalent positions of the space group. For instance, the methiodide of 7-acetoxy-7-H-yohimbine crystallizes in the orthorhombic system with space group P212121. There are four equivalent positions in the space group and four molecules in the unit coll. This should give 16 heavy atom heavy atom interactions with four superimposed at the origin and the other 12 distributed through the cell (Table 1). The symmetry elements present cause eight of these to coincide in pairs when viewed in projection, thus giving four doubleweight and four single-weight peaks. In the unique area of the projection on (010) there are two doubleand one single-weight peak:

 $\frac{1}{2}$ ,2z;  $\frac{1}{2}$ -2x,  $\frac{1}{2}$ ; 2x, $\frac{1}{2}$ -2z; where x, and z are the fractional coordinates of the atom in the unit cell.

### 1.4. DECONVOLUTION OF THE PATTERSON FUNCTION

Several methods are described in the literature for deconvoluting the Patterson function to derive the crystal structure. These are reviewed briefly by Lipson and Cochran (1957) and by McLachlau (1957), and very comprehensively by Buerger (1959). These techniques are most easily understood as attempts to recover the fundamental set from the vector set by superimposing one vector map on the other so that the origin of the

first map lies on a peak in the second. The coincidences are marked and this gives a "reduced vector set". This is illustrated for a fundamental set of <u>points</u> and a vector set of <u>points</u> in figure 1.

If a single weight peak is chosen as the origin for superposition, and if the fundamental set contains a contre of symmetry, that set can be recovered by a single superposition. If the fundamental set is not controsymmetric, the reduced vector set" contains the fundamental set and its mirror image. The penalty incurred for placing the origin of one vector map on an "n"-fold multiple peak in the second is to recover the fundamental set and its congruent image repeated "n" times in different positions. However, successive superpositions...vill reveal a single solution if sufficient single weight peaks are found.

This need for single weight peaks limits the applicability of the method. A Patterson peak representing a vector between like atoms will have a half-width / 2 greater than that of the corresponding atomic peak. If there are N resolved peaks in the Fourier synthesis, the N<sup>2</sup> Patterson peaks are unlikely to be resolved in the same volume. In a structure of reasonable complexity with atoms of similar seattoring power, the single weight peaks will be among the smallest

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in the map and may be lest in a background of vector density. This lack of resolution causes a further difficulty when we attempt to recover the density function If we record only the sites by superposition methods. where the two functions overlap when superposed, chance coincidences will produce false atomic positions. McLachlan suggests that the resolution of a superposition function can be improved if the vector density is summed or multiplied at each point in the cell. This method does enhance the correct peaks but it does not eliminate "ghost" peaks caused by chance coincidences. Buorgor claims that the best resolution is obtained when the minimum value is recorded at each point. The resulting minimum function has a value only when both functions have a positive value and is zero elsewhere.

One of the simplest applications of the method is in the solution of structures with space group PL. The parabromobenzoate of benzil monoxime will serve as an example. In a controsymmetric space group, vectors between atoms not related by the inversion centre are constrained to overlap in pairs to give peaks of weight  $2Z_i Z_j$  at  $\underline{r}_i - \underline{r}_j$ . Atoms related by the centre of symmetry give rise to vectors of weight  $Z_j^2$  at  $2\underline{r}_j$ . Thus the heavy atom-heavy atom interaction produces a single veight peak. Once this has been identified, the electron

\_\_\_\_\_\_

density can be recovered in a single superposition. These conditions apply in the case of the benzil monoxime derivative. Two copies were made of a sharpened, threedimensional Fatterson map. The origin of one map was placed on the centre of the Br-Er vector peak of the second, and the various sections were paired to maintain this displacement. A "minimum function" was derived by tracing minimum contours where peaks in the two sector overlapped. The origin of coordinates in the "minimum function" was defined at the mid-point of the displacement; half-way along the bromine-bromine vector.

### 1.5. SUMMARY

The Patterson function is a very powerful tool in the early stages of structure analysis. It can be used with both centered and non-centered structures, for molecules with atoms of similar scattering power or for molecules with one or more heavy atoms. When combined with the heavy-atom method, it provides a very elegant means of circumventing the phase problem.

## <u>Part II</u>

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The Structure of Yohimbin e:

The Crystal Structure of the Methiodide of 7-Acetoxy-7-H-Yohimbine

### 2.1. INTEDDUCTION

Nohimbine is an indele alkeloid isolated from yohimbehe bark and used as an aphrodisiac in voterinary medicine. The structure of yohimbine, I, was first established by chemical means. The final link in the long chain of evidence leading to the complete structure was provided by Swan (1950), when he determined the position of the hydroxyl group. The work of Witkop (1949) and others established the storeochemistry of the molecule. This work is well documented in "The Alkaloids" (Manakevols. II and VII).



I

The molecule is characterized by five asymmetric control. Hydrogen atoms at C(3), C(15) and C(20) are in axial positions; those at C(3) and C(15) are in the  $\alpha$ configuration. Rings D and E are <u>trans</u> - fused.

The present structure analysis of the methiod.de of 7-acctory-7-N-yohimbine confirms both the gross structure and the stereochemistry of yohimbine itself. However, its importance thes in its connection with research on oxidative transformations of indole alkaleids by N. Finch and V.I. Taylor of the GIBA Pharmaceutical Company. By correlating the results of this X-ray analysis with their own optical rotatory dispersion data. Finch and Taylor have been able to draw some fairly definite conclusions about the mechanisms of certain reactions of indole alkaloids, and to establish the stereochemistry of several alkaloids of this class.

# 2.2. OXIDATIVE TRANSFORMATIONS OF YOHIMBINE AND RELATED ALKALOIDS

In 1962, Finch and Taylor reported that yohimbinoid alkaloids with <u>trans</u>-fused DE ring systems could be converted to their oxindole equivalents by methanolysis of the 7-chloro derivatives (II) to give the imino-ester (III). When the ester is refluxed with aqueous acetic acid, the reaction product is a mixture of oxindoles A and B (IV).



When this series of reactions was attempted using an alkaloid with <u>cis</u> fused DE rings, the results were markedly different. Instead of a 50:50 mixture of epimers of the chloro-derivative, there was a marked prependerance of the Cl-axial epimer. This resulted in a reduction from 40% to 4% in the yield of the methanolysis product, the imino-ester. Thus the method is of practical utility in synthesising only those exindels alkaloids that have <u>trans</u> DE storeochemistry. (Finch and Taylor, 1962a; 1962b).

A later paper (Finch et al., 1963) explored the reactions of Pb<sup>IV</sup> salts of carboxylic acids with alkaloids of the yohimbine type. In all cases, the reactions of the parent alkaloid in methylene chloride with one polar equivalent of the lead salt yielded the corresponding acylexyindolenine (V). When treated with acetic acid, the indolenines eliminated the acylexy residue to give the dehydro-compound (VI). However, if the indolenines are refluxed with methanol and a few drops of acetic acid, these with <u>trans</u>-fused DE ring systems do not react but these with <u>cis</u>-fused systems rearrange to give the corresponding oxindole (IV).

en Station of







parent alkaloid



dehydro compound





Okindole B is the major product of the rearrangement regardless of the stereochemistry at C(3). The reaction mechanism proposed depends on the protonation of the indolenine nitrogen and rearrangement of the resulting carbonium ion with neighbouring-group participation by This mechanism would explain why the acctory molety. grans-fused systems do not rearrange under the conditions of the reaction. The bridge-head nitrogen is more basic than the indolening nitrogen, therefore protonation cannot However, in acylony indolenines occur at the latter site. with cis DE storeochemistry there is a fairly large 1:3disaial interaction with the leas pair of the bridge-This decreases the basicity of this konā nitrogan. attrogon relative to the indelening nitrogen and allows prevenenties of the labour sile. (Finch of ol., 1953).

At this stage, it was uncertain whether the Pb(OAc)<sub>3</sub> group is displaced with retention or inversion of configuration in the formation of the indolenine. On the basis of chemical evidence, it was not possible to determine whether the rearrangement at C(3) occurs before or after the formation of the oxindole. It was hoped that a detailed knowledge of the stereochemistry of at least one of the acyloxy indolenines would be helpful in interpreting these reactions. The compound chosen for X-ray analysis was the acetoxy-indolenine of yohimbine.

### 2.2.1. X-RAY ANALYSIS OF 7-ACETOXY-7-H-YOHIMBINE

The methiodide of 7-acetoxy-7-H-yohimbine was prepared by Dr. Taylor and his colleagues by dissolving 7-acetoxy-7-H-yohimbine in methyl iodide and recrystallising the resulting salt from methanol. This gave colourless, well-formed needles, suitable for X-ray crystallography, m.p. 225-227°C. The crystal structure of the methiodide was determined by the heavy atom method. Approximate locations of all light atoms, other than hydrogen, were found from a series of three-dimensional electron - Pigure 2

-23-

Storeochemistry of the Mothiodide of 7-Acetery 7-N-Yohimbine.



density syntheses. Refinement of positional and thermal of parameters by the mothed, least squares gave a final R-value of 14.60% over 1059 independent reflections.

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Molecular dimensions are given in sub-section 2.4 after a description of the experimental work. The storeochemistry of the mothiodide of 7-acetoxy-7-Uyohimbine is shown in Figure 2. The molecule has a <u>trans</u>-fused DE ring system. The hydrogen at C(3) is in the  $\alpha$  position as is the acetoxy group at C(7). This corresponds to the structural formula VII for the acetomy indolenine tentatively proposed by N. Finch at the I.U.P.A.C. Symposium on the Chemistry of NaturalProducts Prague, 1962.



acetoxy indolcaine of yohimbine

#### 2.2.2. CORNELATION OF X-PAX RESULTS WITH O.R.D. DATA

Optical rotatory dispersion curves are available for 7-acotoxy-7-H-yohimbine, VII; two epimeric 7-chloro-7-H-yohimbans, VIII; methyl-7-hydroxy-7-H-reserpate, IX; and the acetoxy indolenines of pseudoyohimbine, X, and aracine, XI. The formula of these compounds are shown in Figure 3; their O.R.D. curves are shown in Figures 4 and 5. These curves have been drawn from diagrams published by Minch et al. (1965).

The optical rotatory dispersion curves of the two chloro-yohimbans bear a mirror image relationship to one another (Figure 4). This indicates that the C(7)substituent makes the largest contribution to the Cotton The two curves have nearly identical amplitude offect. but opposite sign because the CD ring fusion constrains the C(7)-C(8) bond to be equatorial or pseudo-equatorial in the two epimers. This means that the substituent at C(7) is either axial or pseudo-axial and will have a large interaction with the chromophore group in either the a or the 8 configuration. Any contribution to the dispersion curve from the hydrogen at C(3) would be expected to be small.







X, 7-β-acetoxy-7-Hpseudoyohimbine











XI, 7-a-acetoxy-7-H aricine







-27-


The dispersion curves of the yehinbine and axising derivatives (Figure 5) have the name sign and similar shape in spite of different configurations at the DE ring junction. On the assumption that the C(7)contribution dominates the dispersion curve, the « configuration was assigned to the acetoxy group in the aricine derivative. The curves of the other two compounds have the sign expected of indolenines with a C(7) substituent in the  $\beta$  position. (Finch et al., 1965).

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Those results three some light on the mechanism of the reactions of load salts of carboxylic acids to give acyloxy indolenines. The configuration of the C(7)substituent can be predicted if the position of the hydrogen at C(3) is known. If that hydrogen has an a configuration, the C(7) substituent will de likevise. This observation is compatible with the reaction mechanism proposed in an earlier paper (Finch et 2). (963). Although the storeochemistry of this reaction is not astablished, there is still no information about the subsequent rearrangement of the indolenines to existence. The fact that the observed product is the one expected to be most stable in the reaction mixture would seen to support the thesis that opimerization at C(3) occurs after recreengement.

### 2.3. EXPERIMENTAL

The sample provided by Dr. Taylor consisted of small white needle-like crystals that showed sharp extinction along the needle direction. Rotation, Weissenberg and precession photographs were taken with the crystal mounted so that the needle axis was perpendicular to the X-ray beam. These photographs were used to determine cell dimensions and space group.

-<u>90</u>-

### 2.3.1. CRYSTAL DATA

Formula	C24 <sup>H</sup> 31 <sup>N</sup> 2 <sup>0</sup> 5 <sup>I</sup> F.V. 554-43 ga
System	orthorhombic
Lattice constants	$a = 13 \cdot 16 \text{ \AA}$
	$\mathbf{b} = 26 \cdot 82 \overset{\circ}{\mathbf{A}}$
	$\mathbf{c} = 7 \cdot 33 \mathbf{A}$
Absent Spectra	h00  with  h = 2n + 1
	0k0  with  k = 2n + 1
	$00^{\ell} \text{ with }^{\ell} = 2n \div 1$
Space Group	P2 <sub>1</sub> 2 <sub>1</sub> 2 <sub>1</sub>
Volume	2587 Å <sup>3</sup>
Density (obs.)	1.44 gm./cc
Molecules/cell	<b>4</b>
Density(calc)	1.42 gm./c.c.
F(000)	1128
is ( Culla )	110 cm <sup>-1</sup>

#### 2.3.2. INTENSITY DATA

The OKA zone was photographed by the precession method. Since the crystal-to-film distance is critical in this method (Buerger, 1944), only one film can be placed in the casette at a time. A timed series of six photographs was taken using a stabilised power supply. Exposure times ranged from 11 minutes to 45 hours. Mo Ka radiation was used. The precession angle was set at 30°. Lorentz-polarisation corrections were applied to these data by means of a template. (Euerger 1960, p 185).

Reflection data for the six reciprocal lattice < nots hk0 to hk5 were collected using an equi-inclination Weissenberg camere and multiple film packs (Rebertson, 1943). The intensities of these reflections were estimated visually by comparison with a calibrated intensity scale. The film factor between successive photographs of a series was 3.3 for the zero layer. This was modified for upper layers to allow for the increase in path length through the film with increasing inclination angle (Resemann, 1956). These data were collected with CuKa radiation. They were corrected for Lorentz and polarization factors according to the formula

 $I(hkl) = \frac{1}{\sin 2\theta} + \frac{1 + \cos^2 2}{2} + \frac{1}{|p(hkl)|^2}$ 

The form of the Lorentz factor shown here applies only to zero layer photographs of notating crystals. For upper layers, this must be modified by a factor 1/D, where

$$D = (\cos^2 \mu - \cos^2 \theta)^{\frac{1}{2}} / \sin \theta$$

the rotation factor defined by Tunnell (1939) for the The symbols  $\mu$  and  $\theta$  are the equi-inclination case. equi-inclination angle and the Bragg angle, respectively. No absorption corrections were applied.

The structure amplitudes, |Fobs |, were given the same relative scale by comparison of axial reflections from Neissenberg photographs with common reflections from precession photographs of the Ok # zone.

#### 2.3.3. ANALYSIS OF THE STRUCTURE





The I-I vectors were easily recognized in the Patterson Fractional coordinates of the iodide ion vere; maps.

> y = 0.1026, z = 0.1453.z = 0.1772

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Structure factors were calculated with phases determined by the heavy atom alone. These were placed on absolute scale by comparison of the sums of the observed and calculated structure amplitudes.

 $k\Sigma|F_{o}| = \Sigma|F_{o}|$ 

The R-value, EAAEF, was 37%

The iodide ion is very "heavy" in comparison with the light atoms in the structure. Using Sim's (1957) method, it is possible to predict how well the phase of the heavy atom approximates to the true phase. Sim has calculated the fraction,  $N(\alpha)$ , of structure factors whose phase is within  $\pm \alpha$  of the phase of the heavy atom for various values of x, where x is defined by

$$\mathbf{r} = (\Sigma \mathbf{f}^2_{\mathrm{H}} / \Sigma \mathbf{f}^2_{\mathrm{I}})^{\frac{1}{2}}$$

In general, the probable phase-angle error decreases For this compound, the value of x as y increases. Interpolation on Sim's graph shows that only is 1.46. about 20% of the structure factors will have phase errors greater than ± 35° in a calculation based on the position This estimate is probably a of the heavy atom alone. pessimistic one, as it is strictly valid only for space The space group 22121 has three centrosymmetric group Pl. About 88% of the data from these zones can projections. Since nearly all be expected to have the correct sign.

of the data were considered to be reliably phased by the heavy atom, only these planes with very small values  $|F_{obs}|$  were excluded from the first Fourier calculation. The resulting map revealed the positions of all of the atoms in the skeleton of the molecule. When these were included in the next structure-factor calculation with scattering factors appropriate to carbon atoms, the R-value dropped to 29.6%. After the second Fourier synthesis all 31 light atoms, excluding hydrogen, had been located and identified. They were included in the next cycle of Fourier refinement with scattering ourves corresponding to their chemical types.

After a third cycle, the coordinates were corrected for termination of series errors by a modification of This involves calculating Booth's (1946, 1947) method. two Fourier syntheses with the same coordinates. Observed structure amplitudes are used as coefficients in one map. and calculated structure amplitudes for the other. IC a peak appears at a position, Zp., in the penultimate Fourier map, it should appear in the same position in an Fe map calculated from those coordinates. Termination of series errors cause these two sets of coordinates to differ by Arg = Eg = By for the jth atom. If rog is the position of this atom in the final Fourier map, its carrect position is given by  $y_0 = a_{Z_j}$ .

- 3. <u>.</u> .

If atom peaks appear at different heights in the  $F_0$ and  $F_c$  syntheses, the temperature factor is in error by an amount that can be estimated from the difference in peak height. In the two maps calculated here, peaks in the  $F_c$  synthesis were systematically lower than these in the  $F_c$  synthesis, indicating that the temperature factor  $c = 2^{-B}$ , had been underestimated. When c was increased from 2.8 to 3.0 (B = 4.94 Å<sup>2</sup>), structure factors calculated with corrected coordinates showed an R-value of 21%.

#### 2.3.4. REPINEMENT OF THE STRUCTURE

The computer program used for least-squares refinement on DEUGE was written by J.S. Rollett and is described elsewhere. (Rollett, 1961). The program uses a three by three matrix for the atomic positions and a six by six matrix for the atomic vibrations. In order to provent overflow of the least-squares totals, it was necessary to reduce the scale of the scattering curves and of the structure amplitudes to 1/10 of the normal value. The weighting scheme used gives maximum value to the "average" structure amplitude:

$$\begin{split} \text{if } \left| \begin{array}{c} P_{0} \right| \left\langle \begin{array}{c} P^{*} & \int v &= \begin{array}{c} P_{0} / P^{*} \\ P_{0} \end{array} \right\rangle \\ \left| \begin{array}{c} P_{0} \right| \right\rangle \\ P^{*} & \int v &= \begin{array}{c} P^{*} / P_{0} \\ P^{*} &= \begin{array}{c} 8 \\ P \end{array} \\ \text{vhore} \\ \end{array} \\ \end{split}$$

There is no facility for isotropic refinement in this program. Anisotropic temporature parameters are defined by the equation

exp  $(-B \sin^2 \theta / \lambda^2) = 2^{-(b_{11}h^2 + b_{22}k^2 + b_{33})} + b_{12}hk + b_{23}k^3 + b_{13}h^4).$ 

After each cycle the b<sub>ii</sub> torms from the diagonal of the vibration matrix were used to calculate the thermal parameters parallel to each of the principal axes using the following formulae:-

$$B_{11} = b_{11}/1 \cdot 4427,$$
  
$$B_{X} = 4B_{11}/a^{2},$$

and

(Rollott, 1961; Rossmann, Jacobson, Hirshfeld and Lipscomb. 1959; Cruickshank, 1956).

The components of B<sub>6</sub> in the three axial directions were averaged to give an isotropic temperature factor for each atom. These individual temperature factors were averaged over all members of each chemical type. After two sycles of least squares, the R-value was 19.9%. The average B<sub>6</sub> values for the different chemical types were

 $I = 4.79 \text{ }^{2} C = 4.53 \text{ }^{2} N = 3.74 \text{ }^{2} O = 5.49 \text{ }^{2}$ No further calculations were done on DEUCE. The parameters obtained from the second cycle were used when refinement was continued on KDF9 eight months later.

When the Glasgow SFLS program became available, two rounds of isotropic least squares were calculated with unit weights. The R-value was 17.8%. Cruickshank's (1961) weighting scheme has been recommended for film data and seemed appropriate for this case. It has the form

$$w = 1/(p_1 + |F| + p_2 |F|^2 + p_3 |F|^3).$$

The initial values given to these parameters were as follows:

 $p_1 = 17 \quad p_2 = 8.9 \times 10^{-3} \quad p_3 = 9.9 \times 10^{-5}$ 

After a further three cycles of isotropic block-diagonal least squares, the R-value was 15°4% but the weighting analysis showed that some adjustment in the weighting One more cycle with the constants parameters was required. in the weighting function changed to 8, 1.04 x 10<sup>-2</sup>, and 1.34 x 10<sup>-4</sup>, respectively improved the weighting analysis but had little effect on the R-value. Reflection data for the short and median axes were sorted from the full. scaled, three-dimensional data that had been used for the least-squares refinement. These data were used in a "minimum residual" refinement by the method of Bhuiya and Three cycles of refinement reduced the Stanley (1963). R-value in the short-axis projection from 19°5 to 17°5%. A single refinement cycle in the modian -axis projection produced an improvement of about 1% to give a final R-value 02 17.9%. When the new positions were used for a three-

dimensional structure-factor calculation, the R-value was 16.5% but this dropped to 14.6% after three further cycles of block-diagonal least squares. Refinement was terminated at this point. Unobserved reflections were not included in the analysis.

# 2.4. MOLECULAR DIMENSIONS AND DESCRIPTION OF THE STRUCTURE

Table 2 shows the final atomic coordinates and the isotropic temperature factor for light atoms other than hydrogen. The coordinates and anisotropic thereal parameters of the iedide ion are listed separately at the foot of the table. In all cases, the coordinates are expressed as fractions of the crystallographic exce. Their estimated standard deviations have been coloulated from the least squares residuals and are shown in percentioned to the same number of places of decimal as the coordinates with which they are associated. Both the coordinates and their standard deviations are listed with one figure that is not significant to permit derived quantities to be calculated without less of accuracy.

The isotropic temperature factor used here is the mean square vibration amplitude, U<sub>ine</sub>, where

$$U_{180} = \overline{u}^2 = B/8\pi^2$$

and B is the Dobye factor. Both U and B have the units of A?. Anisotropic vibrational parameters are defined by the equation

$$\exp (-B \sin^2 \theta/\lambda^2) = \exp -2\pi^2 (U_{11}h^2a^{*2} + U_{22}k^2b^{*2} + U_{33}t^2e^{it} + 2U_{12}hka^4b^4 + 2U_{23}k^4b^5e^{it} + 2U_{13}h^4a^*e^{it} )$$

Bond lengths and angles with their estimated standard deviations are listed in Tables 3 and 4, respectively. Figure 6 is a schematic diagram showing bond lengths and the numbering used in the Tables. From Table 3, it can be seen that the average standard deviation of bond lengths in this structure is 0.05 to 0.06 Å. In the benzene ring. the average carbon-carbon bond length is 1.39 Å. The C(sp3)-C(sp3) single bonds in the CDE ring system have an average length of 1.55  $\stackrel{\circ}{\Lambda}$  as do the  $C(sp_2)-C(sp_3)$  bonds at the junction of that ring system with the indolenine These lengths do not differ significantly from the moiety. expected values of 1.397,81.525 A, respectively quoted by Sutton et al. (1958).

The carbon-nitrogen bonds fall into three categories. The C=N bond is  $1.24\pm0.04$  Å long, not appreciably different from the value of 1.27 Å found in dimethyl-glyoxime. The  $C(sp_2)$ -N bond length of  $1.36\pm0.05$  Å is in agreement with the value of 1.371 Å reported in <u>para</u>-nitroaniline (Trueblood. Goldish and Donohue, 1961); 1.38 Å in 2-chloro-4-nitroanaline (McFhail and Sin, 1963); 1.39 and 1.40Å in ibogain hydrobromide (Arai, Copputs and Jeffrey, 1960). The value of 1.54Å for the average  $C(sp_3)$ -N bond length is semathat longer than the value given by Sabton of al., but the long

bonds of this type are frequently reported in alkaloids. In the structure of schitzmine (Hamilton, Hamor, Robertson and Sim, 1962) there are three bonds of this type with a longth of 1.54 Å and a fourth with a length of 1.56 Å. After comparing the results of several analyses of different alkaloids, these authors suggest that 1.52 Å is perhaps a better estimate of the length of such a bond.

Positional standard deviations and thermal parameters are greater for atoms in the side chains of this molecule. The bond lengths show broad agreement with the expected values The carbon-exygen single bond lengths can be placed in two groups; these adjacent to a carbonyl group are shorter ther these removed from the carbonyl. Average lengths are 1.33 and 1.44 Å respectively.

The equation of the weighted mean plane through the bousene ring is;

0.0195%  $\pm 0.1207$  -0.99252 = 0.1275Each above was given a weight inversely proportional to the estimated variance of the atomic position. Deviations from this "boot"plane are shown diagramatically in Figure 7. The units are  $\mathring{A} \times 10^3$ . The only atom in the indolonian modely that deviates significantly from the plane through the boundary that deviates significantly from the plane through the boundary what deviates significantly from the plane through the boundary that deviates significantly from the plane through the boundary what deviates between the bourses ring and the caushes bound would to keep the G=N group in the plane

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of the bonzone zing. C(7) would be expected to deviate from this plane because it is fully substituted and cannot take part in resemble interactions.

Intermolecular contacts under 3.6 Å are listed in Table. 5 and are shown in the packing diagram (Figure 8). For clarity, these have been divided into two groups and are shown as contacts from molecules I and I'; i.e. The molecule whose coordinates are given in Table 2, and the same molecule in the next cell in the positive x-direction. The only contact of particular interest is that between N(25) and O(30). This is a hydrogen bond of the O-H---type between the hydroxyl group at C(17) and the indefendance nitrogen of a symmetry-related molecule. The C...N distance (2.90:0.03 Å) is more reliable than many of the bend distances in the structure. The C-O...N angle is 103.8:1.7°.

Table 6 and Figure 9 show the environment of the indide ion. The ion has more than a desen contacts closer then the quoternary nitrogen with which it is formally associated. The N<sup>4</sup>,...I<sup>7</sup> contacts are quite long and some of the contests between the iodide ion and the atoms of the acctory group seem quite short.

The icdide ion is class? sidway between the quaternamy nitrogen of one molecule and that of the same molecule in the next coll in the negative z-direction. The N<sup>2</sup> . I

distances are 4.59 and 4.68 A sumeratively. However, the iodide ion dees not lie on a direct line between the two nitrogens but is displaced towards the acctory group. 143 closest contacts are with C(23), C(24) and O(31) of that Expected van der Waals separations are shown in group. the table. These wore calculated in the following way. Pauling's figures for the ionic radius of iodide (2.16  $\stackrel{\circ}{\Lambda}$ ). oxygen (1.40 Å), and a mothyl group (2.00 Å), product a minimum separation of 4.15 Å between iodide and a(24). and 3.66 Å between the ion and O(31). The van der Wacle saddes of a carbon with no hydrogen would probably be about half the inter-layer spacing in graphite, or 1.70 Å. Thus C(23) would not be expected to be closer than about 3-86  $\dot{\Lambda}_{\pm}$ Several of the observed contacts are very near the expected van der Vaals distances. The contact with C(24) is very In view of the high standard deviation in the shore. experiment and the uncertainties in the van der Weals distances very little significance can be placed on any individual observation.

Noth the long  $N^{+}$ ... I distances and the short contrate between iodide and carbon and iodide and oxygen seem to be observed frequently in compounds of this type. Silvers and talkesky (1963) report  $N^{+}$ ... I distances of 4.6 and 5.1 Å in eksemuidine; and is correction-II distributedide. MoPhall (1963) reports four distances between 4.42 and 4.36 Å.

Short contacts of iodide with carbon and oxygen have been observed by Przybylska (1961 a, 1961 b), Jellinek (1957), McPhail, Robertson and Sim, (1963) and others. This is probably a storic effect. It is possible that with an ion the size of an iodide ion, packing consideration and van dor Waals forces play a greater role in determining the environment of the ion than do forces of Coulombic attraction.

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# 

 $= \sum_{i=1}^{n} \int_{\mathcal{C}_{i}} d\mathbf{r}_{i} = \sum_{i=1}$ 

# Enhaddo el laberer al-box ebien.

Final about coordinates and bequietre factors. Standard deviations are given as write in the last place.

Atom	5		27.	172 172 174 174	$\mathbf{U}_{i \leq \mathcal{O}}$
	5 0.0482 0.0182 0.0534 0.0534 0.0534 0.0534 0.0162 0.02849 0.00000000000000000000000000000000000	9	Z 5774 1225 200 000000000000000000000000000000		
17(29) 07(20) 07(32)	~0.0338 ~0.8555 ~0.0793 ~0.32005		0.1949 0.1949 0.39633	-0.1078 -0.0855 -0.14516	0.051 0.058

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and the second s			SULS	803I	80 <sub>11</sub> .
	$\mathbf{u}_{\mathbf{r}}$ (3.2.5)	0.0205	0.0186	a0,0191	ad), (33)(34

# e \$ 3,40

# Table 3

Nothichids of 7-dectory7-H-Yonimbins: Bond Longths (A). Notimated standard Ceviations are given as units in the last place.

NÍ	25)	C(S)	1.242(36)
C	21	6/3	1.777(20)
cì	· ā (	N(26)	1.504/52
N	261	c/221	1. 602/621
54 W	26	0 ( <b>a</b> )	i soallaol
84 H	eo!	AR	1 E07 E01
1	2	A T	9 112 71
6	24	×\4	a enalist
			1.301 SU/
6	el el	6(9)	1 0114 321
5	91	5 16	1.201 391
G	, IU	U(II)	1.304 131
6		C(12)	1.0/1(KE/
6	18)	G(13)	1.3/0(4/)
C(	13)	C(C)	1.903(90)
U	13)	H (29)	1.301((40)
G	31.	G(14)	1.220(40)
G	[]恭]	6(15)	1.924((50)
6	15)	c(50)	1.393(29)
6		<b>C</b> (10)	1.5/3(2)
G	10)	GUY	1.592(40)
G	17)	C(18)	1.013(02)
G	(16)	G(19)	1.452(20)
G	197	C(20)	1.010((49)
G ((	(20)	G(21)	1.051(45)
G	(ZZ)	N(20)	1.602((62)
53(1	$(\underline{\mathcal{I}})$	G(2)	1.005((42)
5(	(7)	U(31)	1.401((43)
<u>O</u> (	(31)	0(23)	1.300 001
6(	(23)	C(24)	1.312(01)
6()		Q(z())	1.233()24
6		C(4)	1.405(0)
6	(46)	U(20)	1.803 (49)
<u>C</u> ()	(45)	u(29)	1.300((41)
<u>o</u> (	(23)	C(1)	1.412001
<b>C</b> !	(17)	U(30)	1.4301.431

# 

# Hothicaide of 7-Asetour-7-U-Yohnoice

Bond :	angles	azd	(cheir	stands	দেৱ ৫৬া	(iakions
	Nanocccccccccccccccccccccccccccccccccccc				- a Q M A + M A M A + M A M A A A M M A A A M M A A A M A	

r - 2 - 2 zz



Bond lengths in the Methiodide of 7-Acetoxy-7-H-Yohimbine.





Methiodide of 7-Acetoxy-7-N-Yohimbine: Weighted Mean Plane.

Deviations from the plane through the benzene ring are shown in units of angstroms  $\times 10^{-3}$ . The equation of the plane is:

 $0.0195 \times + 0.1207 \times -0.9925 Z = 0.1275$ .







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# <u>Table 5</u>

Methiodide of 7-Acetoxy-7-H-Yohimbine:

Inter-molacular Mon-bonded Distances (A)

Contacts are between the molecule in equivalent position 1 and a molecule related to it by the symmetry operation indicated by the equivalent poition number. The three numbers in brackets refer to the unit cell translations. The equivalent positions are given.

1.	R <sub>c</sub> y, Z <sub>c</sub>	
2.	1/2-x, -y, 1/2+2	Z
	1/2+E, 1/2-y, -2	Z
ħ.	-x, 1/2+y, 1/2-2	5

Contacts from equivalent position 1, cell (0,0,0).

		e.p.	cell	
C(10)	C(11)	N N N M	(0,1,0)	3.39 A
C(11)	O(27)		(0,1,0)	3.40
N(25)	D(30)		(0,0,0)	2.90
D(28)	C(19)		(0,0,0)	3.58

Contacts from equivalent position 1, cell (1,0,0).

C(5)	C(24)	2	(0,1,0)	3.52 A
C(17)	N(25)	3	(0,0,0)	3.53
C(19)	0(28) 0	3	$(O_s O_s O)$	3.58
C(24)	C(5)	2	(0, 1, -1)	3.52
D(30)	N(25)	3	(0,0,0 <b>)</b>	2.90



Environment of the Iodide Ion:

Contacts with atoms in equivalent position 1.

C(23)	3.84	A
C(24)	3.93	
0(31)	4.09	
C(3)	4.14	
C(5)	4.16	

Closest centacts with quaternary nitrogen

N(26)	1	(0,0,0)	4.59
N(26)	1	(0, 0, -1)	4.68

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# <u>Table 7</u>

The methiodide of 7-acetoxy-7-yohimbine HKL [Po] |Po] HKL [Po] |Po] HKL [Po] [Po] HKL [Po] [Fo]

2 27	0 10.3	13.8	630 0	16.0 11.9	12 4 0	12.2 13.6	341	108.9 117.2	7 11 1	28.6 27.7
2 28 2 30 2 31	0 11.2 0 10.0 0 18.2	9.0 11.9 16.8	7 2 0 7 3 0	8.8 9.3 38.5 43.7	12 5 0 12 6 0 12 7 0	9.6 2.3 12.1 15.3 27.5 32.2	361	96.2 91.7 113.3 111.5 14.6 20.8	7 12 1 7 13 1 7 14 1	41.5 41.8 51.6 50.7 19.0 14.0
2 33 3 1 3 2	0 7.9 0 71.8 0 124.6	10.5 75.6 145.9	740 750 770	29.6 26.3 68.5 71.6 49.4 51.4	12 11 0 12 12 0 12 13 0	18.7 18.3 10.3 14.5 22.9 20.6	3 8 1 3 9 1 3 10 1	57.4 55.5 52.6 52.9 38.5 52.7	7 16 1 7 17 1 7 18 1	15.7 18.3 29.1 34.1 26.2 27.1
334	0 25.6	33.1 150.3	780 790	56.4 63.9 83.3 73.1	12 14 0	11.3 8.6 18.3 20.7	3 11 1 3 12 1	66.3 57.7 26.1 21.0	7 19 1 7 20 1	16.3 15.5 10.3 9.3
367	0 54.6 0 21.8	54.5 28.8	7 11 0 7 12 0	42.8 45.4 16.8 18.2	12 22 0	8.4 9.7 8.4 11.6	3 14 1 3 15 1	75.9 75.6 37.8 48.1	7 22 1	35.0 27.1 15.7 13.3
3 8 3 9 3 10	0 75.2 0 70.6 0 51.4	75.8 74.7 42.0	7 13 0 7 14 0 7 15 0	73.1 60.9 22.2 14.6 52.9 47.4	13 2 0	13.3 13.1 13.2 15.8 14.7 15.3	3 16 1 3 18 1 3 19 1	44.8 46.8 54.0 45.4 35.3 41.4	7271 801 811	28.2 23.7 21.2 23.9 26.7 27.3
3 11 3 12 3 14	0 48.2 0 115.1 0 18.2	47.9 149.0 19.3	7 16 0 7 18 0 7 19 0	11.6 8.7 15.0 17.1 39.8 42.8	13 9 0	12.8 16.2 9.9 5.9 15.2 18.3	3 20 1 3 23 1	19.5 23.2 29.6 29.7 33.2 32.9	8 2 1 8 3 1 8 4 1	61.5 60.1 46.6 47.7 39.5 33.7
3 15	0 50.1 0 58.4	48.8 60.0	7 21 0	24.0 22.7 26.7 28.3	13 13 0	9.3 11.1	3 25 1	21.4 19.8	8 5 1	17.0 20.6
3 19	0 41.2 0 22.8	38.3 26.7	7 25 0	15.4 12.9 53.4 56.5	0 4 1	93.8 100.2 40.4 65.6	4 1 1 4 2 1	88.9 87.8 31.3 32.1	8 8 1	49.5 48.6
3 22 3 23	0 33.0 0 71.9 0 9.6	57.3 7.1	8 2 0 8 3 0	10.9 8.6	071	92.6 96.9 57.6 61.2 91.9 99.7	4 3 1 4 4 1 4 5 1	57.1 48.0 96.4 113.3	8 11 1 8 12 1	30.4 28.7 56.2 43.6
324 325 326	0 13.6 0 12.1 0 21.0	8.4 11.7 28.4	840 860 870	40.7 43.2 41.7 48.0 18.2 11.9	0 9 1 0 10 1 0 11 1	66.6 60.9 31.1 21.2 132.0 132.9	461 471 481	61.0 60.5 43.0 35.4 29.5 30.1	8 13 1 8 15 1 8 16 1	37.3 41.9 32.5 21.5 31.1 32.4
3 28	0 8.7	9.5	8 8 0 8 9 0 8 10 0	40.9 34.9	0 12 1 0 13 1	94.7 116.8 41.2 48.0	4 9 1 4 10 1 4 11 1	125.2 124.1 83.8 75.2 61.3 69.7	8 18 1 8 22 1 8 27 1	33.2 25.7 22.0 28.7 22.8 19.4
3 32	0 9.8 0 22.6	9.0 16.8	8 11 0 8 12 0	31.7 33.6 22.3 19.7	0 15 1	31.7 29.2 53.6 51.9	4 12 1 4 13 1	31.0 35.4 36.1 31.6	9 0 1	27.2 26.5
4 1 4 2 4 3	0 102.2	55.0 119.8	8 15 0 8 16 0	9.4 8.0 39.5 37.9	0 17 1 0 18 1 0 21 1	28.6 34.1 35.5 32.8	4 14 1 4 15 1 4 16 1	63.3 17.7 51.9 47.4	931 941	34.6 30.1 41.7 39.8
44 54 46	0 10.3 0 22.2 0 58.7	11.7 18.2 44.9	8 17 0 8 18 0 8 20 0	21.4 16.4 12.2 8.2 21.3 15.9	0 22 1 0 23 1 0 24 1	21.1 24.2 22.9 26.9 10.3 8.2	4 17 1 4 18 1 4 19 1	29.5 29.7 26.6 22.2 44.7 41.6	9 5 1 9 6 1 9 7 1	33.4 29.0 35.5 35.2 28.8 26.9
4 7 4 8	0 87.9	96.3 118.9	8 23 0 8 24 0 8 25 0	8.9 1.8 8.6 12.2	0 25 1	24.5 26.3 39.5 27.0 38.6 21.9	4 20 1 4 21 1 4 23 1	22.4 29.5 14.4 17.6 26.5 25.8	981 991	15.1 15.7 27.8 28.5 17.4 20.3
4 10 4 11	0 <b>36.9</b> 0 <b>49.7</b>	39.2 63.5	8 26 0 8 27 0	17.1 15.4	0 28 0 32	12.0 13.5 8.8 5.5	4 24 1	20.4 19.8 25.3 20.9	9 11 1 9 12 1	35.9 32.4 30.7 27.0
4 13 4 14	0 20.0 0 42.8 0 19.0	23.3 53.7 16.3	8 29 0 9 1 0	7.8 8.8	1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1 1	55.1 30.6 80.7 83.0	5 0 1	35.2 41.9 61.5 65.5	914 1	21.4 17.3
4 15 4 16 4 17	0 21.5 0 42.8 0 56.0	14.9 33.4 60.9	920 930 940	48.9 49.0 14.9 18.0 12.0 7.4	141	143.9 135.9	531	62.9 76.8 75.2 61.6	9 19 1 9 20 1	22.6 21.7 14.0 14.9
4 19 4 20 4 21	0 29.1 0 48.3 0 62.2	27.8 26.8 56.9	950 960 980	44.6 43.4 40.3 43.4 51.7 49.9	1 7 1	123.6 100.5 146.6 130.1 56.4 65.7	5 6 1	67.4 69.8 39.9 43.8	9 23 1 9 24 1	13.6 13.8
4 22 4 23 4 24	0 6.8	12.9 15.8 16.5	9 9 0 9 10 0 9 11 0	20.0 16.3 22.8 19.7 33.4 31.1	1 9 1 10 1 11	71.1 82.1 80.9 80.4 67.9 78.4	581 591 5101	36.8 41.2 53.8 58.0 76.6 86.0	10 0 1 10 1 1 10 2 1	24.9 23.3 31.1 32.9 15.3 12.1
4 27 4 28	0 28.8	30.1 13.8 62.4	9 12 0 9 13 0 9 15 0	38.9 35.6 27.0 19.6	1 12 1 1 13 1 1 14	71.7 77.1 29.0 39.3 55.3 55.9	5 11 1 5 12 1 5 13 1	65.4 62.8 32.9 25.6 24.0 26.7	10 3 1 10 4 1 10 5 1	19.3 18.3 17.3 18.9 51.4 42.4
5 3	0 124.6	127.0 76.9	9 16 0 9 17 0	40.1 40.8	1 15	45.9 38.5 53.9 46.2	5 14 1 5 15 1 5 16 1	56.8 51.6 58.7 63.5	10 6 1	28.5 29.7 15.8 13.8 14.2 17.3
556	0 31.0	39.9 100.0	9 19 0 9 20 0	14.9 13.8	1 18 1 19	34.0 33.8 37.7 27.3	5 17 1	27.3 25.1 28.0 25.5	10 9 1	46.0 47.7 18.0 20.5
5 7 5 8 5 9	0 53.0 0 14.5 0 122.9	33.4 17.8 100.3	9 21 0 9 22 0 9 23 0	0 11.1 8.8 0 17.5 16.7 0 10.6 13.8	1 20	51.5 63.7 28.3 24.8	5 21 1 5 22 1	14.5 20.5 16.3 14.4	10 13 1	24.5 25.1
5 10 5 11 5 12	0 31.6 0 35.6 0 79.2	30.9 47.4 89.4	9 25 0 10 0 0 10 1 0	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1 23 1 24 1 25	12.8 17.1 37.3 35.8 8.1 12.7	5 23 1 5 24 1 5 25 1	12.7 12.5 22.3 25.1	10 17 1	15.9 11.1 19.6 19.9
5 13 5 14 5 15	0 15.4 0 41.6 0 54.4	19.5 42.3 39.8	10 2 0 10 3 0 10 4 0	$\begin{array}{cccccccccccccccccccccccccccccccccccc$	1 26 1 27 1 28	11.4 16.9 17.5 20.1 13.5 12.9	6 0 1 6 1 1 6 2 1	19.5 18.2 56.8 72.8 69.5 74.3	10 19 1 10 20 1 11 0 1	19.2 21.0 11.7 13.5 41.3 42.9
5 16 5 17 5 18	0 39.8	33.5 24.7 36.9	10 5 0 10 6 0 10 7 0	5.4 $7.311.2 7.923.2 21.8$	132 20 21	13.9 15.5 48.4 59.2 60.8 76.2	631 641 651	57.8 61.0 46.0 45.7 20.5 15.2	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	10.2 16.9 16.1 9.2 18.0 15.0
5 19	0 48.2	30.7	10 8 0 10 9 0	10.4 10.6 21.0 14.9	2 3	152.5 149.4 161.6 140.3	6 6 1 6 7 1 6 8 1	61.0 60.0 62.5 64.5 44.8 43.0	11 4 1	18.0 14.0 21.4 20.6 29.5 32.1
5 23	0 10.7	10.5	10 11 0	13.5 15.9	220	85.8 102.0 91.6 97.9	6 9 1 6 10 1 6 11 1	41.2 39.9 37.7 40.2 36.0 37.5	11 7 1	16.3 17.7 21.6 21.0 29.0 27.2
5 20	0 13.0	7.6	10 14 0 10 15 0 10 16 0	10.0 $15.021.4$ $20.919.0$ $19.2$	28	41.8 48.7	6 12 1	56.3 56.9 25.4 22.0		29.0 22.6 12.9 15.4
066	0 18.1	30.7 125.9	10 17 0 10 18 0 10 21 0	0 11.7 7.9 0 12.0 14.5	2 10 2 11 2 13	13.0 12.1 34.1 33.9	6 15 1	48.9 45.8 39.0 34.1	11 15 1	20.8 12.4
667 68	0 102.9 0 93.9 0 32.9	104.2 99.9 29.7	10 23 0 11 1 0 11 2 0	) 9.6 12.3 ) 9.4 10.3 ) 34.2 35.9	2 14 2 15 2 16	59.9 62.0 52.5 55.9 46.4 44.4	6 17 1 6 18 1 6 19 1	30.5 33.5 29.3 23.8	12 2 1	16.3 15.8 14.6 14.2
6 9 6 10 6 11	0 16.2	21.5 65.9 35.2	11 4 0 11 5 0 11 6 0	6.7 8.9 15.0 19.0 34.6 34.8	2 17 2 18 2 19	25.1 25.1 42.3 46.5 35.7 36.2	6 20 1 6 21 1 6 22 1	16.3 23.6 26.5 21.6 24.5 25.3	12 4 1 12 5 1 12 6 1	16.3 12.2 18.2 24.4
6 13 6 14 6 15	0 22.2	19.5	11 7 0 11 8 0	21.4 23.4 34.8 41.9	2 20 2 21 2 22	45.2 36.0 40.4 37.7 22.7 24.4	6231 701 711	20.2 16.4 21.7 12.7 55.7 56.3	12 8 1 12 10 1 12 11 1	18.1 18.7 16.0 12.0 17.7 16.1
6 16	0 15.7	19.2	11 12 0	26.1 23.8 11.9 7.1	2 23	20.4 16.7	721 731 74	56.0 63.2 56.5 63.6 31.5 27.4	13 0 1 13 1 1 13 3 1	15.9 16.3 14.2 16.6 21.0 21.4
6 19	0 27.1	21.5	11 18 0	13.8 14.5	2 30	13.2 13.1 8.1 11.5	7 5 1	15.6 15.8 55.2 46.7 67.2 57 9	14 1 1 14 2 1 14 3 1	13.2 9.4 13.2 12.4 13.1 9.4
6 25 6 27 6 20	0 10.5	9.3	12 1 0	) 10.7 13.0 ) 19.2 18.5 ) 34.9 35.2	312	99.3 117.7 34.0 25.0 99.3 96.4	7 8 1 7 9 1 7 10 1	61.4 53.8 12.0 10.4 19.5 15.4	0 5 6 C	54.2 23.8 142.6 176.2 61.8 46.1

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	Tab	<u>te 7</u>		
The meth	iodide of 7-	-acotor -7-	H-yohimbine	Ctd.
H K L FO FC	H K L Fol Fc	H K L FO FC	H K L FO FO	H K L FO FC

# Pert III

The Structure of Benzilmonoximes The Crystal Structure of the <u>Para</u>-Bromobenzoate of Benzilmonoxime

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#### 3.1 INTRODUCTION

The currently accepted structures of the two monstanes of benzil are shown here



This configurational assignment, made by Meisenheimer (1921). was based on the fact that the benzoate of the Sociate, IV, is formed from the econisation of triphenylsexcase, III.



There is come evidence that the configurations should be reversed. The acctate of S-benzilmonoxime is stable at 190°C, whereas that of the a-oxime undergoes thermal decomposition at 140°C. When the reaction products are treated with vater, acctic acid, benzoic acid and benzonttrile are formed, indicating that the initial products were probably benzonitrile and acotic-benzoic achydrade. Such thermal decompositions of the occur by introductor -- 57

cyclic Mochanians, vis.



This would imply that the hydroxyl group in the actine was <u>cis</u> to the carbonyl group.

Some spectroscopic ovidence would seem to support An investigation of some a- and posting this argument. by infra-red spectroscopy (Palm end Vorbin, 1959), shows that with each isomeric pair, the B-oxime has the higher molting point and the higher conductivity. This incomes in acidity of the 0-M bond is accompanied by a shift we lover frequencies of the I.R. absorption bandassectated In most isomeric oximate with the 0-H strotching mode. the Georine absorbs at lower frequency than the Geissere. Bausilmonomimo is an exception on two counts. In solvaios of boussne and chloroform, the B-oxime absorbs at higher The molting point of the frequencies than the Goorine. p<del>ner-bronoboneasis of the</del> 3-oxime is only 107°C as computed with 195°C in the case of the C icomer.

Other anomalies in the behaviour of these were according to a character include an abarrand Booksam scarrangement (Block. 1990) and abarrand Booksam scarrangement (Block. 1990) and abarrand marking of the original market and abarrande model and the hold of the 1996 is the cond Minumbial 1996 is 1999 is N.N.E. studies undertaken at Ohic State University seemed to support a change in the configuration but these were not conclusive (M.S. Newman, private communication). Professor Newman requested an X-ray analysis and sent a sample alleged to be the para-bromobonzcate of  $\beta$ benzilmonoxime (m.p. 105-106°C). Barnes (1938) reports malting points of the para-bromobenzoates of  $\alpha$ -and  $\beta$ menoxime. His values are as follows:

> a form 105°C , 8 form 147°C .

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That the sample provided was the <u>para</u>-bromobenzoate of the a-form of the exime was further confirmed by the fact that attempted hydrolysis of the <u>para</u>-bromobenzoate <u>did not</u> yield the parent exime. (M.S. Newman, private computation). This behaviour is typical of derivatives of a-exists.

The crystals provided by Professor Norman were long white moddles and showed sharp extinction along the direction of the needle axis. Unfortunately, they showed a very strong tendency to gree together as twins. The specimen chosen for analysis had one individual of the twin sufficiently smaller than the other to allow necurate indexing. It had been intended to solve the structure by the usual heavy atom method. This approach failed The atmature was finally solved by combining a heavy

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atom-phased Fourier synthesis with a "minimum function" derived from a three-dimensional, sharpened Patterson function. Three electron density distributions were calculated to refine the atomic positions. Further improvement of the positional and thermal parameters was obtained by the method of least squares. The final R-value was 16-0% for 1258 independent, non-zero reflections.

The results of this analysis show that the <u>page</u> bromobenzoate of the a-benzilmonoxime (m.p. 105-106°C) has the <u>para</u>-bromobenzoate group <u>trans</u> relative to the carbonyl group (Figure 10).



This confirms Meisenheimer's assignment of structure to the eximes, I and II, of benzil.

#### 3.2. EXPERIMENTAL

The para-bromobenzoate of *c*-benzilmonoxime yields long colourless, needle-shaped crystals when crystallised from ethanol solutions. These melt sharply at 105°C on a hot-stege microscope. When examined under a polarizing microscope, they show sharp extinction plane the zoodle direction.

#### 3.2.1. CHOICE OF CRYSTAL

Preliminary photographs showed that the compound crystallises in the triclinic system, with space group Pl or PL (Subsequent calculations have shown that the space group is PI). Weissenberg photographs of the hkO zone were taken in the usual manner with multiple film packs (Robertson, 1943). Photographs of the hkl and hk2 reciprocal lattice nets showed that the crystals were twinned. This twinning occured in such a way that there was exact registry of points in the zero layer of the reciprocal lattice, but complete disregistry in other In the crystal chosen for analysis, one individual lavers. of the twin was sufficiently larger than the other to allow The spot shape of the reflections in accurate indexing. the photographs was very poor.

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#### 3.2.2. CHOICE OF UNIT CELL

In the triclinic system, the choice of unit cell is not dictated by symmetry considerations. For this reason there are an infinite number of ways of choosing a primitive cell in a triclinic lattice. However, for any crystal, there is one cell that is unique. This particular cell is called the "reduced cell" and is defined by the three shortest, non-coplanar translations in the lattice. In such a cell, each cell edge is shorter than reciprocal lattice points into the sphere of reflection. Thus both the  $\underline{a}^{n}\underline{b}^{n}$  plane and  $\underline{tbe}\underline{a}^{n}\underline{c}^{n}$  plane can be photographed with the same mounting. The dihedral angle between the planes is the complement of the angle between the plane normalst e.g.  $[\underline{a}^{n}\underline{b}^{n}] \perp \underline{c}$  and  $[\underline{a}^{n}\underline{c}^{n}] \perp \underline{b}$ ; therefore, the angle between the planes is 180-c. In the monoclinic case and in systems of higher symmetry  $\underline{a}^{n} = 180-\alpha$ . In the triclinic case

-57-

 $\mathbf{a}^* = \cos^{-1} (\cos \beta^* \cos \gamma^* - \sin \beta^* \sin \gamma^* \cos \beta).$ Thus the six reciprocal cell dimensions are available from a single mounting. The real cell dimensions can then be calculated from the formulae given by Buerger (1942. p.360).

### 3.2.4. CRYSTAL DATA

Formle	C <sub>21</sub> NO3Br	F.W.	<b>395</b> ,23	(3 <sup>72</sup> -
System	triclinic			
Space group	Hor R			
Lattice constants	$a = 8 \circ 78 \text{ Å}$			
	b = 17°01 Å			
	$c = 6 \cdot 27 \text{ Å}$			
	$\alpha = 95 \circ 11^{\circ}$			
	β = <b>91°08°</b>			
	- 102°340			

reciprocal lattice points into the sphere of reflection. Thus both the  $\underline{a}^{*}\underline{b}^{*}$  plane and the  $\underline{a}^{*}\underline{c}^{*}$  plane can be photographed with the same mounting. The dihedral angle between the planes is the complement of the angle between the plane normalst e.g.  $[\underline{a}^{*}\underline{b}^{*}] \perp \underline{c}$  and  $[\underline{a}^{*}\underline{c}^{*}] \perp \underline{b}$ ; therefore, the angle between the planes is 180-a. In the monoclinic case and in systems of higher symmetry  $\underline{a}^{*} = 180-a$ . In the triclinic case

 $\mathbf{d}^{*} = \cos^{-1} (\cos \beta^{*} \cos \gamma^{*} - \sin \beta^{*} \sin \gamma^{*} \cos \alpha).$ Thus the six reciprocal cell dimensions are available from a single mounting. The real cell dimensions can when be calculated from the formulae given by Buerger (1942, p.360).

#### 3.2.4. CRYSTAL DATA

Formule	C21NO3BE	F.W.	395.23	(3 <sup>53</sup> -
System	triclínic			
Space group	P. or PI			
Lattice constants	a = 8°78 Å			
	$b = 17 \circ 01 \text{ \AA}$		r	
	$c = 6 \cdot 27 \text{ Å}$			
	$\alpha = 95^{\circ}11^{\circ}$			
	β = 91°08°			
	$= 102 \circ 34^{\circ}$			

 Volume
 931 Å

 Density(obs)
 1.47 gm/cc

 Molecules/coll
 2

 Density(colc)
 1.49 gm/cc

 F(000)
 416

 µ(Gu Ka)
 35.9 cm<sup>-1</sup>

i , A Bas

## 3.2.5. INTENSITY DATA

Intensity data were obtained for the recipional lattice nots hkl to hks using a Nonius Veissenborg converand multiple film packs. To avoid the spot shape correction necessary if data are obtained from upper conlover halves of the same film, two sets of photographe were taken of each layer, one displaced from the other by a motation of 180°. The two sets were scaled to and another by comparison of common reflections. Wite magine recipreeal lattice acts were given a common relation such by comparison by axial reflections from the Velacoborg deta with common reflections from the procession photograph: Lorentz and polarization corrections of the Okl zone. had been calculated on the DEUCE computer for the Vetnession a Rease corrections had been applied to the processies data. data by means of a template. (Beerger, 1960, p.185).
## 3.2.6. ANALYSIS OF THE STELLTR

Projected Patterson syntheses. P(UV) and P(VV). vere calculated to determine three coordinates of the heavy The median axis projection, calculated from atom. precession data, showed two large peaks. Neither of these gave satisfactory agreement with the major peak in the short axis projection. A three-dimensional vector map. P(UVN), showed two major peaks corresponding to the two peaks in the projected function, P(UV). The larger of these was a diffuse peak, very much extended in a direction roughly parallel to the median axis. Since the smaller peak was the usual Gaussian shape it was assumed to represent the Brabe interaction. The bromino coordinates derived from this peak were;

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x = 0.3246, y = 0.0826, z = 0.3573.Structure factors were calculated with phases determined by the heavy atom. Siz's ratio, r, for this compound is  $r = \left(\frac{\Sigma r^2 Br}{\Sigma r^2 L}\right)^{\frac{1}{2}} = \left(\frac{1225}{1012}\right)^{\frac{1}{2}} = 100$ .

indicating that about 20% of the data will have the wrong sign in a calculation based on the heavy atom alone. (Sim, 1957). To avoid introducing errors in to the electron density distribution by including terms with the wrong sign. all structure factors with

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wore excluded from the first Fourier synthesis. The

R-value for the terms actually included in the summetion was 53°7%.

The first Fourier map was calculated in sections along the median only. There were 40 peaks with heights between two and six electrons/ $A^3$ . Other than the bromine atom, none of the known features of the molecule could be recognised.

A minimum function was derived from the three-dimensioned Patterson function in the manner described in Part I (1.4). Both the Fourier map and the minimum function were plotted on glass sheets, but the only obvious similarities between them were the positions of the bromine atoms and very pronounced streaks of electron density running diagonally across the maps. Although the minimum function was not well resolved, it was possible to reduce the number of atomic sites by comparing the two functions, one section at a time. Only 27 of the 40 peaks in the Fourier map corresponded to positive regions in the minimum function.



These 27 peaks were plotted in three-dimensions using long wires supported by a cork base. Two coordinates were defined by the positions of the wires relative to a grid drawn on the base. The vertical coordinate was marined by a small piece of fibre-glass electron a insulation shouthing. Tontative locations were assigned to 9 of the atoms in the selected. Hive of these were used in a structure-factor calculation with hk0 reflection data. The B-value increased. The three-dimensional data were "sharpened" by applying the Lorentz-polarisation function to the Patterson coefficients, P2. This gives



a sharpening function of the form shown on the left, The sharpened Patterson function was used to give an improved potition for the bromine atom and to derive a minimum function The bromine-phased electrondensity calculation and the Patterson synthesis ware consend in sections along the share

axis because it seemed likely that the strocking above to in the provious maps represented unresolved, steeply tilted bonzone rings. The resolution of the minimum function had been much improved by the sharpening process.

The electron density distribution and the similar function were superimposed by plotting equivalent sections from the two functions on the same glass sheet. This composite map revealed the positions of the 25 light struct Since the structure of the molecule was known, all of the stone could be identified. Then they were included

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in the next atracture factor calculation with scattering factors appropriate to their chemical types, the R-value dropped to 34.6%. Another cycle of Fourier refinement gave a further improvement to 30.1%. No further calculations were done on the DEUCE computer at Glasgov University.

#### 3.2.7. EEPINEMENT OF THE STRUCTURE

The Glasgow Structure Factor Least Squares program was available before the Neurier program. For this reason refinement was continued by least squares in spite of the high R-value.

The first cycle of losst squares used unit weights for each reflection and served to provide a rough over all scale factor for the data. The R-value was 35°7%. For all subsequent calculations, Cruickshank's (1961) weighting scheme was used:

 $fv = (1/p_1 + F + p_2F^2 + p_3F^3)^{\frac{1}{2}}$ 

The initial values of the parameters were

1

 $p_1 = 6.0, p_2 = 9.9 \times 10^{-3}, p_3 = 1.25 \times 10^{-4}$ These values were adjusted slightly during the course of refinement. The final values were

$$p_1 = 4 \cdot 2$$
,  $p_2 = 1 \cdot 6 \times 10^{-2}$ ,  $p_3 = 3 \cdot 8 \times 10^{-4}$ 

After four cycles of isotropic block-diagonal least-squares refinement, the R-value had been reduced to 26.6%. The date were rescaled. After two more cycles the R-value vas 24.6%.

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A progress for refining projecticus by the "minimum residual" method of Bluiys and Stanley (1963) vas veikkes by K.V. Muir of this Department. Much this program became available, it was used to refine the projections along the short and median axes. The R-velues for the two zones improved from 25% to 23% and from 27% to 22% Three-dimensional structure factors based respectively on positions from the two projections showed an R-value of 21.9%. Three cycles of least squares reduced this to 18.2%. Anisotropic temperature factors were allowed to refine for 4 cycles of block diagonal least squares. Refinement was terminated at this point. The R-value was 16.0% for 1258 independent reflections. Unobserved reflections were not included in the analysis.

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## 3.3 MOLECULAR DIMENSIONS AND DESCRIPTION OF THE STRUCTURE

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The configuration of the para-branabenzoate of bensil monoxime is best displayed when the crystal structure is viewed in projection along the short axis (Figure 10). Van der Waals forces are important in determining the packing arrangement of the molecule, but a very short contact between the oxygen of the carbonyl group and the carbon of the carboxyl group would seem to indicate that dipole-dipole interactions may also be important. Intermolecular contacts are seen most easily in the projection along the median axis (Figure 11). Figure 14 is an calargement of a portion of the median axis projection.

Table 8 lists the final atomic coordinates and their standard deviations as given by the least-squares refinement. The coordinates and their standard deviations are quoted with the same number of significant figures. No significance should be attached to the last figure as it has been included only to allow derived quantities to be calculated without less of accuracy. It should be noted that the standard deviations of the heavy atom coordinates are an order of magnitude smaller than those of the light atoms.

Anisotropic temperature parameters are shown in Table 9. These are the parameters of the equations

$$\sup \left( -\mathbb{R} \sin^2 \theta / \lambda^2 \right) = \exp \left[ -2 u^2 \left( V_{11} h^2 e^{\pi^2} + V_{22} h^2 b^{\pi^2} + U_{33} e^{\pi^2} e^{\pi^2} + 2 U_{33} e^{\pi^2} e^{\pi^2} e^{\pi^2} + 2 U_{33} e^{\pi^2} e^{\pi^2} + 2 U_{33}$$

The factor of 2 has been included in the cross-terms. Bond lengths and their standard deviations are given in Table 10; bond angles are shown in Table 11. These are shown diagramatically in Figure 12.

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Table 12 is a list of  $p_{ij}$ , the correlation coefficients between the coordinates of each atom. The correlation coefficient between the i<sup>th</sup> and j<sup>th</sup> parameters in a leastsquares calculation is defined as

 $\rho_{ij} = (a^{-1})_{ij}/[(a^{-1})_{ij}(a^{-1})_{ij}]^{\frac{1}{2}}$ , where  $(a^{-1})_{ij}$  is the j<sup>th</sup> element of the i<sup>th</sup> row of the matrix inverse to  $a_{ij}$ , the matrix of the normal equations. The intra-atomic correlations arise primarily from non-orthogonality of the area and from anisotropy of the thermal parameters.

Standard deviations of bond longths in this structure lie between 0.03 and 0.04 Å. Most of the observed bond longths are very near their expected values. Individual bond longths in the benzene rings vary between 1.30 and 1.51 Å. with an average longth of 1.41 Å. The C-Br bond longth is 1.85 Å. Both of these values compare favourably with the expected values of 1.397 and 1.85 Å, respectively (Saturn, 1965). There are three C(sp<sub>2</sub>)-C(sp<sub>2</sub>) single bonds in the Obtained which an evenego longth of 1.49 Å, and two 9=0 bonds with an everage longth of 1.21 Å. More measurements are an agreement with the values of 1.492 and 1.210 Å. respectively for bonds of these types in bonsil (Brown and Sadanaga, 1965). The C-O bond length of 1.38 Å is close to the value of 1.358 Å reported for the long bond in enters and acids (Subbon et al., 1958). The N-O bond length of 1.42 Å is similar to the value 1.408 Å, reported in Secondization (Levine, 1969), and to the value 1.415 Å, reported in Secondization extens (Hall, 1965).

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In a summery of oxime structures, Hamilton (1964) Hith various longths for the C=N bond between 1.225 and 1.28 The value predicted from Pauling's radii is 1.265 measured bond longth of 1.35+03 A seems to be signal Phonenia langer than the predicted length. This may be real of the Hay reflect an under-estimate of the standard deviation of Who boad length. Hedgeon and Rollott (1963) have show they when the wariances of parameters and calculated from the tuvozene of the block metrices, stenderd deviations of bork Asspend one be underestimated by as much as 80%. X8 43.0 AARCONORCE is real, it implies that this is not a prave double hand but that it is taking part in some serie of Theman .... antoreethon. This is not the ease, Resonance interve while the honzone sing is provinted because both H(b) of Title desings significantly from the plane through th

ring (Table 13). Becomence with the carbonyl group is impossible. The root mean square distance from the "best" plane through C(8), C(9) N(22) and O(23) is more than (0.25 Å). This would indicate that the difference is not real and that the standard deviation of this band length and been underestimated.

"Bost" planes have been calculated for the three benzene rings and for the carboxyl group. Deviations from these planes are shown in Table 13 and Figure 13. It is interesting to note that G(17) and Br(26) lie in the plane of benzene ring A, but that the engloss of the carboxyl group have been displaced from this plane by a rotation of 13<sup>2</sup>.

Peuling's figures for the Van der Weals radii of carbon (1.70 Å), nitrogen (1.50 Å) and oxygen (1.40 Å) predict that the shortest inter-molecular contacts will be as follows:

CC	3.40 Å	GN	3020	A
ØO	3.12 Å	00	2-80	e Ac

The only contact that is significantly shorter then due convergending Van der Waals redius is that between O(24) and O(2). The calculated standard deviation of this measurement is just under 0.03 Å, but the difference between the observed and calculated distances is large enough to be significant even if the standard deviation were as high as 0.04 Å. Cantacts as short as 2.77 Å between  $sp_2$ -hybridized enclose care and charter of a carbonyl group have been observed to address

compounds where there is no possibility for hydrogen bonding. (Davies and Blue, 1955); Ohn. Joffrey and Sakurai, 1962; Bolton, 1963, 1964, 1985). Bolton has studied this effect and has suggested that it occurs when dipole-dipole or dipole-induced-dipole interactions are quite strong. The shortest contacts occur when the bond of the carbonyl group is approximately collinear with the W-bond of a carbon atom in an adjacent molecule. In the compounds studied by Bolton, the carbon atom is in an aromatic system with at least two carbonyl substituents. The average longith of the 0....C contact is 2.82 Å, and the average value of the Column angle is 1M. In the present study, the carbon ston involved in the contact is an sp-carbon in a carboxyl group. Thore is no large inductive effect present, but neither of the C=O groups in the molecule seem to be involved in any appreciable resonance interaction. It seems reasonable to expect a cortain amount of dipole-dipolo interaction between the two groups. This interaction would not be as strong as in alloxan and the other molecules studied by Bolton. the C=9...C angle (116°) is less favourable and the contant is longer then those he observed.

Kihara (1963) has shown that the structures of multipolar molecular crystals are governed by electrostatic forces between the molecules. No has determined correct erystal structures for a number of relatively simple molecules by building models with shape and magnetization appropriate

to the particular molecule and allowing the models to "exystallize". This would seem to indicate, that electrontatic forces play an important tole in determining the crystal structures of a fairly broad category of molecules. This general observation, coupled with the fact that the carboxyl group in the present structure has been rotated out of the plane of the benzene ring in a direction that would make a more favourable angle between the M-orbital of the carbon atom and the q-orbital of the oxygen, suggests that the close contact is real.

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Packing diagram of the <u>para</u>-bromobenzoate of benzil monoxime, viewed along the <u>c</u>-axis

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a







Br(26) 0.31968(41)

0.08090(18)

1.3570(49)

			*
	0.2293(35) 0.3015(30) 0.2336(29) 0.0974(27) 0.0197(31) 0.0534(31) 0.0534(31) 0.0534(31) 0.0504(27) 0.0604(27) -0.0601(27) -0.0601(27) -0.1020(29) -0.2271(36) -0.3264(33) -0.3040(36) -0.3040(36) -0.3040(36) -0.3110(23) 0.3110(23) 0.4064(24) 0.5643(33)	=0.0307(14) $=0.0077(16)$ $=0.0692(17)$ $=0.1249(14)$ $=0.1250(14)$ $=0.1892(14)$ $=0.2860(13)$ $=0.4829(14)$ $=0.4829(14)$ $=0.4829(13)$ $=0.4829(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4839(14)$ $=0.4859(14)$ $=0.2459(14)$ $=0.2459(14)$ $=0.2459(14)$	$\begin{array}{c} 1.1503(35)\\ 0.9457(45)\\ 0.7852(45)\\ 0.7852(45)\\ 1.0551(37)\\ 0.8263(37)\\ 0.8263(37)\\ 0.8263(37)\\ 0.8263(37)\\ 0.83637(37)\\ 0.83635(37)\\ 0.0527(35)\\ 0.0527(35)\\ 0.0557(37)\\ 0.0555(37)\\ 0.055($
C(17) C(18)	0.4064(24) 0.5693(33)	-0.2459(14) -0.2455(16)	
C(19) C(20)	0.6244(25)	-0.2042(10) -0.2907(15)	0.4507(43)
C(81)	0.3674(20) 0.0373(21)	-0.2909(19) -0.2576(11) -0.3356(10)	<b>0.35</b> 70(28) -0.1263(25)
U(23) D(24) D(25)	0.1143(16) -0.1135(19)	~0.2309(10) ~0.2309(10)	0.5207(23) 0.7229(28)

Final fractional atomic coordinates with their estimated standard deviations. No significance should be attached to the last figure as it has been included only to allow derived quantities to be calculated without loss of accuracy.

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# Table 9

<u>Para-bromobenzoate of Benzilmonoxime:</u> Anisotropic Temperature Parameters

	U 40	U <sub>22</sub>	U <sub>33</sub>	20,3	20 <sub>31</sub>	20 <sub>12</sub>
C(1)	0.156	0.073	0.063	-0.064	0.035	0.085
c(2)	0.079	0.086	0.116	-0,010	-0 009	0.028
č(3)	0.060	0.103	0.139	0.010	-0.004	0.009
C(4)	0.074	0.079	U.U82	0.006	0.030	0.045
C(5)	0.095	0.110	0.087	0.014	0.024	0.069
C(6)	0.115	0.063	0.082	-0.07	0.025	-0.007
C(7)	0.072	0.087	0.079	-0.039	0.038	0.023
C(Å)	0.067	0.075	0.064	-0.007	0.022	0.029
C(9)	0.087	0.083	0.052	0.041	0.040	-0.004
C(10)	0.083	0.074	0.076	0.011	0.078	0.039
C(11)	0.088	0.046	0.114	0.005	0.041	0.022
C(12)	0.120	0.077	0.126	0.001	0.022	-0.022
C(13)	0.090	0.089	0.153	-0.008	0.013	0.022
C(14)	<b>0.</b> 108	0.123	0.120	0.002	0.012	0.050
C(15)	0.053	0.138	0.077	-0.062	-0.023	-0.005
C(16)	0.054	0.080	0.086	-0.020	-0.020	0.016
C(17)	0.047	0.088	0.084	0.019	-0.001	-0.044
C(18)	0.105	0.092	0.076	0.020	-0.016	-0.036
<b>C</b> (19)	0.083	0.097	0.069	-0.018	=0.010	0.001
C(20)	0.050	0.089	0.116	-0 <b>.</b> 026	-0.002	-0.003
C((21)	0.061	0.090	0.097	-0.020	0.026	0.012
N((55)	0.074	0.071	0.070	0.022	0.010	0.005
0(23)	0.080	0.110	0.088	0.015	0.009	-0.003
0((24)	0.059	0.093	0.083	-0.042	0.010	-0.004
0(25)	0.090	0.098	0.110	-0.020	0.065	=0.004
Br(26)	0.136	0.096	0.102	a0.060	-0.011	0.014.

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## Table 10

Fara-bromobenzoate of Benzilmonorime:

Bond lengths in angstroms, with standard deviations given as units in the last place.

Br(26)	C(1)	1.845(25)
C(1)	C(2)	1.448(37)
C(2)	C(3)	1.390(39)
C(3)	C(4)	1.446(35)
C(4)	C(5)	1.449(35)
C(5)	C(6)	1.301(36)
C(6)	C(1)	1.477(39)
C(4)	C(7)	1.507(33)
C(7)	0(25)	1.197(28)
C(7)	0(24)	1.377(27)
o(24)	N(SS)	1.417(23)
N(22)	c(8)	1.354(28)
C(8)	C(16)	1.495(30)
C(16)	C(17)	1.433(32)
C(17)	C((18)	1.387(35)
C(18)	C(19)	1.316(35)
C(19)	C(20)	1.475((35)
C(SO)	C(21)	1.377((32)
C(21)	C(16)	1.357((34)
C(8)	C(9)	1.456((30)
C(9)	0(23)	1.217((26)
C(9)	C(10)	1.469((32)
C(10)	C(11)	1.420((34)
C(11)	C(12)	1.447((38)
C(12)	C(13)	1.375((44))
C(13)	C(14)	1.380((46)
C((14)	C(15)	1.511((4))
C(( 15)	C(10)	1.323((33)



<u> 2ara-Sromobenzoate of Benzilmonoxime: Bond Lengths and Angles.</u>

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# <u>Fable H</u>

## Para-Bronobanzoato of Menzilmonomimes

Bond angles with their artimated obsect. deviations given as units in the last place.

		angle	E.s.d.
BCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC		573 7984779795730768848305778888483057 99487615638890895873076884830577588184184818 111111111111111111111111111	annan an an an an an an ann an an ann an a

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# Table 12

# Para-Bromobenzoate of Benzilmonoxime:

# Correlation Coefficients

C(1)	0.29998	0.14586	-0.16655
cla	0.11687	0.06582	0.01953
clas	0.08878	0.04642	0.07248
č14)	0.18409	0.17467	0.07385
C(5)	0.25684	0.12183	0.09247
<b>ด้ไอ้</b> โ	0.00550	0.12405	0.05836
C (7)	0.13869	0.18289	-0.07714
c(8)	0.16573	0.12154	0.03710
C 95	0.01546	0.16854	0.21003
C(10)	0.18099	0.30766	0.08628
C(11)	0.11643	0.18964	0.06149
<b>C</b> (15)	-0.02487	0.09343	0.09301
C((13)	0.12420	0.09403	0.05286
C((14)	0.17601	0.07558	0.06221
C((15)	0.03727	-0.03303	-0.13106
C((16)	0.09594	-0.00379	0.00554
C((17)	-0.10309	0.05197	0.13659
C((18)	-0.09101	-0.02412	0.15814
C((19)	0.05061	-0.01672	0.01787
<b>C</b> ((20)	<b>0.</b> 05333	0.07642	~0.03691
C((21)	0.07686	0.13350	⇔0.01227
N((SS))	0.07201	0.09923	0.13074
D((53)	0.05114	0.08290	0.11402
0((24)	0.04473	0.08099	au.00900
0((25)	0.04323	0.24747	0.02312
er(50)	0.04472	0.00410	-0.10339

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## Table 13

Para-Brom	benzoate	of Benz	ilmonox	ine: dev	iations	in
Atoms inc.	luded in	om prane B plane	0 611265	C	197 C M M M L C +	D
C((1)) C((2)) C((4)) C((4))	- 19 8 30 - 60 50 - 9	C(16) C(17) C(18) C(19) C(20) C(21)	2 -29 12 7 18 36	C(10) C(10) C(12) C(13) C(14) C(15)	19 -19 -13 -19 -13 -19 -19 -19 -19 -19 -19 -19 -19 -19 -19	0(24) 0(24) 0(27)
Root mean	square d 34	eviation	: 27		3 <b>4</b>	
Atoms not C(7) Br(26) O(24) O(25)	included -43 70 312 154	in plan C(8) C(9) N(22)	e: 205 ~559 1347	C(9) D(23)	159 525	

Equations of planes. X,Y,Z are atomic coordiates orthogonalized according to the following formulae:

	$\begin{array}{l} \mathbf{X} = \mathbf{a}\mathbf{x} + \mathbf{c}\mathbf{z} \cos \beta, \\ \mathbf{Y} = \mathbf{b}\mathbf{y}, \\ \mathbf{Z} = \mathbf{c}\mathbf{z} \sin \beta \end{array}$
A	-0.5217X + 0.7644Y - 0.3788Z = -4.4150
B	-0.1411X -0.9093 Y -0.3914 Z = 3.8857
С	0.5917 X -0.7604 Y -0.26772 = 4.5940
D	

-



<u>Fare-</u>Dromobenzoate of Benzilmonoxime: Deviations from the planes through the benzene rings are given in unit of angstroms X10<sup>-3</sup>

## 资源处理:20 14

Tera-Bronchesurate of Beneilescoorime-

inter-molecular Coatse is < 5.60 A.

Equivalent position I refers to the molecule whose coordinates are listed in Table S; position 2 refers to the polecule related to this one by a centre of inversion.

C(9)	0(25)	Ť	( 0, 0,-1)	3.44
C(13)	C(14)	2	(-1,-1, 0)	3.49
<b>6(</b> 14)	C(14)	2	(-1,-1, 0)	3.53
C(19)	0(25)	-	(1,0,0)	3.53
C(S0)	0(25)	Ţ	(1,0,0)	3.19
)«(SS)	C(5)	1986	(0,0,-1)	3.30
0(23)	0(4) C(7) C(12) C(13) C(21) O(24) O(25)		$ \begin{pmatrix} 0_{2} & 0_{2} & -1 \\ 0_{p} & 0_{p} & -1 \\ 0_{p} & 1_{p} & 0 \\ 0_{p} & -1_{p} & 0 \\ 0_{p} & 0_{p} & -1 \\ \end{pmatrix} $	A. 99 A. 99 A. 339 A. 349 A. 3
0(24)	C(6)	3	$(0, 0_{s}-1)$	3.42



Inter-molecular contacts (median axis projection)

-87-Table 15

Perg-Bromohopro to of Renzilmonoxime

HE KE LL	Fo Fo	E K L	Fo Fc	6: K L	Fo Fc	进放工	Fo Fc
5 8 J 4 11 U	11.7 -11.2	3-1 0 7-2 0	24.4 13.7 5.1 5.6	5 13 -1 5 14 -1	4.7 -6.6 4.4 -4.6	6 -9 -1 6-12 -1	12.5 13.6
1-12 () 1-14 ()	9.1 -9.1 5.8 8.0	3-4 0 3-5 0	32.6 -31.6 18.7 -18.7	6 1 -1 6 2 -1	11.9 -12.0 10.8 10.3	6-13 -1 7 0 -1	3.3 -5.2
22–165 ⊍ 33–33 ©	4.2 -4.8 5.2 -5.0	3-6 u 3-7 0	22.8 -25.6 30.7 -31.0	64-1 65-1	7.8 -9.3 10.0 11.0	7 -5 -1 7 -6 -1	11.5 -13.4
3-11 ú	14.5 18.4	3,48 ú 7,48 ú	10.9 11.9	6 6 -1 6 7 -1	19.0 19.0 6.5 7.0	7 -8 -1	3.3 2.2
4≟14ű ű 4a–15 ű	84 7.2	3,–1ບັບ 3–12 ບ	5.9 6.1 10.2 12.8	58-1 69-1	5.5	ŝ -3 -1	5.8 5.6
7 – É u u 2 u	4.1 -3.9 9.5 -8.1	3-13 0	5.6 6.7	6 11 -1 6 12 -1	3.5 -4.6	8 –6 –1 8−10 –1	7.3 -7.6
5 <b>3</b> 5	18.7 13.2	3-15 U	5.9 5.7	7 1 -1	2.0 4.5	0-12 1 0-11 1	6.7 -8.1
	63.2 -61.6 24.7 -26.3	4-20	51.8 49.2 33.1 30.7	73-1	5.8 6.2 10.3 -10.2	v⊸10 1 v ⊸0 1	15.1 -15.9
យាយំយ យដា យ	15.6 13.7	4,⊶£i 0 4,⊶5, 0	10.5 12.5	77-17	5.1 5.5	ນ - ອີ້ 1 ບ - 71	16.3 -17.9 24.0 29.5
013 C 14 0	3.9 6.7 8.3 No.9	4-60 4-70	25.0 27.3	711-1	3.2 4.7 3.3 1.8	⊎–6,1 ಲ~5,1	51.1 49.8 13.5 11.8
U 15 J 1 Q U	5.9 5.8 38.0 43.2	4-10 L 4-11 C	5.4 -7.6 5.4 -4.3	83-1	5.7 -4.5	્ર-4 1 ચ-3 1	10.6 9.5 68.9 34.9
120 140	56.1 -44.7 35.2 -30.1	5-1 U 5-2 U	12.8 -12.1	-16	5.2 5.0	ତ <b>- ଅ</b> ା ତ ଅା	43.0 -38.0
150	10.7 9.9	5-3 0 5-7 0	21.2 -24.9 17.7 17.6	S−16 1 S−15 P	7.3 7.1	ະ <b>4</b> 1	42.2 27.7
170	33.2 31.7 12.4 13.3	5-8 u 5-11 u	21.3 20.0 9.8 NO.4	1-3-1	75.0 69.0	2 <b>0</b> 1	14.2 11.0 22.4 19.2
1 9 0 1 10 9	51.2 59.5	- <u>5</u> -3°°	7.3 -6.9	1-8-1	3.2 -3.3	ິນກາງ	7.4 -9.1
1 12 0	3.9 4.8	ອ64 ສຸຊຸ	17.0 -20.2	1-10 -1	15.4 15.2	0131 0141	10.2 -12.4
114 0	10.2 -13.6	5-10 0 6-10 0		1-12 -1	15.6 14.2	0151 1-131	6.7 -5.8 7.9 -10.1
210	25.7 19.1	6-11 0 7 -7 D	11.8 14.8	1-14 -1	ê.î 8.4	1-10 1 1-9 1	14.9 19.4
2 3 3	26.5 22.3	7-4 0	14.5 12.6	1 0 -1 2 0 -1	54.6 68.7 9.4 -11.9	1-81 1-71	11.4 14.6 21.1 21.4
19.00	22.0 18.7 11.5 12.5	541	8.6 -9.1 6.0 4.5	2 -1 -1 2 -2 -1	3.8 30.2 15.3 12.8	1-6 1 1-5 1	31.2 35.7 44.4 -43.9
2702	15.6 16.2	1 -2 -1 2-15 -1	13.5 -12.8 7.2 7.1	2 -3 -1 2 -4 -1	77.6 65.7 20.1 18.7	1-4 1 1-3 1	92.3 -111.5 56.5 -76.1
2 9 3 2 9 3	9.8 -14.0 19.2 -14.6	4-12 -1 1 2 -1	4.8 3.6 35.5 39.0	2 -5 -1 2 -7 -1	71.8 62.7 36.0 -27.0	1 -1 1	3539.3
2110 2120	14.5 -15.5 7.0 -10.7	13-1 14-1	32.4 36.1 26.9 -23.2	2 -8 -1 2 -9 -1	15.1 -12.0 33.1 -23.0	121	51.7 -45.3 66.1 56.4
213 G 214 O	5.9 -8.1 7.3 -11.5	1 5 -1	58.4 -54.0 39.8 -34.2	2-10 -1 2-11 -1	50.2 -35.3 29.4 -20.8	151	6.2 9.3
3 <b>1</b> 8	70.2 59.9 33.6 37.6	1 7 -1	5.5 -9.0 42.1 -36.3	2-12 -1	5.55	171	11.4 12.9
1997 1997	39.0 -44.4 21.5 -22.8	1 11 -1	13.3 8.1	2-17 -1	62 77	191	10.6 -13.5 18.3 -23
1991 1997 1997	7.0 -10.5	1 13 -1	16.6 13.5	3 -1 -1 3 -2 -1	13.2 -12.5 3-5 -23.4	1 11 1	24.5 -26.6 9.4 -11
311 0	9.1 15.2	2 1 -1	2727.3	3 -3 -1 3 -4 -1	33.0 -49.3 4.9 4.8	1131 1151	4.6 -3.9 6.6 6.1
4 5 0	33.4 -30.2	2 3 -1	50.4 -53.4 4.6 -6.2	3 -5 -1 3 -6 -1	21.2 17.7 46.7 33.6	2-12 1 2-11 1	11.9 16.2 7.1 11.9
4 3 0 4 4 3	25.9 -31.9 13.2 -16.6	2 5 -1 2 6 -1	46.0 -48.0 21.7 -22.3	3 -7 -1 3 -8 -1	23.8 -15.5 47.1 35.0	2-51	28.2 - 32.1
4 5 J 4 6 J	6.7 -8.4 3.5 2.9	2 7 -1 2 8 -1	3.4 €.3 28.2 29.2	3 -9 -1 3-11	14.8 8.5 6.1 4.7	2-07	43.2 44.9
4 7 U 4 8 C	8.2 11.8 9.4 17.0	29-1 210-1	17.6 15.9 10.1 6.9	3-11 -1 3-12 -1	13.9 -12.4	2-31	29.5 -25.4
4 9 5 4 10 5	8.0 9.7 7.1 11.4	2 11 -1 2 13 -1	5.8 -5.9	3-13 -1	(.9 -0.) 9.1 -10.1	2-11	27.9 25.4
520	35.0 -34.9	2 14 -1 2 15 -1	6.4 -4.9	4 0 -1	25.4 25.1	211	42.3 4.2 19.5 2.0
5 3 4 i 2	10.2 12.7	3 1 -1	17.2 -16.3	4 -2 -1	5.3 -6.4 12.9 -29.9	231	11.8 12.0 16.0 -25.4
997 1995	7.7 10.5	3 3 -1	21.7 21.6	4 -4 -1 4 -5 -1	16.4 -12.3 38.8 -33.1	251	440.5 -39.2 38.2 -40.3
510 0	5.9 -7.7 8.4 19.1	3 5 -1	4.1 5.4 14.2 15.3	4 -6 -1 4 -7 -1	27.1 -25 47.4 -35.8	271	32.7 -34.4 24.2 -21.5
6 2 4 6 6 4	8.6 11.0	39-1 310-1	10.5 -11.1 7.6 -6.9	4⊷10 -1 4-11 -1	17.4 17.0 17. <u>3</u> 16.6	291	10.7 9.5
670720	5.9 -4.1 8.3 -10.1	3 11 -1 3 12 -1	11.4 -13.2 9.1 -7.8	4-13 -1 4-14 -1	5.7 0.7	2131	9.4 6.7 9.4 6.7
730	8.4 -10.0 31.8 -25.2	3 13 -1 3 17 -1	2.9 -3.4 3.6 3.9	5 6 -1	5.5 -6.8 6.5 -6.8	3-91	9.1 -13.4
1-2 0 1-3 0	27.5 21.9 51.1 45.3	3 19 -1	2.5 1.1 40.4 41.7	5-2-1	10.2 11.4	3-6 1	40.7 40.4 34.6 34.1
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1-80 1-90	27.0 -20.7 6.6 -9.1	4 5 -1	6.4 6.4	5 -6 -1	3.0 -3.8 20.3 -16.8	3-11	24.6 -19.5 8.4 -5.2
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1-15 U 1-16 (	7.3 7.4	4 9 -1 4 12 -1	9.8 -11.3 5.8 6.3	5-10 -1 5-12 -1	10.2 -8.9 5.7 -5.8	337	40_0 -40_0 6_4 -6_4 10_3 11
2 3 5	29.4 -26.7 38.6 34.2	4 14 -1 5 2 -1	6.1 5.2 25.3 -23.6	5-13 -1 5-14 -1	7.3 9.0 2.9 2.9	371	15.7 17-4 18.6 10 3
2 -6 0 2 -7 0	8.9 -5.6 20.2 21.2	5 3 -1	6.9 -8.1 7.9 -8.6	5-15-1 6 9-1	2.0 2.7	391 391	27.7 32.1 18.2 22.5
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$\begin{array}{cccccccccccccccccccccccccccccccccccc$	19.8 -10.3 -6.7
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$\begin{array}{cccccccccccccccccccccccccccccccccccc$	-11.9
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Table 15 Ctd.

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нкг	Fo Fc	нкі	Fof Fc	нкс	Fo Fc	нкг	Fol Fe
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-89-Table 15 Ctd:

SECTION II

A Comparative Study of 1,l'-Binaphthyl and the β-form of Perylene

### INTRODUCTION

Coulson has pointed out that a bond in a conjugated system may be regarded as a superposition of a s- bond and a fractional T- bond. For a particular type of hybridization, the length of a bond can be expected to vary with change in T- bond order. Various empirical relationships have been suggested to correlate bond length with bond order. Since most of the accurate structure determinations have involved bonds of informediate T- bond order, agreement between theory and experiment is very good in this region but becomes less reliable for extreme bond orders.

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Binaphthyl, I, and perylene ,II, are examples of molecules that contain bonds of very low T- bond order. Both of these molecules can be regarded as composed of two naphthalone residues connected by long bonds.



There are mine unexcited Kekulé structures for each molecule, representing the various combinations of three Kekulé structures of each napthalone nucleus. None of these structures attributes double bond character to the linking bonds. Thus simple valence bond theory predicts that the long bonds are pure  $sp_2 - sp_2$  single bonds with no 11 bond character. Molecular orbital theory predicts that the We bould order will be low, but it cannot be zoro since this theory assumes that there is always some interaction between We bonds on adjacent atoms.

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A study of these two compounds provides a fairly stringent test of the hypothesis that the aromatic character of the perylene molecule is localized in the nephthalene residues. Perylene is planar but binophyhyl is not. The angle between the planes, of the nephthalenic residues is about  $68^{\circ}$ , sufficiently large to proclude any  $\pi =$  bonding interaction between the two ring systems. If there are significant differences between the dimensions of the two molecules, this could be interpreted as evidence for interaction between the nephthalenic residues in perylene.

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## THE CRYSTAL AND EDLECHAR STEDOTURE OF 1.1' BINAPHTHYL

The crystel structure of 1,1 - binaphthyl (n.p. 148°C) was first established from projection data by Brown, Trotter and Bobertson (1981). They have reported that the molecular configuration is <u>cig</u> with an angle of about 73° between the planes of the two ring systems, and have suggested that the variety of melting points reported in the literature corresponds to different crystalline forms. Brown has refined the structure in two projections by differential synthesis but his results have not been publicated.

Badar, Chua, Cooke and Marris (1965) have shown that the compound has at least two distinct crystalline vector. The higher molting form (m.p. 197-159°C) was propared by deamination of (+)-4,4°-naphthidine and is eptically active. When molted and resolidified it is converted into the low molting form (m.p. 145°C). Infra-red spectra of the two forms are identical in carbon disulphide, in carbon tetrachloride and in cyclohexane, but there is an extra peck in the solid state spectrum of the low molting form. The authors attribute this to C-H out-of-plane vibrations cauced by  $H_{co.H}$  and  $G_{co.H}$  interactions inherent in the gig configuration. They suggest that the higher molting form

## 1.1. EXPERIMENTAL

Dr. Marris kindly provided samples of two exystalling modifications of 1.1'-bitarbibly1. The high-molting form does not show extinction when examined under a polarising microscope. When a "single" crystal is photographed by the usual rotation method the result is a rather "spotty" powder diagram, typical of coarse, randomly-oriented crystallites. Recrystallisation from light petrolema did not improve the crystals.

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The low-melting form crystallises as colourless plates that show sharp extinction and can be cut without apparent deformation. Weissenberg and procession photographs vore of one of these crystals mounted about its unique axis, used to check cell dimensions and space group. Cell dimensions were further confirmed by measuring reciprocal lattice spacings of high angle reflections on the linear diffractometer.

#### 1.1.1. <u>GRYSTAL DATA</u>

Formula C <sub>20</sub> H <sub>14</sub>	R. V.	<b>25</b> 4+3			
System	monoclinic				
Lattice constants	$a = 20.97_8 \text{ Å}$				
	$b = 6 \circ 35_1$ Å				
	c = 10.128  Å				
	β =105°17°				
Absent Spe <b>ctra</b>	hk\$ if h+k odd				
	h0 <i>% if %</i> odd				
	Oh) if k odd				
Space group	Ce or C2/e				
Volume	1302°3 A				

Density (obs.) 1.299 gm/ceMelecules/cell 4 Density(calc.) 1.297 gm/ceLinear absorbition coefficient,  $\mu = 0.86 \text{ cm}^{-1}$ (No Se,  $\lambda = 0.7107 \text{ Å}$ )

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### 1.1.2. <u>REFLECTION DATA</u>

Reflection date from eight reciprocal lattice nets, h04 to b74, were collected with No KG rediction on a Hilger-Natts linear diffractometer of the type described by Arndt and Phillips (1961). Peak height and background were measured at least twice for each reflection. Flames with intensities smaller than ten counts per minute were considered unobserved and were not included in the refinement. Lorents-polarization corrections were applied to the observed data by computer program using the formals expression to equi-inclination Veissenberg geometry.

#### 1.1.3. THE HILGER-MATTS DIFFRACTONETER

The diffractometer is of the moving-crystel-stationarycounter type and can be described as a mechanical model of the reciprocal lattice. Its main components are three slides driven by a system of tapes, pulleys and gears. The motions of the crystal and the slides are coupled im such a way that any plane, hkd. can be brought into the reflecting position by setting the coordinates of the corresponding reciprocal-lattice point, he<sup>4</sup>, hb<sup>2</sup>. Act, on -96-

on the three slides. The oscillation motor then retates the crystal through a small arc. During each oscillation, three counts are recorded. The first and third are background counts (time, t) recorded at the two extremes of the oscillation. The middle count (time, 2t) is recorded while the crystal oscillates through a small, pre-set arc. A minimum of two oscillations, one with each of the SrO-ZrO, belanced filtors, is required to measure peak and This type of counting background for a single reflection. mechanism should produce similar standard deviations for MOSt measurements. However, spring and backlash in the mechanisms cause setting errors that increase with decrease in the cylindrical coordinate, 5. This has two consequencess in terms of operation of the diffractometer, it means that reflections from planes with low sin & values must be measured by hand; in terms of quality of the data, it implies that absolute errors in the intensity data will show some systematic dopendence on sin \$-

#### 1.2. REPIMEMENT OF THE STRUCTURE

Brown's (1961, unpublished) coordinates from differential synthesis refinement of two projections were used as starting point for the three-dimensional refinement of 1,1°-binephthyl. The initial R-value was 36°2%, but this was reduced to 18°7% by three cycles of isotropic fullmatrix least squares refinement with Cruickshank's (1961) woighting scheme;

 $\sqrt{v} = (1/p_1 + P_0 + p_2 P^2 + p_3 P^3)^{\frac{1}{2}}.$ 

The values of the parameters wores

p<sub>1</sub> = 4.20, p<sub>2</sub> = 0.011, p<sub>3</sub> = 0. One further isotropic refinement cycle improved B to 14.6%.

The coordinates of the seven hydrogen atoms were calculated by hand. When these were included in the structure-factor calculation, the E-value was 15.8%. One further refinement cycle reduced R to 14.6%. At this point, the parameters of the weighting scheme were changed slightly:

 $p_1 = 2.56, p_2 = 0.011, p_1 = 0.$ 

The refinement continued with isotropic temperature factors for hydrogen atoms and anisotropic temperature factors for carbon atoms. Because of the size of the matrix involved, the block-diagonal approximation was used. After five cycles, the refinement seemed to be near convergence: the N-value had fallen to 9.37%; indicated parameter shifts for carbon atoms were between 1/5 and 1/10 of the corresponding estimated standard deviations; and the weighting analysis showed that the average value of  $2\pi 4^2$  was reasonably compared when analysed as a function of either ain  $\theta/\lambda$  or  $|F_{obs}|$ . However, none of the hydrogen atoms showed unlikely temperature factors and had moved to unacceptable locations.

Stowart, Devideon and Simonon (1965) have pointed out that if hydrogan atom form factors calculated from a simple 1s wave function are used in least-squares refinement. isotropic temperature factors will be abnormally low and the C-H bond distances will be short. The scattering factors used, the present calculations are those tabulated by Hanson, Hormon, Lee and Skillman (1964). Another possible source of difficulty in the refinement of the hydrogen parameters is the choice of weighting scheme. Cruickshank's (1961) veighting scheme was designed for film data and is probably not applicable in this case. 1 East particular parameters used tend to give high relative weights to planes with small intensitios.

The positions of the hydrogen atoms were receleuleted by a computer program. Each atom was placed 1.10 Å along the external bisector of the angle subtended at the corresponding carbon atom. Further refinement was attempted using these calculated positions with a scattering enve appropriate to bonded hydrogen (Stewart et al., 1965) and a new weighting scheme of the form,

 $\sqrt{w} = \left[1 - \exp(2\sin^2\theta / \lambda^2)\right]^{\frac{1}{2}}$ 

The anisotropic refinement of carbon parameters converged in two cycles to an R-value of 9.36%, but the weighting analysis was not satisfactory. A more successful weighting scheme was devised after examination of the

unweighted structure factors of a function of the magnitude of  $\mathbf{F}_{obs}$ , the magnitude of cin 3/A and the index of the recipirocal lattice not, as shown in Table 1. The **R-value** is high for reflections with small values of  $|\mathbf{F}_{obs}|$ . The increase in R with increase in layer index is caused by the large number of weak reflections observed at high angles. The absolute error is the structure factors is constant for most categories of reflections, but increases for large values of  $|\mathbf{F}_{obs}|$  and for low values of sint /A. The weighting scheme adopted after this examination of the data gives unit weights to most reflections, but

if  $\sin^2 \theta / \lambda^2 < 0.16$  then  $w = (\sin^2 \theta / \lambda^2)/0.16$ ; if  $|F_{obs}|$  > 20 then  $w = 20/|F_{obs}|$ . Reflections with  $|F_{calc}| < 1/3 |F_{obs}|$  were considered to be unreliably phased and were given a weight of zero. This effectively removes 45 of the 1315 reflections from the refinement. The value of  $\Sigma w \Delta^2$  was 1545.86.

After two full-matrix calculations involving only the parameters of the carbon atoms, E was  $9\cdot30\%$ . Positions and isotropic temperature factors of hydrogen atoms were allowed to refine for four block-diagonal cycles. Befinements was terminated at this point with  $R = 3\cdot50\%$  and  $5wA^2 = 1097\cdot60$ . Indicated shifts of carbon parameters were less than 3/10of the corresponding estimated standard deviations; shifts of bydrogen parameters were between 1/3 and 1/5 of their 0.5.6.

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6	1150-58	1095-6	194-83	235	0-169	0-63
8	1278.04	1255-3	136.21	183	0.107	0•7.0
20	933-00	905-8	73.73	105	0-079	0.70
12	694-00	699-63	53-15	64	0-077	9-20
14	826-93	836-49	66.63	64	0-030	1 × 0 5
17	1004.84	956069	54-95	65	<b>0-</b> 050	Q.C.
24	1487 - 58	1493-25	81 • 81	74	0-655	$\frac{1}{2} = \frac{1}{2} \langle i \rangle$
33	1850.24	1800.14	$102 \cdot 14$	66	0·056	2150
67	2193+30	2162°35	78•63	50	<b>0</b> -036	197
110	648 • 36	666•89	23.71	\$	0°03₹	2-96
	As function	of sin0/	'λ			
• 20	845-15	829•49	50.72	27	0-060	3,200
• <u>30</u>	2190.96	2149.28	125.62	93	0°052	
• 4 <b>(</b> )	2399•49	2270.33	204.60	284	<b>0</b> • 085	
• 45	1705.99	1706-12	109-37	<u>11</u> 9	<b>0</b> •864	
• 50	1476-63	1520.34	131.63	<u>1</u> 79	0.000	0000
° 55	1241 • 24	1257-40	105 • 52	251	0.002	9×70
°60	1196-40	1157.24	144.68	184	0.133	Q = 75
•63	1061 • 34	1024.54	137.36	181	0-129	3.20
a Z Ç	646•91	58 <b>1.</b> 60	109.53	114	<b>0</b> •269	<b>0</b> *93
• <u>7</u> 5	342• <b>37</b>	303 <i>•4</i> 7	67 • 06	68	<b>0</b> ~396	0 < 0
< <u></u>	246.09	98•1 <b>3</b>	53.70	34	0.368	
2,00	42 • 87	29047	12.36	11	0.367	オッドウ
	An Sumetion	of recip	rocal la	ookoo	inden	
()	1642-53	1636.59	116.78	94	0.013	
÷.	2897-20	2316.72	555°64	313	0.017	2 × 2 3
2	2653•17	2632.49	<b>203 • 3</b> 8	500	0~07S	() 스탠스
<b>?</b>	10.8625	1741.73	155-26	267	0.081	0.53
4	2674-48	1642-12	152.26	180	0.035	0+03
5	114 <b>2 • 88</b>	5106085	137.057	197	0.150	(Des C)
4Ĵ	83 <b>4•94</b>	778.61	152+54	<u>)</u> 449	0-127	See See
	077-23	503+92	146-65	<b>1</b> 50	0-212	

#### 1.3. MOLECULAR DEREGIOUS AND DESCRIPTION OF THE STRUCTURE

Final atomic coordinatos and temperature parameters are listed in Tables 2 and 3, respectively, with standard deviations given as units in the last place. Bond lengths and angles and their standard deviations are shown in Figure 1 and in Tables 4 and 5.

The molecular structure of  $1,1^{\circ}$ -binaphthyl viewed in projection along the g axis is shown in Figure 2. The planar maphthalene residues are linked by a bond across the diad axis, h. The length of this bond  $(1 \cdot 475 \pm 0 \cdot 005 \text{ Å})$ is very close to the value of  $1 \cdot 479 \text{ Å}$  predicted by Devar and Schmeising (1959) and to the value of  $1 \cdot 477 \text{ Å}$  predicted by Pauling's method with a new order-length curve (Cruickshank and Sparks, 1960).

The equation of the best plene through the nephthelene nucleus is

0.0512X + 0.5596X - 0.8272Z = -1.6754where X, X and Z are orthogonal axes defined in the following formulae:

 $X = ax + cscos\beta;$ 

X = by;

 $2 = czsin\beta;$ 

Deviations from this plane are shown in Table 8. A  $\chi^3$ , test with 7 degrees of freeden has a significance level of 0.10, indicating that the maphthelene ring systems are essentially planes. If a mean plane is calculated through either of the beasene rings of the asphthelone system, three of the carbon atoms in the other ring show very large deviations from their plane. This does not imply that the asphthalene system is buckled; it just shows that the "best" planes through the individual rings are slightly different from the best plane through the whole system.

In Tables 5 and 7, some of the more important intraand inter-molecular contacts are shown. Primed numbers refer to atoms whose positions are generated by the symmetry operation of the died axis. It seems likely that the most important intra-molecular forces are carbon-combon repulsions between  $C(2)...C(2)^{\circ}$ ,  $C(8)...C(8)^{\circ}$ , and  $C(8)...O(9)^{\circ}$ whereas, the most important inter-molecular forces arise from the very large number of  $C_{**}$ . H interactions. Each naphthalene residue has 32 non-bonded  $C_{**}$ . H contacts shorter than 3.50 Å. The compromise between inter- and intramolecular forces produces an angle of 68° between the planes of the two ring systems.

Badar of al. (1965) have interpreted the solid state (nujol mull) infra-red spectrum of the low molting form of  $1_{c}1^{c}$ -binaphthyl in terms of a short non-bonded contact between H(17) and C(8<sup>c</sup>). Present calculations show that H(17) has two short intra-molecular contacts, and four inter-molecular contacts, and four

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The extra peak in the spectrum of the low melting form may be caused by C...H interactions but both inter- and intra-molocular contacts are involved. This would seem to cast some doubt on the argument that the high melting form must have a grans- configuration since it lacks this peak in its spectrum. The authors have reported that the high melting form is optically-active. This implies that it cannot crystallise in a controsymmetric space Any crystal modification that reduced the number group. of C...H contacts would probably reduce inter-molecular repulsion forces. If this allowed the angle between the planes to expand, it would decrease the intra-molecular interactions involving  $II(17)_{\circ}$ It is not necessary to postulate a change in molecular configuration to explain the differences in the two spectre. However, the point will remain "not proven" until suitable crystals of the high melting polymorph become available for X-Ray analysis.

#### 1.4. COMPARISON OF THEORY AND EXPERIMENT

Cruickshank and Sparks (1960) have shown that Pauling valence bond theory and Huckel molecular orbital theory give very good agreement with experiment for naphthaleno, anthresens and various alteraent hydrocarbons. They have recommended that a correlation of the form,

$$r = 1.477 - (1.477 - 1.337) - 1.333p = 0.333p + 1$$

be used for valence-bond calculations, and have suggested that a linear variation between m=0.40 at r = 1.46 and m = 0.85 at r = 1.34 would improve the performance of simple molecular orbital theory. In these relationships, r is the bond length, p is the T-bond order calculated by the Pauling superposition method and m is the T-bond order calculated by the molecular orbital method. Tratter (1964) has recently used these order-length correlations to compare theoretical predictions with experimental results for a large number of aromatic molecules.

Table 9 shows m-bond orders for the various bonds in the asymmetric unit of 1,1:-binaphthyl. The mashering used corresponds to the normal convention for asphthalene. The H.M.O. bond orders quoted here have been taken from Coulson and Streitwieser (1955)

Valence bond theory predicts that the bond-lengths in the maphthalene nuclei of 1,1°-binaphthyl will be the same as those in maphthalene itself. For this reason, it is interesting to compute the lengths obtained in the present study with the values obtained by Gruicksheak and Sparks for maphthalene and with those predicted by the two

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theories. Mean lengths are compared in Table 10, and the root mean square differences are shown at the foot of each column.

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A X<sup>2</sup>test with 10 degrees of freedom shows that the dimensions of the mepthalenic residue are not significantly different from those of naphthalene. Hovevez, the differences between observed and theoretical values are highly significant. In both cases, agreement is vorse for the short bonds than it is for the long ones and the major source of disagrooment is the very short C(5)-C(6) It is likely that the discrepancies will be bond. reduced when the experimental values are corrected for vibrational oscillation. Some of the bonds that show vorst agreement are the ones likely to be most affected by this correction.

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V. T-Binaphthyl: Final Fractional Atomic Coordinator. Standard deviations are given as units in the last place

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C(1)	0:03550(10)	0.1057(44)	0.25298(25)
c(2)	0:06095(12)	~0.05065(44)	0.18985(29)
C(3)	0.12841(14)	~0,05107(48)	0.19519(31)
c(4)	0.17047(12)	0.08794(48)	0.26514(29)
C(5)	0:18919(11)	0.41283(47)	0.10531(29)
C(6)	0.16698(13)	0.55895(47)	0.47141(30)
C(7)	0.09985(13)	0.57849(44)	0.46830(29)
c(8)	0.05754(10)	0.43030(40)	0.39799(25)
C(9)	0.07905(10)	0.26431(38)	0.39733(23)
C(10)	0 14708(10)	0.25500(41)	0.33171(23)

Final coordinates and isotropic togerature parameters for the hydrogen atoms.

H(11)	0:0301(15)	-0:1522(49)	0:1337(30)	0.0385
H(12)	0:1443(17)	-0:1573(55)	0:1494(35)	0.0517
H(13)	0:2179(18)	0:0827(53)	0:2679(36)	0.0579
H(14)	0:2320(18)	0:2007(60)	0:4073(37)	0.0590
H(15)	0:1965(17)	0:6831(55)	0:5195(36)	0.0590
H(15)	0.1965(17)	0.6831(55)	0.5195(36)	0-0534
H(16)	0.0827(16)	0.6920(53)	0.5167(33)	0-0485
H(17)	0.0152(16)	0.4328(57)	0.3985(3~)	0-0516

#### Table 3

Alloctropic temperature parameters of the called acord.

	U	U	U	2023	2931	20 12
G(1)	0.0301	0.0360	0.0370	0.0071	0.0158	0.0051
C(S)	0.0379	0.0432	0.0501	-0.0071	0.0217	0.0114
2(3)	0.0435	0.0526	0.0542	~0.0078	0.0321	0.0325
C(4)	0-0262	0.0566	0-0492	0.0180	0.0149	0.00年9
C(6)	0.0359	0.0501	0.0529	-0.0042	6310.0	~0.0195
C(7)	0:0380	0.0461	0.0457	-0.0115	0-0185	~0.0089
C(8)	0.0270	0.0435	0-0407	~0°00255	0.0241	-0.0013
C(9)	0.0278	0.0395	0.0332	0~0093	0.0158	0-0020
C(10)	0-0580	0.0451	0-0370	0.0174	0.0.85	0.0062





Figure 2



Table 4

# 1, 1-Binaphthyl: Bond Longths.

Standard deviations are given as

valto C(1) C(1) C(2) C(3) C(3) C(4) C(5) C(5) C(5) C(7)	in the last C(1) C(2) C(3) C(4) C(5) C(5) C(6) C(7) C(8)	place, 1.4748(46) 1.3645(37) 1.4036(38) 1.3598(41) 1.4127(39) 1.4135(37) 1.3450(42) 1.4037(38) 1.3590(37)
C(9) C(9) Bonds C(2)	C(9) C(10) C(1) involving hy H(11)	1.4120(35) 1.4162(29) 1.4328(32) /drogen. 0.983(31) Å
	H(12) H(13) H(14) H(15) H(15) H(17)	0,929(35) 0,989(37) 0,904(37) 0,995(36) 0,957(34) 0,889(34)

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### Table 5

1,1~B	inaphti	yl:	Band	Anglee.(°	)
C(1) C(2) C(2) C(2) C(2) C(2) C(2) C(2) C(2	CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC	00000000000000000000000000000000000000	)	122.01 119.84 120.79 119.33 122.43 122.30 119.81 119.66 122.11 119.66 123.07 119.04 119.04	

# Bond angles involving hydrogen.

C	(1)	C(E	e) H(11)	118.1(1.8)
H	(11)	C(a	e) C(3)	119.7(1:8)
C	(2)	<b>C</b> (3	) h(12)	120.3(2.1)
H(	(12)	C	) C(4)	119:8(2:1)
C	(3)	C(L	) H(13)	120.2(2.2)
H(	(13)	C(L	I) C(10)	119.0(2.2)
C	(10)	C	5) H(12)	114:8(2:4)
$\mathbf{H}($	(14)	C	5) C(6)	122.8(2.4)
C	(5)	C	5) H(15)	121.9(2.1)
Н	(15)	C((6	C(7)	118.2(2.1)
C	(6)	C((7	(16) H(16)	120.5(1.9)
H(	(16)	C((7	╯) C((8))	119.8(1.9)
C	(7)	C((8	3) H(17)	120.2(2.3)
H	(17)	C({	3) C((9)	117.7(2.3)

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# Table 6

1, 1-Binaphthyl: Intra-molecular contacts (3.50 A.

:

C	(1)	C(2)	2,456
C	(1)	C(9)	2.534
CÍ	(1)	H(11)	2.595
C	(1)	H(17)	2.631
C	(1)	<b>C</b> (8)	2.971
C	(2)	H(11)	3.008
C	(2)	C(2)	3.107
C	(2)	H(17)	3.470
C	(8)	H(17)	2.999
CÌ	(8)	C((8))	3.321
C	(8)	-C((9))	3.333
C	(9)	H(17)	2.813
C	(9)	C((9)	3.277
H	(9)	<b>H</b> ((11)	2.951
H(	(17)	H((17)	2.905

#### Table 7

Inter-molecular contacts. Equivalent position numbers refer to the following transformations.

10	KoV DED	5.	1/2+x,	1/2+y, z
2.	-X. Y. 1/2-2	6.	1/2-20	1/2+y, 1/2-z
3.	- X, - V, - Z	7.	1/2-20	1/2-90-2
Ц,	X, -y, -1/2+Z	8.	1/2+20	1/2-y, -1/2+2

C-H	contacta	\$ 3.15	A.		
C(1)	) H(16)	Ŭ,		(0,1,0)	3.100
C(2)	) <b>H(16</b> )	29,		$(0_{\theta}1_{\theta}0)$	2.983
C(3)	) H(16)	Ly		(0,1,0)	2.962
C(4)	<b>H</b> (16)	Lţ.		(O, 1, O)	3.040
C(4)	) <b>H</b> (15)	Ľ <u>y</u>		(OaloO)	3.051
C(5)	) H(12)	7		(0,0,1)	2.955
C(5)	) H(13)	6		(0 <sub>0</sub> 0 <sub>0</sub> 0)	3.132
C(7)	) H(17)	3		00000	5.05/
C(8)	) <b>H</b> (17)	3		(Oolel)	20999
C(8)	H(11)	Сў. Ро		0,0,1	5.140
C(10	)) N(16)	E.J.			3.140
<b>A A</b>		1 9 En	A		
	contacts	< 3.00 1	ALC.	(a - 9 a)	3.600
C Z	G	ะ		0.51.0	3,580
22		9 7		0.0.1	3.450
282	63	<b>1</b>		0.1.0	3.580
c 7		จ		0.1.0	3.600
~~~	161	u u		<b>N</b> -0 <b>N</b> -0	-
Н-н	contacta	1 2,50	Α.		
H(17	H(17)	3		(0,1,1)	2.46
· · · · · · · · · · · · · · · · · · ·	0				

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Table 8

1,1-Binophthyl: Moan Planes.

Deviations from the planes are shown in whits of angstroms  $\times 10^{-4}$ .

- A. Plane through naphthalone residue: 0.0512X + 0.5596Y = 0.82722 = -1.6754
- B. Plane through atoms 1,2,3,4,9,10: 0.0567X + 0.5544Y = 0.8303Z = -1.6774
- C. Plane through atoms 5,6,7,8,9,10: 0.0451X + 0.5653Y = 0.8237Z = -1.6536

		A	В	C
C	(1)	99	10	3
C	(2)	-5	14	-222
C	(3)	-97	1	-400
C	(4)	-105	-40	-371
C	(5)	197	32	16
C	(6)	-6	-197	12
<b>C</b> (	(7)	-127	396	-23
<b>C</b> (	(8)	-62	299	6
C	(9)	77	-48	22
<b>C</b> (	(9Ó)	110	63	-32

Root mean square deviation from planes,



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# Table 9

1, 1-Binaphthyl: T-Bend Orders,

Normal naphthalons numbering is used.

÷		V.B.	11. M.D.
C(1)	C(2)	0.667	0.659
C(S)	C(3)	0.333	0.621
C(3)	C(4)	0.667	0.712
C(4)	<b>C</b> (10)	0.333	0.550
C(10)	C(5)	0:333	0.554
C(5)	C(6)	0:333	0.722
C(6)	C(7)	0.333	0.607
C(7)	<b>C</b> (8)	0:667	0.717
C(8)	C(9)	0:333	0.559
C(9)	C(10)	0:333	0.523
<b>C</b> (1)	<b>C</b> (9)	0.333	0.507
C(1)	C(1)'	0.000	0.404

#### Table 10

### A comparison of hond lengths.

	<u>Bingelik</u> yl	naphihalanc	<u>V.B.</u>	H.M.O.
C(1) C(2) C(2) C(3) C(3) C(4) C(10) C(5) C(10) C(5) C(5) C(6) C(6) C(7) C(7) C(8) C(7) C(8) C(9) C(1) C(9) C(1) C(1) C(1) Root mean a	1.264 1.404 1.360 1.413 1.413 1.413 1.404 1.359 1.404 1.359 1.416 1.433 1.475	1.364 1.415 1.364 1.421 1.421 1.421 1.421 1.421 1.421	1.375 1.421 1.321 1.421 1.425 1.425 1.421 1.421 1.421 1.427	1.391 1.401 1.377 1.420 1.419 1.419 1.475 1.475 1.459 1.459
		0.009	0.014	0.015

			_114	1_			
	1 1_Bir	anhth	$v_1 S + m$	, leture	Factor	Tablos	• C+a
икт	I Fol Fo	нкт	Fol Re	HKL			
18 2 3	10.7 -11.3	9 3 8 11 3 7	4.9 4.2	7 3 1	28.4 27.4 27.3 -26.3	-22 4 4 2.6	2.0
H 8225864 x 5 24468 22 24 28 28 28 28 28 28 28 28 28 28 28 28 28	1.       Pro       3.1.4.2.0.48.84.5.0.7.7.9.5.6.4.7.6.3.98.0.0.0.3.7.9.2.5.7.6.8.1.5.8.0.7.6.8.1.0.0.4.0.4.8.9.0.3.8.0.0.1.7.5.0.8.7.6.8.7.6.8.1.0.0.3.7.2.0.9.0.3.4.0.0.1.8.8.4.7.4.6.0.7.8.6.1.0.4.0.4.8.9.0.3.8.0.0.1.7.5.0.7.6.8.1.0.0.1.8.8.4.7.4.6.0.7.8.6.1.0.4.0.4.8.9.0.3.8.0.0.1.8.8.4.7.7.4.6.0.5.4.8.0.0.1.8.8.4.7.7.4.6.0.5.4.8.0.0.1.8.8.4.7.7.4.6.0.5.4.8.0.0.1.8.8.4.7.7.4.6.0.5.4.8.4.8.4.6.2.4.8.9.2.3.8.0.7.6.4.0.4.8.9.2.3.8.0.7.7.4.6.0.7.8.6.1.0.0.3.7.8.7.8.1.3.4.8.0.0.1.8.8.4.7.7.4.6.0.5.4.8.0.0.1.8.8.4.7.7.4.6.0.7.8.0.1.3.0.7.8.6.7.8.6.1.0.0.3.8.7.8.7.8.1.3.4.8.0.0.1.8.8.4.7.7.4.6.0.5.4.8.0.0.7.8.0.1.7.7.4.6.0.7.8.0.1.0.0.3.7.8.7.8.1.7.5.8.7.8.0.1.0.0.1.8.8.4.7.7.4.6.0.7.8.0.1.0.0.1.8.8.4.7.7.4.6.0.7.8.0.1.0.0.1.8.8.7.8.1.1.8.7.3.2.1.2.2.3.8.7.8.1.7.5.8.7.8.8.1.7.6.7.8.6.1.0.8.8.7.8.7.8.1.1.8.7.3.2.1.2.2.3.8.7.8.1.7.5.8.7.8.8.1.7.6.8.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.6.7.8.8.7.8.7		Fe 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-9 5 2 -7 5 2 -1 5 2	5.0 -4.9 4.4 -4.8 16.3 17.0	-24 6 4 -22 6 4 -20 6 4	3424	-3.1 4.4 7.6	15 7 8 15 7 7 15 7 7	4.5 2.8 6.8	4.7 3.3 7.5	-23 7 1 -25 7 1 21 7 0	3.0 3.7 3.6	0.4 -1.6 3.5
1 5 2 -3 5 2 3 5 2	31.4 32.6 8.8 -8.7 6.5 -6.2 4.6 4 4	-10 6 4 -12 6 4 -8 6 4	2.7 4.1 5.6	1.9 4.4 -6.3	977 577 377	3.8 4.7 3.7	-2.1	19 7 0 17 7 0 15 7 0 13 7 0	3.2 3.7 2.6	-1.6
7 5 2 9 5 2 11 5 2	4.4 4.3 11.2 -11.9 8.0 7.1	-6 6 4 -4 6 4 -2 6 4	2.5	-2.1 3.1 3.2	$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	2.3	2.2	11 7 0 9 7 0 7 7 0	7.2 4.2 6.7	7.2 -3.6 -6.6
13 5 2 15 5 2 23 5 2	$\begin{array}{ccc} 8.3 & -7.5 \\ 21.2 & -21.1 \\ 7.4 & 7.1 \end{array}$	264 664	4.0 4.0 7.3	-4.2 -7.2	-11 7 7 -13 7 7	4.4 7.9	-3.4 6.6	570	1.4	-1.4

1,1'-Binaphthyl Structure Factor Tables Ctd.

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#### THE CENSTAL AND MOLECULAR STRUCTURE OF PERTLEME

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In 1963. Tanaka discovered a second crystalline form of perylene which he designated as the S-form. From studies of absorption and flucroscence spectra of the two polymorphs, he concluded that the previously known «-form is a ground-state-excited-state dimer. To support his arguments, he dotormined the crystal structure of  $\beta$ -pozyleno in two projections. The hOA projection is well resolved and was refined to an E-value of 17%. The y-coordinates were derived by assuming that the molecule is planar and has a structure similar to the Because there is serious overlap in the Ok& a-form. projection, the y-parameters were not refined. The 2-value for this zone was 27%. In view of the current interest in molecules containing bonds with extreme mebond orders, it seemed worth-while to extend Tanzka's study to three dimensions.

The crystal structure of G-perylene was first established from projections along two axes by Donaldson, Robertson and White in 1953. Their measurement of 1.5040.03 Å for the length of the peri bonds was in very good agreement with Coulson's (1951) theoretical prediction of 1.50 Å. Camerman and Trotter (1964) have made a threedimensional study of G-perylene using counter data. They report the value of 1.47140.006 Å for the longth of the long bonds and state that the molecule is significantly en en gren El region gren e

distorted from planatity.

There are no other debailed analyses of pervious itself in the literature, but the molecule forms denoracceptor complexes with a wide variety of acceptor molecules including fluoranil and pyromellitic disubydride. The crystal structures of the T-molecular complexes of perviews with these two compounds have been determined (Kanson, 2063; Budyeas and Merbotein, 1965).

The two polynopphs exhibit fluorescence when strained under ultra-violet light. This provides an easy coshed for distinguishing between them, since the fluorescence from the a-form is red and that from the f-form is prov-

#### 2.1. ENPERIMENTAL

There were a few exystels of the  $\beta$ -form is the couple of analytical grade perplene obtained from Kosh-Marko Laboratories. Since these were not suitable for single crystal studies, they were set aside to be used as seed crystals. When perticus of the sample were possibilities. from benzene and from petroleum other, the first "erep" consisted mainly of the 6-form. If the supermaters liquid was decanted and seeded with crystels of the  $\beta$ -form it yielded a good erep of very this hemagenal plates that were to brittle to be cat. However, it was possible to grow very large crystels by allouing bonness solution; to evaporate slowly for several menths. Crystels grow in this menuer were out to the classe of a uniform well were used to obtain cell dimensions and intensity date.

land, en Samerika (Se

#### 2.1.1. CRYSTAL DATA

Formula	<sup>C</sup> 20 <sup>II</sup> 12	F.U.	252+3
System	monoclinic		
Lattice constants	$a = 9.77 \overset{\circ}{\Lambda}$		
	$b = 5 \cdot 90_{A}$		
	c ≈10•58 <sub>9</sub>		
	β <b>≈96∙75°</b>		
Absent Spectra	hOL if & odd		
	0k0 if k odd		
Space group	P21/c		
Volues	607 • 0 Å <sup>3</sup>		
Bensity(obs.)	1•38 gm/cc		
Molecules/coll	2		
Density(cale.)	1.380 gn/cc		
Linser absorption c	officient µ, (	)∘920 cs	,- <b>1</b>
(No Ra; = 0	•7107 Å)		

Since there are two molecules per unit cell, the centre of the molecule must lie on a centre of inversion.

The cell dimensions listed here have been measured from precession and Veissenberg photographs and have been confirmed by checking reciprocal lattice spacings of high angle reflections on the diffractometer. The length of the g-axis is considerably shorter than Tanaka's value -119--

of 11-27 Å but the gattern of intensities is in agreement with the reflection data quoted in his paper.

#### 2.1.2. REFLECTION DATA

Reflection data were obtained on a Hilger-Vatta linear diffractometer using No Ka radiation. The diffractometer is equipped with Sr0-Zr0, balanced filters Two crystals were used. and a scintillation counter. Data from reciprocal lattice nets hOf to h44 were collected on one exystel and date from nets h5% to h7% vare collocted on the second. After Lorentz-polarisation corrections had been applied, the scale of the observed structure amplitudes was approximately twice the absolute All date wore included in the refinement. SF3 20 scale are no "unobsorvod" planes because reflections with negative not counts are given an intensity of +1 by the data processing progress.

#### 2.2. REFINEMENT OF THE STEUCTURE

Although the coll discussions dotermined in the present study vero substantially different from these measured by Tanaka, the intensities of the reflections in the hOS zone agreed with the values quoted in Tanaka's paper. It seemed reasonable to use his coordinates as the stating point for the reflacment of the structure.

When the "minimum residuel" method of Bhuiya and

Shealoy (1963) was opplied to the maigne and project of the R-value fell from 3%5 to 16% in eight opeles. Profess m- and p-coordinates from this projection wast coefficient with Ramaka's calculated y-coordinates and used to be your structure factors for the five reciprocal lattice costs, hOS to h45. The initial R-value was 54%, but this was reduced to 21.5% after three cycles of Fourier reflected After termination-of-series corrections had been equal of to the atomic parameters, refinement was continent by 20. method of least squares. Three isotropic least-space for calculations with unit weights improved the R-walue 10 18.9%. At this points, a weighting function of the form

 $w = 1 - \exp(-2 \sin^2 6/\lambda^2)_{\ell}$ 

was applied. R was reduced to 15.0% by two Surthest Acoustance calculations but was improved by only 0.5% when the temperature factors of the carbon atoms were allowed to refine anisotropically. When bydrogen atoms were thethet in the structure-factor calculations, E foll to 13.6% and dropped to 12.7% after two cycles of emissively testinemet.

Shie from layers h54 to h78 were included at this point. Three full-matrix least-squares cycles were colectated to roline the scale factor between the two batches of data. In order to evold a singularity of the matrix, the V<sub>pp</sub> temperature parameters were bold control. The R-value was 16.4% for 1805 reflections.

.123.

The unweighted structure factors were analysed as a function of sin  $\theta/\lambda$ , the index of the reciprocal lattice layer and the magnitude of  $F_{obs}$  as shown in Table 11. The weighting scheme devised from this examination of the data had the following form:

> if  $\sin^2\theta/\lambda^2 \leqslant 0.16$ , then  $w = (\sin^2\theta/2)/0.16$ ; if  $|F_{obs}| \ge 7.0$ , then  $w = 7/|P_{obs}|$ .

Planes with  $|F_{calc}| \leq 1/3$   $|F_{obs}|$  were considered to be unreliably phased and use given a weight of sore.

A separate analysis of the weak reflections showed that R was over 200% for reflections with  $|\mathbf{F}_{obs}| < 0.9$ electrons and over 80% for reflections with 0.5  $\leq |F_{obs}| \leq 1.0$ . Most of the reflections with structure amplitudes such as then one electron wore among these that had been given an intonsity of one count per minute by the data processing program because the not count was negative. It is therefore surprising to find that |Feale | is larger than |Fobs This may indicate that there for nearly all of them. is some systematic error in the measurement of the weak In order to investigate this point it will reflections. be necessary to alter the programs and to process the data again. As a tomporary expedient the 488 planes with |Pobs|<1.0 were caluded from the refinement. Qf abe remaining 1317 planes, 126 had been given a weight of sovo. The R-value was 13.8% and Two? was 498.31. After

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# 2001c 11

 $\beta$  -Forylono: Analysis of Observed and Calculated Structure Amplitudes for Rollections with F 24 0

As a function of |Pobel

.80 1.00

	S Pol	E Pd	5141	N	B	Δ	
2	930-90	823-96	196-50	933	0.533	0.53	-
3	643.56	547049	172-37	262	0.268	0.66	
4	438-46	461.19	80-93	141	0°166	0°57	
8	1548-57	1534.66	153-33	273	0.099	0-56	
16	1352-16	1297.64	97-74	123	0.072	6.79	
33	1386-67	1355-39	73.08	64	0-053	1.14	
111	381-81	391-05	26.56	9	0.040	2°95	
	As a functi	ion of sin	0/ <b>λ</b>				
-20	491 • 51	494-32	51.01	32	0.104	2.59	
• 40	2170.07	2089-89	234.40	279	0-108	0.84	
50	1469.36	1439.22	150-42	298	0.105	0.50	
° 60	1169.82	1133-33	196-38	441	0°168	0.44	
-70	962.06	857.18	262-68	490	0-272	0.54	
°80	406-60	286-39	173-66	<b>5</b> 58	0-425	0.76	
l • 00	57°71	31.05	32.58	37	0°564	0.88	

As a function of reciprocal lattice index

O	939-05	930-39	99°06	140	0-105	0-71
L	1512-15	1489.49	163-62	260	0.108	0-63
2	1393-97	1351-83	165 • 52	262	0.119	0°63
Э	782.22	740-37	131-03	254	0°158	0.25
4	777.81	752-31	125-58	237	0°162	0.53
5	618-68	595 • 58	140-14	240	0.556	0.28
6	356 • 25	282 • 97	129.42	208	0•363	0.05
7	352-00	268 • 44	146.14	204	0.415	0°73

# three further cycles of least squares, R was 13.3%, WA<sup>2</sup> was 460.56 and the refinement of the carbon stone had converged.

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#### 2.3. INTERIM RESULTS

Atomic coordinates are listed in Table 12 with their standard deviations. Temperature parameters are shown in Table 13. The hydrogen atoms were placed by a computer program and have not been refined; their temperature factors are probably slightly low.

Table 14 is a list of bond lengths and bond angles derived from the parameters of the carbon atoms. All C-H bonds are 1.09 Å leng.

The porylone molecule is remarkably planar as is shown in Table 15. The equation of the "best" plane through the whole molecule is

0.2633X - 0.5793X - 0.7714Z = 0, where X,X and Z are orthogonal axes defined by the equations

 $I = gx + cscos\beta$ ,

 $\mathbf{X} = \mathbf{b}\mathbf{y},$ 

$$Z = ezsin6$$
.

We not not equare deviation from this plane  $(3 \cdot 4 \pm 10^{-3} \text{ Å})$ is large than the avorage optimated standard deviation of aborto position. A more severe test of planatity is provided by the sugations of the deviations from a plane

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#### Table 12

6 -Perylene: Final Fractional Atomic Coordinatos. Standard deviations are given as units in the last place.

CCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCCC	0.05911(30) 0.11785(36) 0.26034(33) 0.34271(31) 0.37200(36) 0.31745(31) 0.17620(30) 0.08948(35) 0.14412(37)	-0.16889(61) -0.32806(75) -0.32329(70) -0.16218(67) 0.17590(68) 0.33318(56) 0.33403(57) 0.17241(68)	0.08899(27) 0.17450(35) 0.21807(34) 0.17535(29) 0.04415(33) -0.03988(27) -0.08412(28) -0.08412(28) -0.04359(34)
C(8)	0:08948(35)	0:17241(68)	-0。04359(34)
C(9)	0:14413(37)	0:00226(78)	0。04423(37)
C(10	) 0:28757(39)	0:00439(76)	0。08855(37)

Calculated poltions of hydrogen atoms (not refined).

H(11) 0.0531	-0.4611	0.2082
H(12) 0:3032	-0: 4486	0.2876
H(13) 0:4531	0.1654	0.2080
H(14) 0:4055	0.0088	0.0767
H(15) 0:3844	0.4609	-0.0748
H(16) 0.1347	0.4640	-0.1514

#### Table 13

.

Anisotropic Temperature Parameters of Carbon Atoms.

	Um	Uaa	U 33	EUss	2U31	2U12
C(1)	0.0307	0.0377	0.0312	0.0001	0.0104	0.0020
C(S)	0.0353	0-0551	0.0452	0.0027	-0.0001	0.0145
C(3)	0.0287	0.0520	0.0439	-0.0123	0.0072	-0.0119
C(4)	0.0294	0.0449	0.0347	-0.0181	0.0079	0.0021
C(5)	0.0405	0.0427	0.0438	-0.0021	0.0084	-0.0237
c(6)	0.0328	0.0338	0.0310	-0.0017	0.0115	0°0002
C(7)	0.0315	0.0341	0.0333	~0.0095	0.0087	ତ <u>୍</u> ଟ୍ରେପ୍ଟ୍ରେ
C(8)	0.0363	0.0437	0.0480	0.0077	0.0117	0-0118
C(9)	0.0352	0.0566	0.0519	0.0055	0.0105	-0.0344
C(10)	0.0436	0.0507	0.0495	0.0271	0.0068	0.0234

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Table 14

 $\beta$  -Perglene: Bend Longths. (?)

Standard deviations are given as units in the last place.

C(1) C(2) C(2) C(2) C(2) C(2) C(2) C(2) C(2	C(2) C(2) C(2) C(2) C(2) C(2) C(2) C(2)	1.3822(54) 1.4154(50) 1.4154(50) 1.3574(53) 1.4093(58) 1.4209(59) 1.3516(51) 1.4052(44) 1.3781(51) 1.4258(50)	
C(9) C(9) <u>peri</u> -l	C(10) C(1) Donde	1.4256(56) 1.4245(56) 1.4758(52)	 

C-H bonds are all 1.09 Å long.

Bond-Angles with their standard deviations.

C(1)	C(S)	C(3)	121.13(37)
C(S)	C(3)	$\mathbf{c}(4)$	120.35(37)
C(3)	C(4)	C(10)	120.64(32)
C(4)	C(10)	C(5)	121.37(35)
C(4)	C(10)	C(9)	119.65(38)
C(9)	C(10)	C(5)	118.99(39)
C(10)	C(5)	C(6)	120.52(34)
C(5)	C(6)	C(7)	121.21(32)
C(6)	C(7)	<b>C</b> (8)	120.68(31)
<b>C</b> (7)	C(8)	C(9)	119.52(33)
C(8)	C(9)	C(10)	119.08(38)
C(8)	C(9)	C(1)	121.80(33)
C(10)	C(9)	C(1)	119.12(38)
C(9)	C(1)	C(2)	119.11(31)

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Pelle 5

B -Perylone: Mean Planes.

Deviations are shown in Ångetross × 10-4

Plane A is the plane through the whole carbon skeleton.

Plane B is the plane through the central ring.

		A	В
C(	1)	29	8
C(	2)	c3 1	~72
C(	3)	-58	-141
C	4)	46	~1Å
C(	5)	-45	-45
C	6)	35	70
C	7}	-16	31
C	8)	~32	-8
C	9)	20	8
c.l	101	90	

Root mean aquare deviation: 34



through the control ving of the solecule. Mose of the atoms Cavintos significantly from this plane, which is defined by the equation.

0.2630X - 0.5773X - 0.7730Z = 0.

The crystal structure of 3-perylene is analogues to those of maphthalene and anthracene (Abrahans, Reburies & White, 1949 ; Sinclair, Robertson & Mathieson, 1950). The orthogonal projections shown in Figures 2 and 3 illustrate this similarity. Figure 2 is a projection of the whole anit coll on to the (100) plane. Figure 3 shows half of the projection on to (010).

It is interesting to compare the results of the analysis with theoretical predictions and with other organic studies. Sable 16 lists the back derivation in β-porylone with these predicted by various theoretical derivation fable 17 summarises the results of X-Day studies derivation is porylone; and Table 13 compares the dimensions of β-porylone with these of quaterrylone, results and have a dimensions of the dimen

The H.N.O. hand-lengths in Table 16 have been obtained by combining the R-band orders tabulated by Stociations and Example (1965) with the order length of suggested by Gaulekeherik and Syachs (1960). Stof harets producted by the Perior Jame-Popte (N.P.R.) worked and the cylit 7 orbited (S.P.O.) method tere behave the s gauge by Doves and Eleicher (1965). Read Mark (1999)







differences between experimental and theoretical languist are shown at the feat of each column. The goodness-of-fit becomes worse as the theory becauss more elegant. The length of the <u>peri</u>-bonds is the main source of disagreement but even when this bond is not considered, the same general trend is observed. The relative w-bond orders predicted by molecular orbital theory follow the observed variations in bond length very closely.

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The more elegant F.P.F. and S.P.O. calculations predict that the control bond in the maphthalonic residue should be shorter than the two outer bonds but the reverse situation is observed in the experiment. The 9-10 bond in  $\beta$ -perylene is longer than either the 2-3 bond or the 6-7 bond. Although these differences are not large enough to be significant, they are probably real slave they are observed in 1,1°-binaphthyl and in the complement of perylene listed in Table 17.

The dimensions of the pervious molecule obtained from various three-dimensional X-Ray studies are shown in Table 17. The average standard deviation of band length appears at the top of each column; the rest seas square difference is shown at the bottom. The parylene molecule is necessarily controlynmotric in  $\beta$ -purvious, in the pervious, in the pervious, in the pervious (Hancon, 1963), and in the pervious (Hancon, 1963), and in the pervious (Hancon, 1963). -129-

B -Ferylene: Bond Langth Comparisons

# <u>Table i6</u>

	ß	-perykue	V.B.	<u>H.M.O.</u>	<u>P.P.P</u> .	<u>S.P.D.</u>
C(1) C(2) C(3) C(4) C(9) C(9) Peri-1	C(2) C(3) C(4) C(10) C(10) C(1) Ponds	1.380 1.410 1.355 1.415 1.426 1.427 1.427 1.427	1 . 421 1 . 421 1 . 421 1 . 421 1 . 421 1 . 477	395 398 378 427 427 4457	1.381 1.413 1.375 1.421 1.404 1.422 1.456	1.379 1.379 1.34 34 375 1.462
Root s	ucan s(	juare á	lfferenc	e: <u>peri</u>	-bonds 1	ncluded.
Root 1	182N 30	juare d	0.010 Lfferenc 0.010	0.013 8: <u>per1-</u> 0.011	0.014 bonds ex 0.012	0.015 cluded. 0.013

Table 17

-

		SPADY SPADY					
		B-scoylane	peryene- fluoranil	Pergiese PMD®	a - persience		
		ଙ <b>=0</b> ₀005	G=0°002	<i>∽=</i> 0.010	0°=0.011		
C(i)	C(2)	1.382	1.380	1.393	1.381		
c(s)	C(3)	1.415	1.419	1.407	i .40i		
C(3)	c(4)	1.357	1.358	1.364	1.360		
c(4)	C(10)	1.409	1.405	1.365	1.412		
C(10)	C(5)	1.421	1,415	1.445	1.395		
C(5)	C(6)	1.352	1.364	1.356	1.362		
C(6)	C(7)	1.405	1.401	1.383	1.405		
C(7)	C(8)	1.378	1.379	1.381	1.384		
C(8)	C(9)	1.429	1.427	1.407	1.427		
C(9)	C(ió)	1.425	1.434	1.443	1.434		
C(9)	C(1)	1.424	1.431	1.423	1.420		
<u>peri</u> -b	onds	1.476	1.473	1.493	1.473		
Root meen sourre difference:							
		6 · · · ·	0.005	0.019	0.009		

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# <u>Table 18</u>

	B-perylene	Susterniene	naphhalene	nt braphthyl
C(1) C(2)	1.380	1.384	1.364	1.343
C(2) $C(3)C(3)$ $C(4)$	1.410	1.425	1.354	1.355
C(4) C(10)	) 1.415	1,392	1.421	1.413
C(9) = C(10) C(9) = C(1)	1.420	1.431	1.418 1.421	1.430
<u>perl</u> -bonds	1.476	1.527		1.475
Root mean s	quare differ	ence includi	ng peri-bond	5:
Hood moon	-	0.025	0.009	0.008
noot meen e	yware allier	0.015	0.009	800. <b>0</b>

bond longths listed for a-perplane are the results of Mason's (private communication) refinement of the date published by Camerana and Frotter. It should be noted that the average standard deviation reported here is approximately twice the estimate quoted in the original paper, despite a much lower R-value.

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In Table 18,  $\beta$ -perylene is compared with quaterrylene (Shrivastava and Speakman, 1960), naphthalene (Cruicksbank and Sparks, 1960), and 1,1°binaphthyl. The root nean square differences between the dimensions of  $\beta$ -perylene and those of the other melecules are small except in the case of quaterrylene. This is likely to be a reflection of the fact that the structure of quaterrylene was determined from two projections; it does not necessarily imply a real difference in melecular structure.

#### DISCUSSION

4

Although the analyzon of these structures are not yot complete, the precision of the measured bond-lengths (0.003) to 0.005  $\stackrel{\circ}{A}$ ) is sufficient to allow a cortain amount of discussion.

Cruickshank (1956) has shown that when libration corrections are neglected, the apparent band lengths are systematically shorter than their true values. In 1,1'binaphthyl, the thermal vibrations are markedly enterthouse and libration corrections have not yet been applied. The bond that is likely to be must affected by whose corrections -131-

is the 5-6 bond. In the pervious structure, libration corrections will probably be very small since there is little anistropy in the thermal parameters. In the present state of the analysis, the molecule appears to have its maximum vibration amplitude almost perpendicular to the plane of the ring. This unlikely situation can be explained in terms of a poor scale factor between the two batches of data.

The dimensions of 1,1'-binaphthyl and \$-perylemo are very similar. The root mean square difference is only 0.008 when "chemically equivalent" bonds are averaged. If the individual bond lengths are compared, the root mean square difference increases to 0.011. However, none of the differences between corresponding bond-lengths is significant.

The lengths of the  $\operatorname{sp}_2 \cdot \operatorname{sp}_2$  bonds reported in this study are  $1 \cdot 475 \pm 0 \cdot 005$  Å in  $1_0 1'$ -binsphtbyl and  $1 \cdot 476 \pm 0 \cdot 005$  Å in  $\beta$ -perylene. These lengths are in good agreement with the value of  $1 \cdot 473 \pm 0 \cdot 005$  Å in the perylene-flueranil complex and the value of  $1 \cdot 473 \pm 0 \cdot 010$  Å in  $\alpha$ -perylene (Mason's refinament) but they are in disagreement with the values of  $1 \cdot 53$  Å in quaterrylene,  $1 \cdot 493 \pm 0 \cdot 010$  Å in the perylene-PMDA complex, and  $1 \cdot 493 \pm 0 \cdot 015$  Å from electron diffraction measurements of perylene in the value phase (Tractaborg, 1965). -332-

The peri-head longth quoted for quoternylane is the avorage value of gin independent head longths and would be expected to be very accurate in spite of the fact that the structure was determined from two projections. The discropancy between this value and the 1.48 A single bond length prodicted by Dever and Schneising (1959) has been attributed to a stretching of the linking bonds caused by storic repulsions between the hydrogen etems However, recent studies suggest in the <u>ortho</u>-positions. that hydrogen contacts of this type are "softer" then Bailey and Mason (1963) have reported had been supposed. N...N contacts of 1.85 Å in 1,8-bisdebydro-[14]-onum.coc. Hirshfeld, Rabinovich and Brogman (1965) have found that the mean separation of the inner hydrogens in [16]-anralese 13 1.97 A. There is no apparent deformation of C-C Agreement between the langthe bonds in either structure. of the linking bonds in the two molecules in the present study would seem to confirm that M ... H contacts of this type do not cause significant longthening of G-C bondo.

Differences between the lengths quoted in this work and those observed in the pervices-MMA complex ave; in general, quite large; but certain experimental details in the latter study should be considered. Thermal vibration amplitudes in the complex are large and endestable limited computer capacity made it necessary to refine the parameters of only 6 atoms in each least-squares cycle; the softwament was touchasted before convergence.
It seems likely that the standard deviations have been underestimated.

The electron diffraction study is difficult to Porylone is a very large molecule to be ovaluate. invostigated by this method, and certain simplifying assumptions must be made. Traetteberg has assumed that the naphthalonic residues in perylene have the same molocular dimensions as nephthalono itself and has detormined the lengths of the peri-bonds by subtracting the meloculor intensity function of nephthalene from the radial distribution function for peryleno. Results obtained from a calculation of this type are extremely sensitive to changes in the model. On the other hand, it is possible that there are real differences in the dimensions of the molecule in the vapour phase  $(300^{\circ}C)$ and in the solid.

It is interesting to note that the trend in 7-bond orders predicted by the Huckel molecular orbital method is very close to the trend in the bond longths observed in this study. (See Figure 4). This suggests that the goodness-of-fit will probably improve as more accurate bond-length determinations become available and better order-length curves are devised. The most serious discrepancy between the experimental bond lengths in  $\beta$ -perylene and the predictions of P.P.P. and S.P.C. celculations is the relative 7-bond order of the central

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bond of the neglethelouie residue. According to these two theories, this bond should be shorter than the outer bonds, but the reverse situation is observed in the four structures involving the parylene molecule and in the structure of 1,1°-binephthy1.

The agreement between the dimensions of the two polymorphs of perylene and between perylene and 1.1'binephthyl is better than one would expect from the size of the standard deviations estimated for the individual structure determinations. Mone of the differences between these molecules is significant at the level of precision of the analyses. The dimensions of the significant nephthalenic residues do not differ from these of maps thalene itself.

#### CONCLUSIONS

The results of the present work show that the band lengths in the maphthalenic residues of  $1,1^{\circ}$ -binaphythyl and  $\beta$ -perylene are the same, within the experimental error, as those of maphthalene itself. This can be atorpreted as direct evidence that the aromatic character of the perylene molecule is localised in the maphthalenic residues.

The average value for the sp2-sp2 single bond longths in these two structures is 1.475±0.005 Å.

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		6											
нкг	Fol	Fo H	K L	Fo	Fc	н	K L	Fol	Fc	нкг	Fo	Fc	
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B-Perylene Structure Factor Tables

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			na ataine na							<u> </u>					
H E	K L 2 10	F0	FC	н к	L	Po	Fc h	H	K L	[ Fo[	Fe	H	KL	Fo	fe 7
***************************************	unanananananananananananananananananana	3082-1763623883993467640373959971747396122118946262176214357643220482812325259914423543635721423729931263114441241559795123183309	1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1, 1	<b>~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~~</b>	10 18849426661466110 10 10 10 10 10 10 10 10 10 10 10 10	10,00,00,00,00,00,00,00,00,00,00,00,00,0	126-1281476545575451-15-148898584515767586555465288868185976555455855545585554558568819785688144765597557757757455858585858585858585858585858	<b>&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;&gt;</b>	๚๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛	41 annrango ann 16arta 6nn 160 o an 107600 o ao 1071 reat ao 1016 na bhann an 1011 an 1010 an 1011 an 100 an 1011 an 1010		<b>。。</b>	ия акаатарана и или полнати и или полнания и или пол ми лий войно полнонования или полнания и или полнани или лий войно полнования или полнания и или полнания	<b>50875005997568151704750101776078805748855145799974048168818145477554865000010841956660000</b>	

<u>B-pervlene Structure Factor Tables</u>

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нк	L Po	Fc	HKL	Fo	Fc	HKL	Fo	Fc	H K L	Fe	Fa	
๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛	578901244 9876552452.12214.457778469288876541012457457122457981932817551882553391435122112987453217553212917298765331554332	399301756223609106905926410800794449419494900555399900400000000000000000000000	๛๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚๚		۲۹۹۹ مالیه ۹۹۹ مالیه والا مالا مالا مالا مالاه و ۵۹۹ ماله ماله ماله ماله ماله ماله ماله ماله	๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛	[a+b] a + b] a	๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛๛	╕╻╖╖┪╗┝╖┝╖┙┙╖┙╖┙╖┙╖┙╖┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙┙	то базата и порединии и п В на порединии и порединии и В на порединии и поредини и порединии и порединии и порединии и порединии и порединии и порединии и поредини и поредини и порединии и порединии и порединии и поредини и по	้างนายระเที่ยนการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบาน เป็นการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการระบานการ	

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# 3-perviene Structure Factor Tables

## Appendix I

Comments on Choice of Unit Cell in a Triclinic Lattice

#### CROICE OF TENIS GELL IN A TELOLITHIC SYSTEM

In the triclinic system, the choice of unit cell is not dictated by symmetry considerations; thus there are an infinite member of ways of chosing primitive cells in a triclinic lattice. However, each lattice is uniquely described by the Bravais "reduced cell"; the cell whose edges are the three shortest non-co-planar translations in the lattice. For purposes of compilation, it is describe to describe the lattice in terms of this reduced cell. The two most common algorithms for deriving a reduced cell from any obstrary primitive cell are the Boleunay reduction and the Divichlet reduction. There are cortain yeasons for proferring the latter.

In any parallelepiped, note of angles eccur in onautiomorphons pairs at opposite corners. One pair of corners has the angles  $\alpha_{,\beta,\gamma}$  homogeneous; that is, they are all obtuse (or all acute). The other three pairs of corners are of mixed types  $\alpha_{,\beta}$ ,  $\gamma'$  or  $\alpha'_{,\beta}$ ,  $\gamma'$ , where  $\alpha'_{,\beta}$ ,  $\gamma'$  are the soute (or obtuse) supplements of  $\alpha_{,\beta,\gamma}$ . If we define vectors  $\underline{\alpha}_{,\beta,\gamma}$  all have the same sign. This allows a coll to be defined uniquely in terms of the homogeneous triplet with the three shortest translations. Nithin the limits imposed by the condition that the angles be homogeneous. there are the paraible parallelophysic - and there and the other sente. Only one of these corresponds to the Brannis-reduced cell with edges cherter time our faceor bedy-diagonal.

The Deleancy reduced coll is tri-obtase and so will correspond to the Bravais reduced cell only when the latter happens to be obtase. When the Bravels-reduced cell is tri-scute, the Delaunay-reduced cell will have disgonals chorter than its adges. In the first satisf of <u>Crystel Pete</u> the Belgunay reduced cell was used as the standard cell. In 20% of the cases reported, the cell has angles that differ from 90° by more than 30°; in some cases the difference is as much as for, and there seens to be no theoretical reason may the Dekounty spin still could not produce a coll with angles approaching NO the Braveis-reduced cell is tri-acute, the Belevery reduction can be applied to the reciprocal cell (Delectors 1967), though Bassger has shown that the cherters of the in reciprecel space do not necessarily correspond to the cherolest A.h.c. in direct space.

The Birichlet reduction vill always produce the coll with the three shortest translations. Given any bright, g.g.g. it forms the reduced triplet by point subbigles of each voctor in turn in attactive to should the other two edges. The stops can be perfound to way order and they are repeated until there is no further change. Ther order convention adopted (a<b<c or c<a>b) has no effect on the reduction.

In the second edition of <u>Gryatal Mate</u> the Braveisreduced cell is used as the standard cell in all cases; but the condition imposed is that  $\alpha$  and  $\beta$  be non-acute. In the case of a tri-acute parallelopiped, this corresponds to a shift of origin to the "right-handed" corner whose angles are  $\alpha^{\circ}$ ,  $\beta^{\circ}$ ,  $\gamma_{\circ}$ .

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